Introduction to Incoherent Scatter 1

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Incoherent Scatter Radar

- Incoherent
- Scatter
- Radar

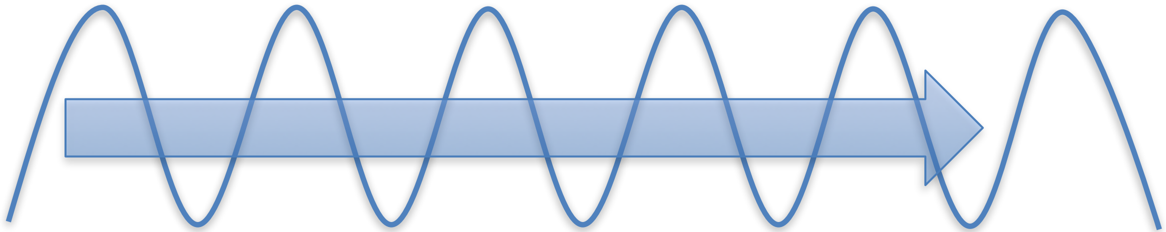
Incoherent Scatter Radar

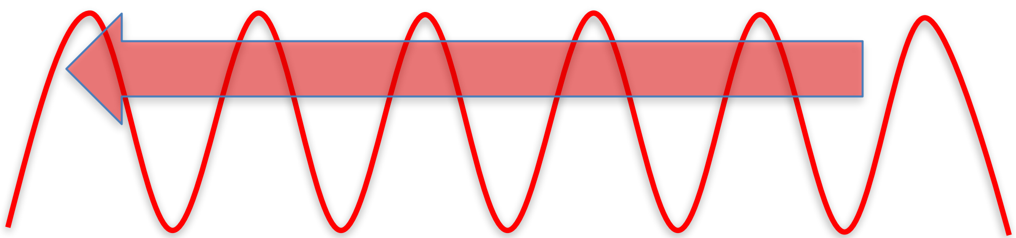
- Radar
- Scatter
- Incoherent

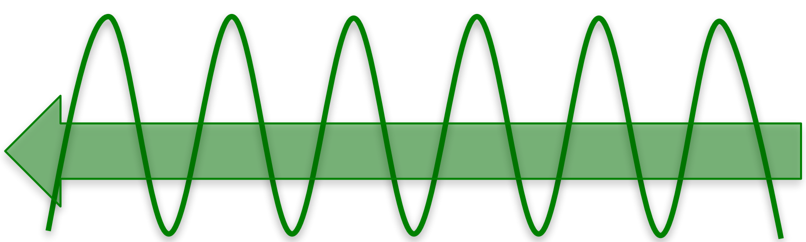
Radar

- RADAR (RAdio Detection And Ranging)
  - is a technique for detecting and studying remote targets by transmitting a radio wave in the direction of the target and observing the reflection of the wave.
  - **Radar** is an object detection system which uses radio waves to determine the range, altitude, direction, or speed of objects. (wikipedia)

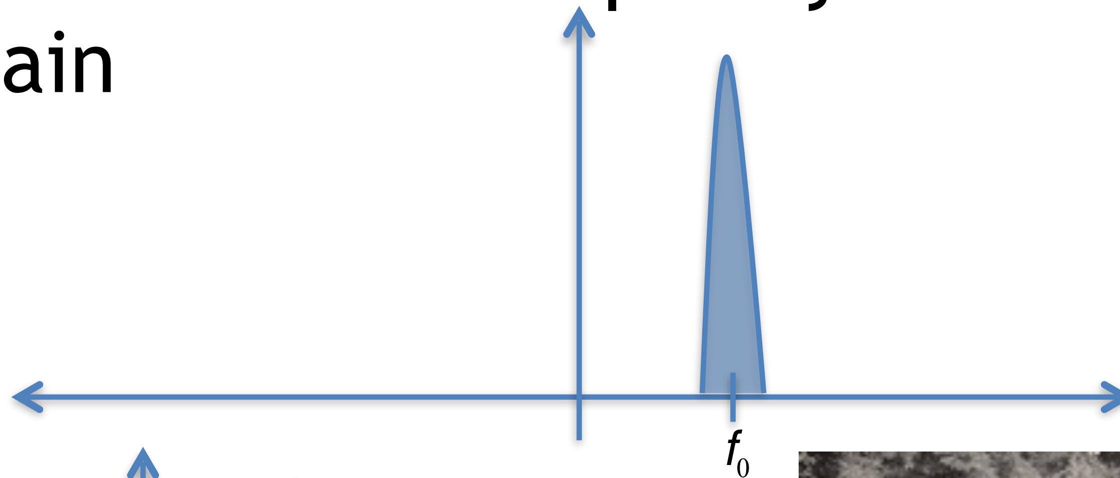
Doppler Radar – time domain

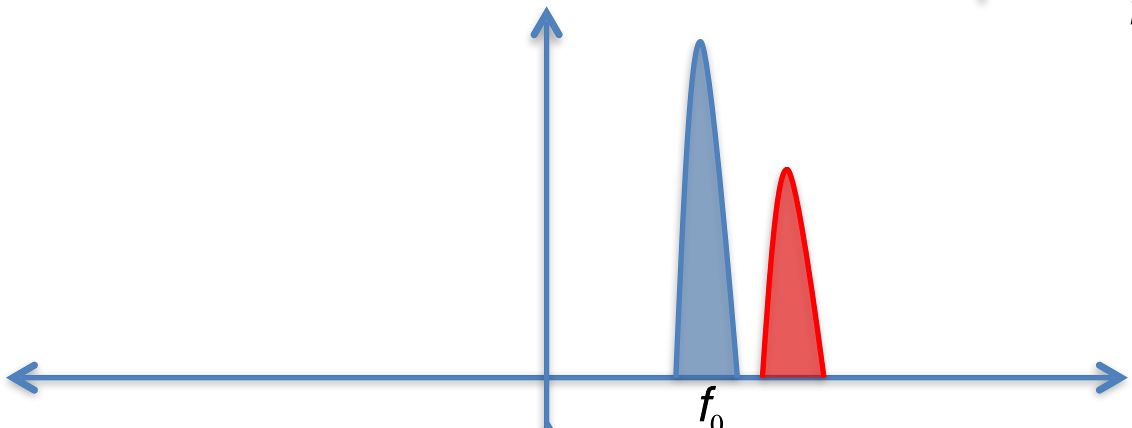


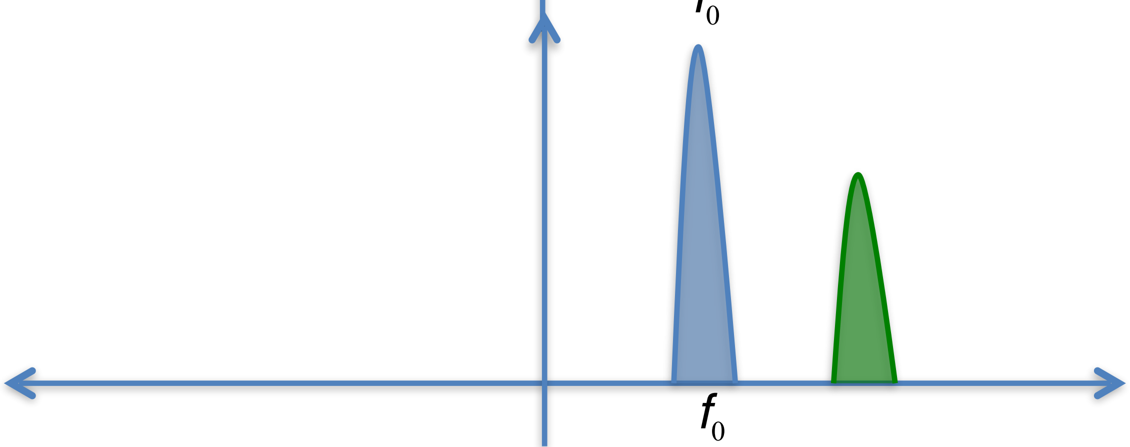


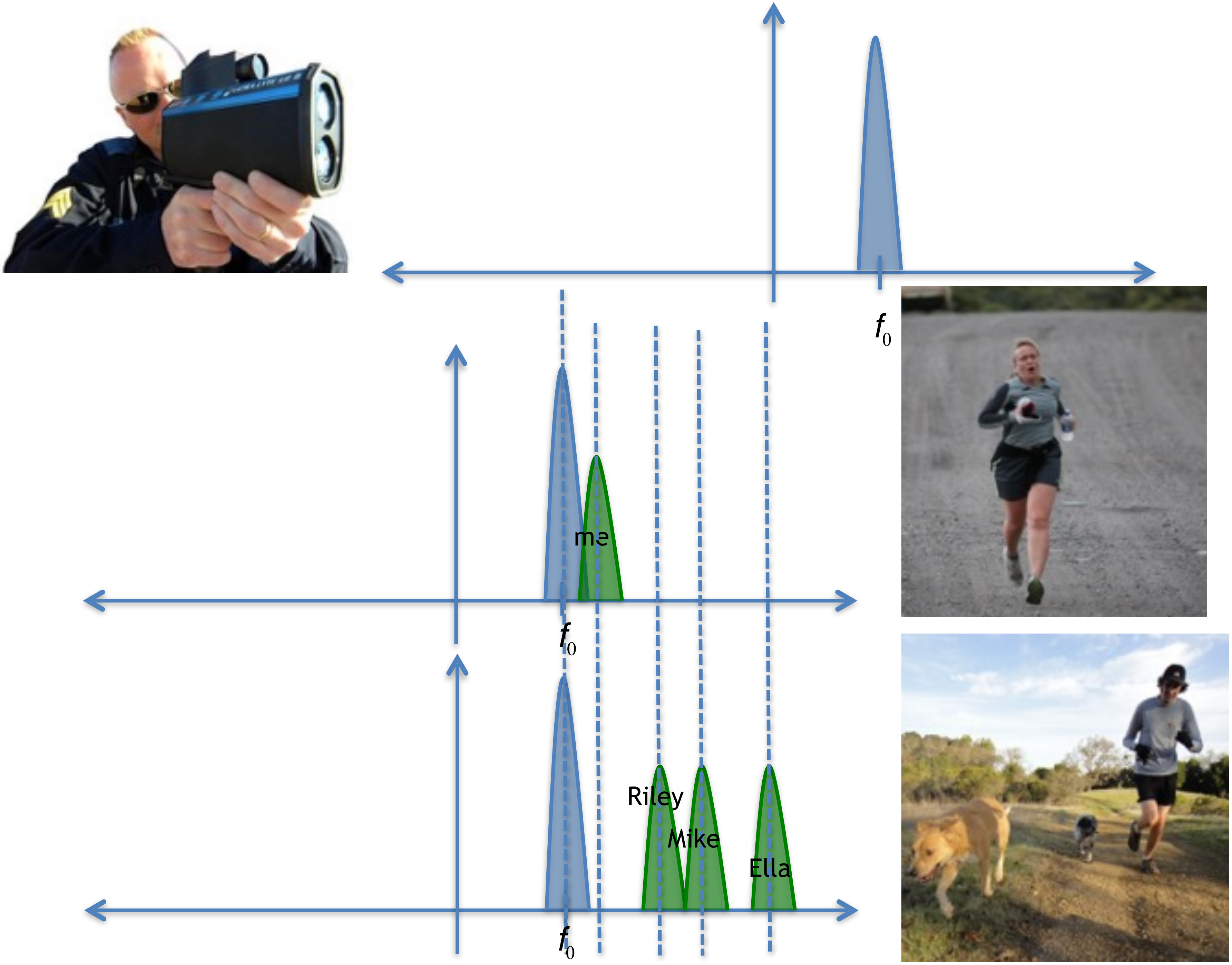


Doppler Radar – frequency domain

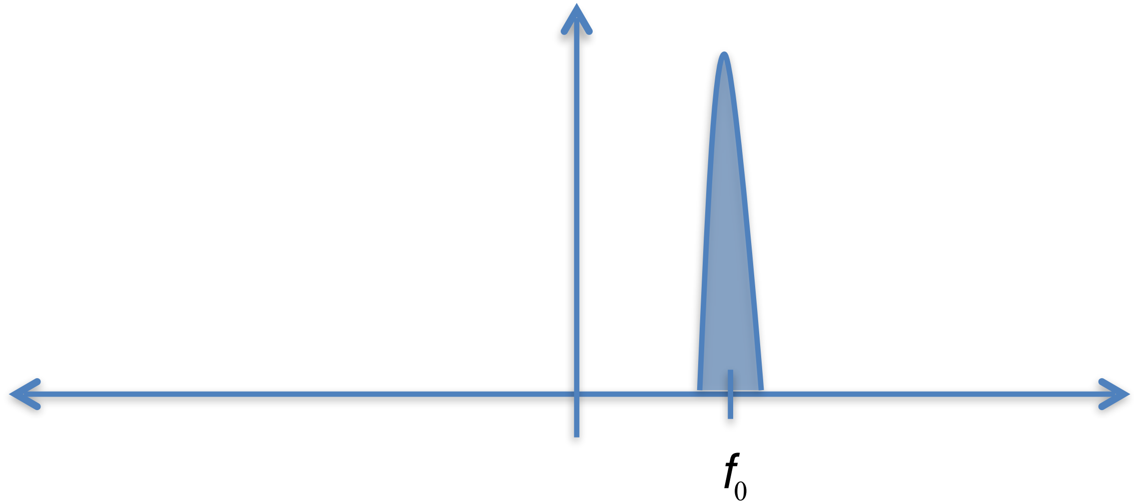


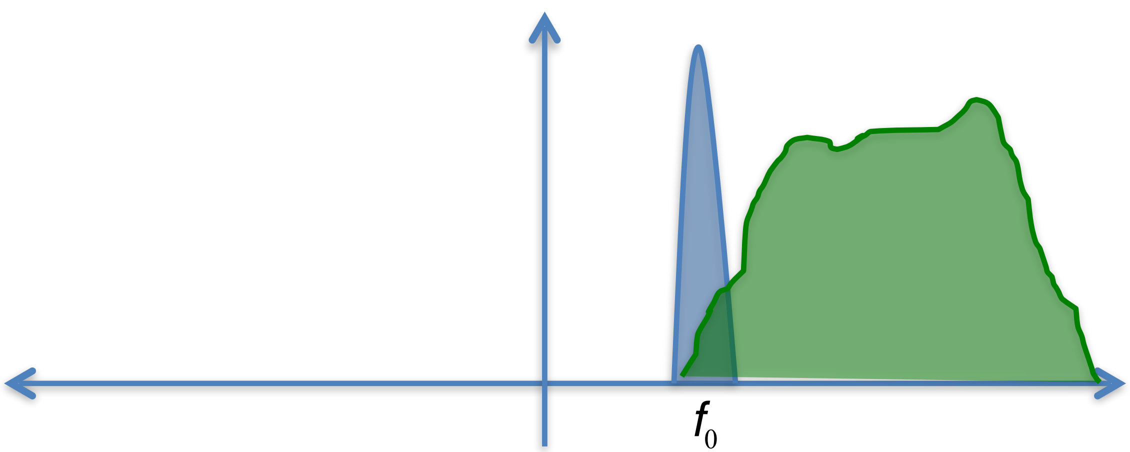






f_0
me
f_0
Riley
Mike
Ella
f_0


f_0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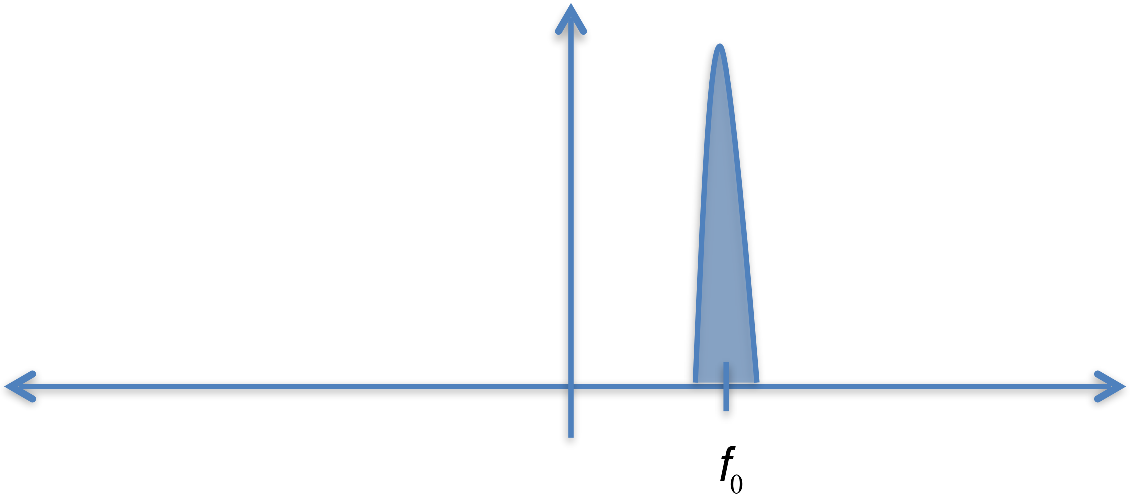
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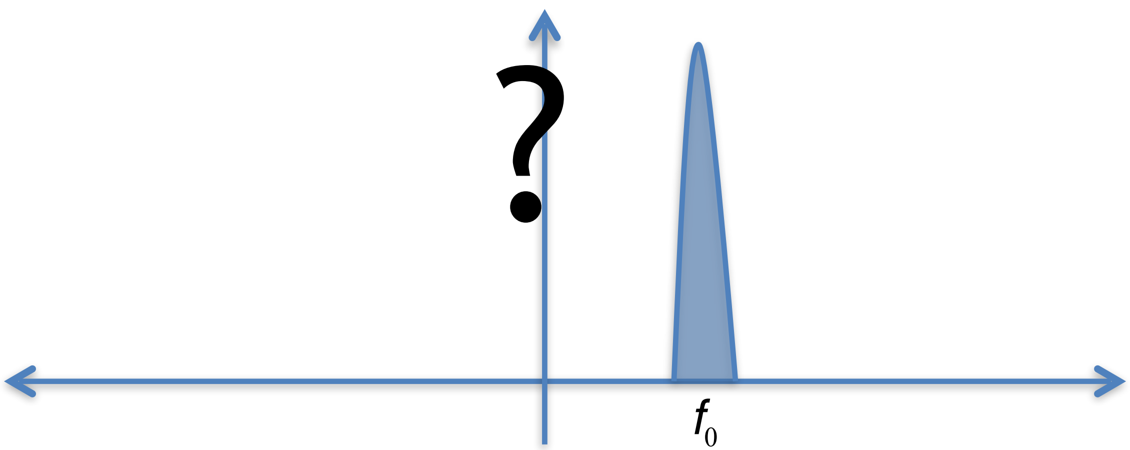
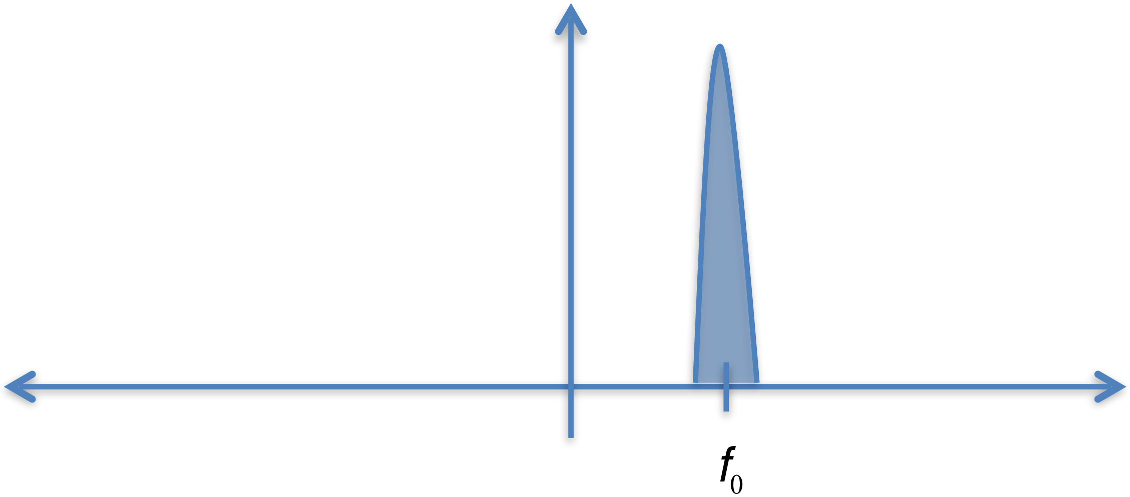

?
f_0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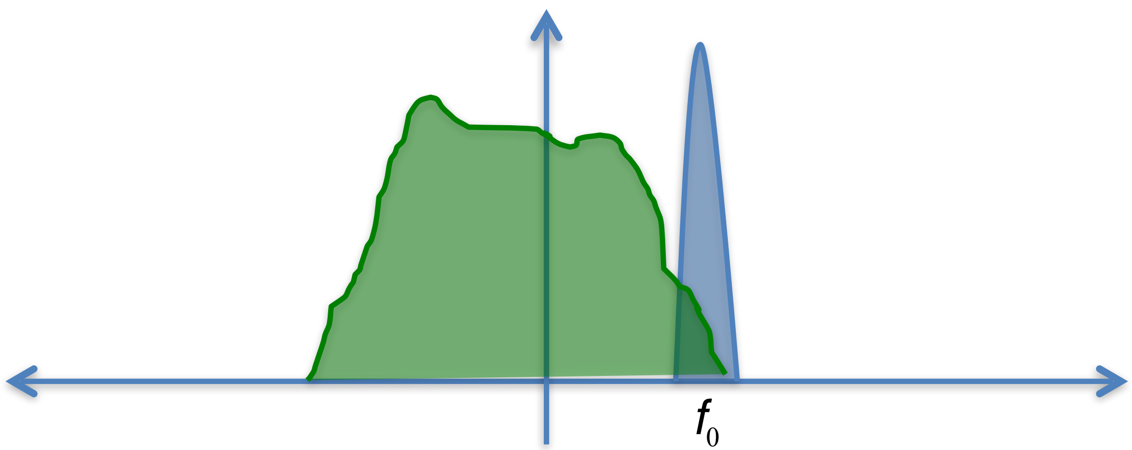

f_0


START
START
START

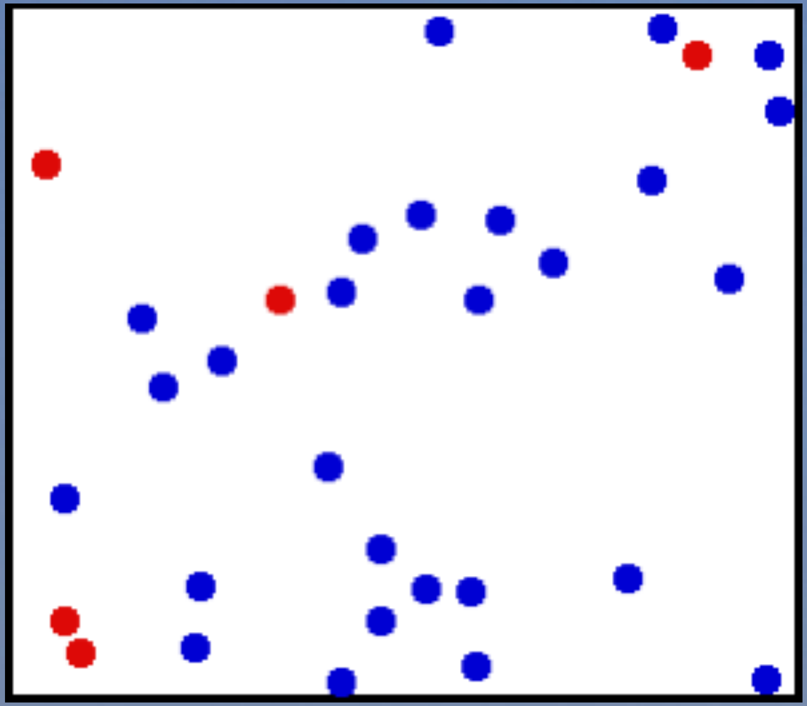
Thomson scattering

- Thomson scattering is the elastic scattering of electromagnetic radiation by a free charged particle, as described by classical electromagnetism.

- In the low-energy limit, the electric field of the incident wave (radar wave) accelerates the charged particle, causing it, in turn, to emit radiation at the same frequency as the incident wave, and thus the wave is scattered.

- As long as the motion of the particle is non-relativistic (i.e. its speed is much less than the speed of light), the main cause of the acceleration of the particle will be due to the electric field component of the incident wave, and the magnetic field can be neglected. The particle will move in the direction of the oscillating electric field, resulting in electromagnetic dipole radiation.

Thermal fluctuating electrons



Radar Equations

Hard target:

$$P_r = \frac{P_t G_t A_r \sigma}{(4\pi)^2 R^4}$$

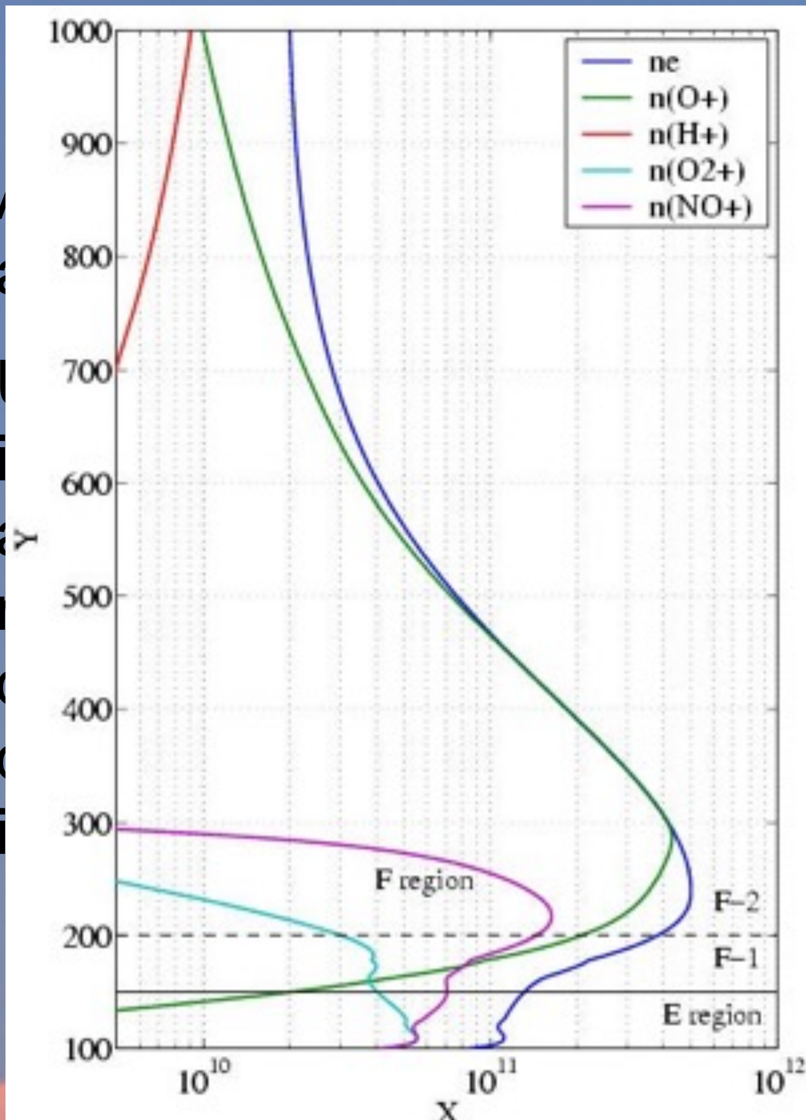
Incoherent Scatter Radar

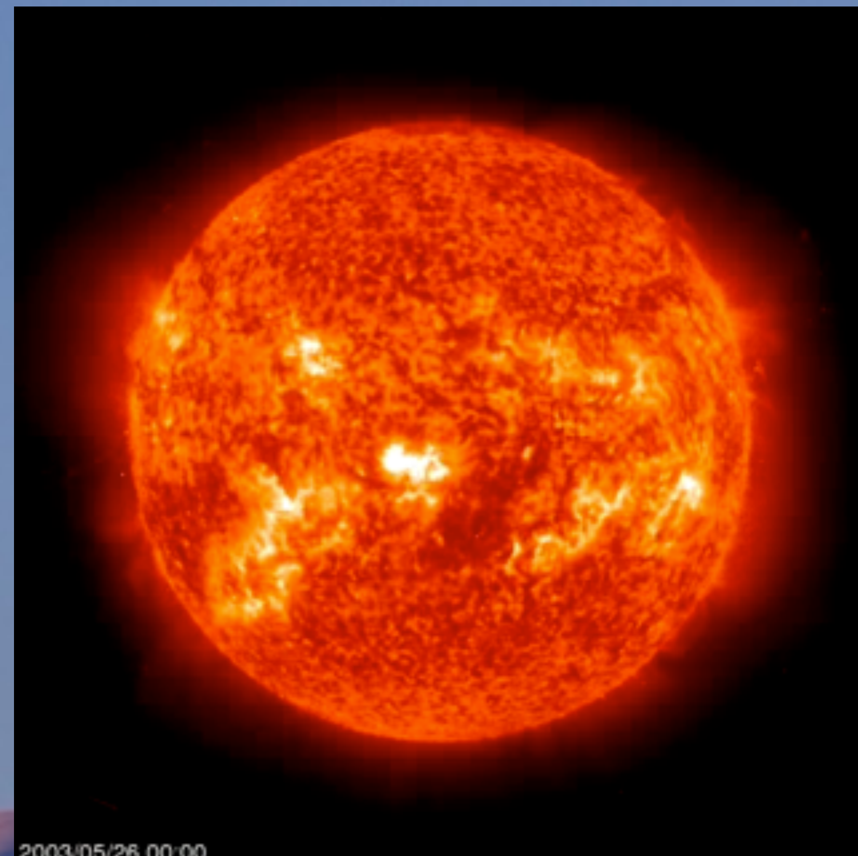
$$P_r = \frac{C c_0 G \lambda^2}{2(4\pi)^2} \frac{P_t \tau_p}{R^2} \frac{\sigma_e n_e(R)}{(1+k^2\lambda_D^2)(1+k^2\lambda_D^2+T_r)}$$

The ionospheric plasma

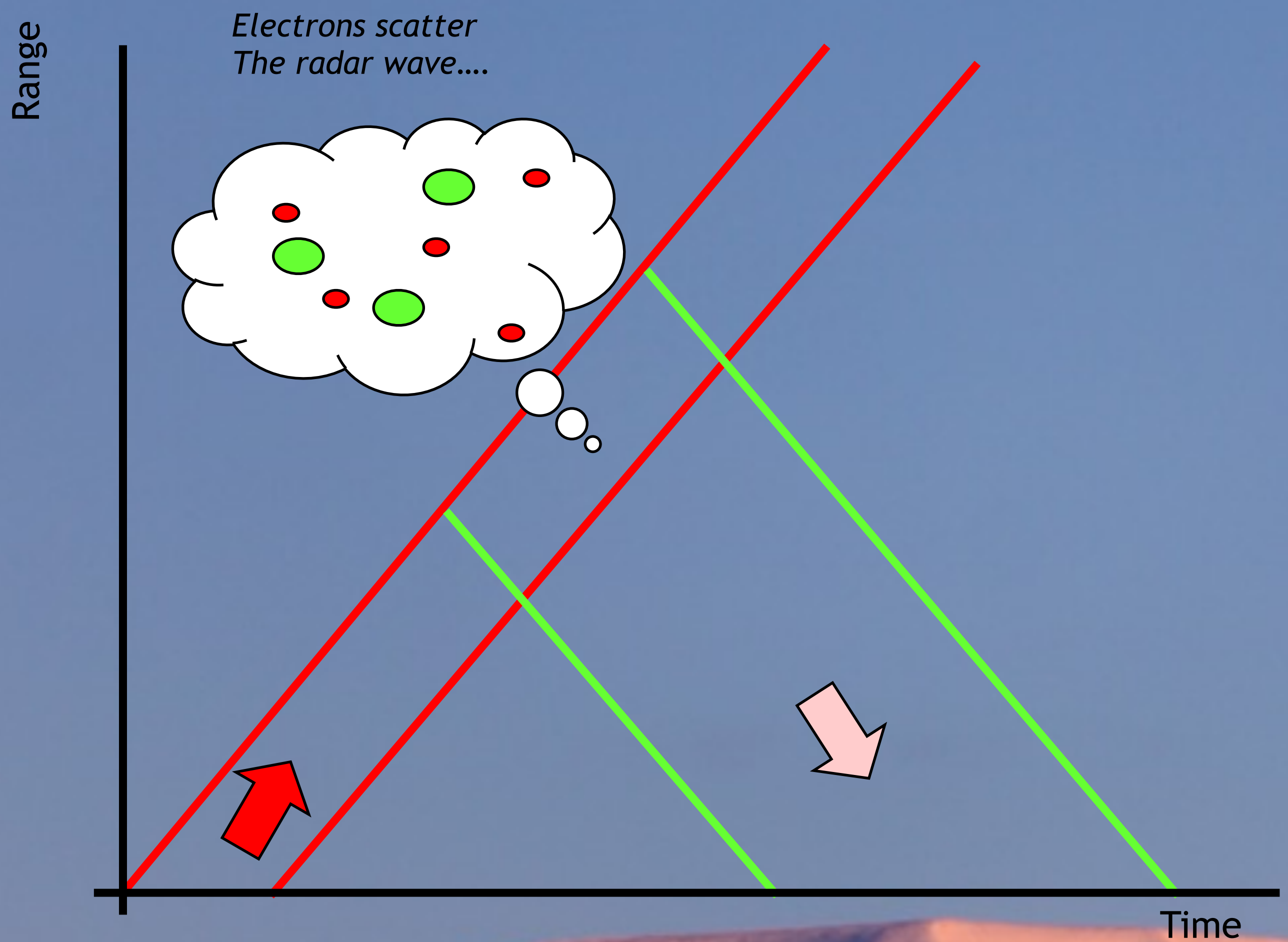
- So now we need electrons in the ionosphere to scatter the radar wave off…

…what Anita said this morning…





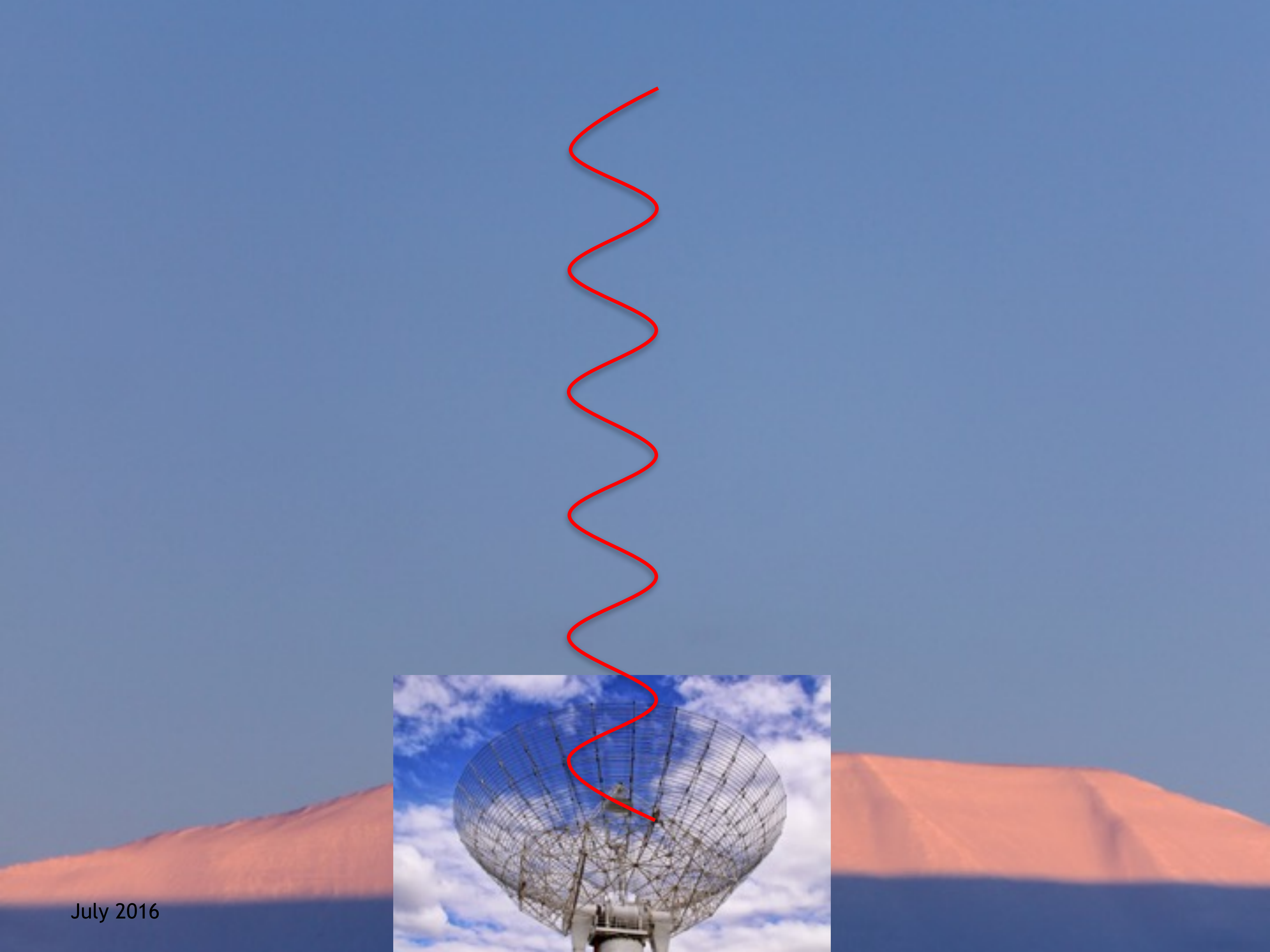
How ISRs work...

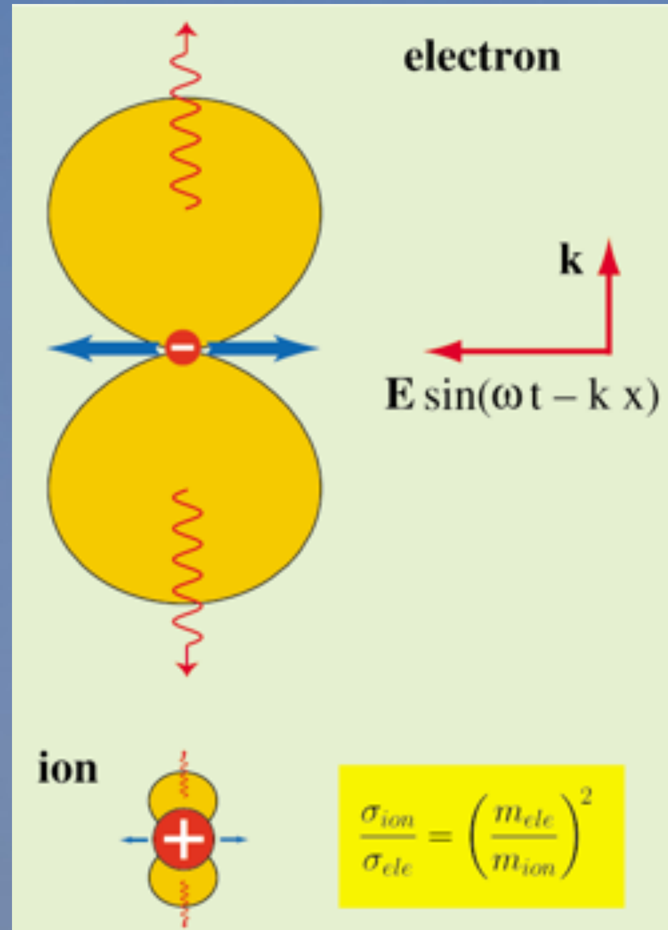


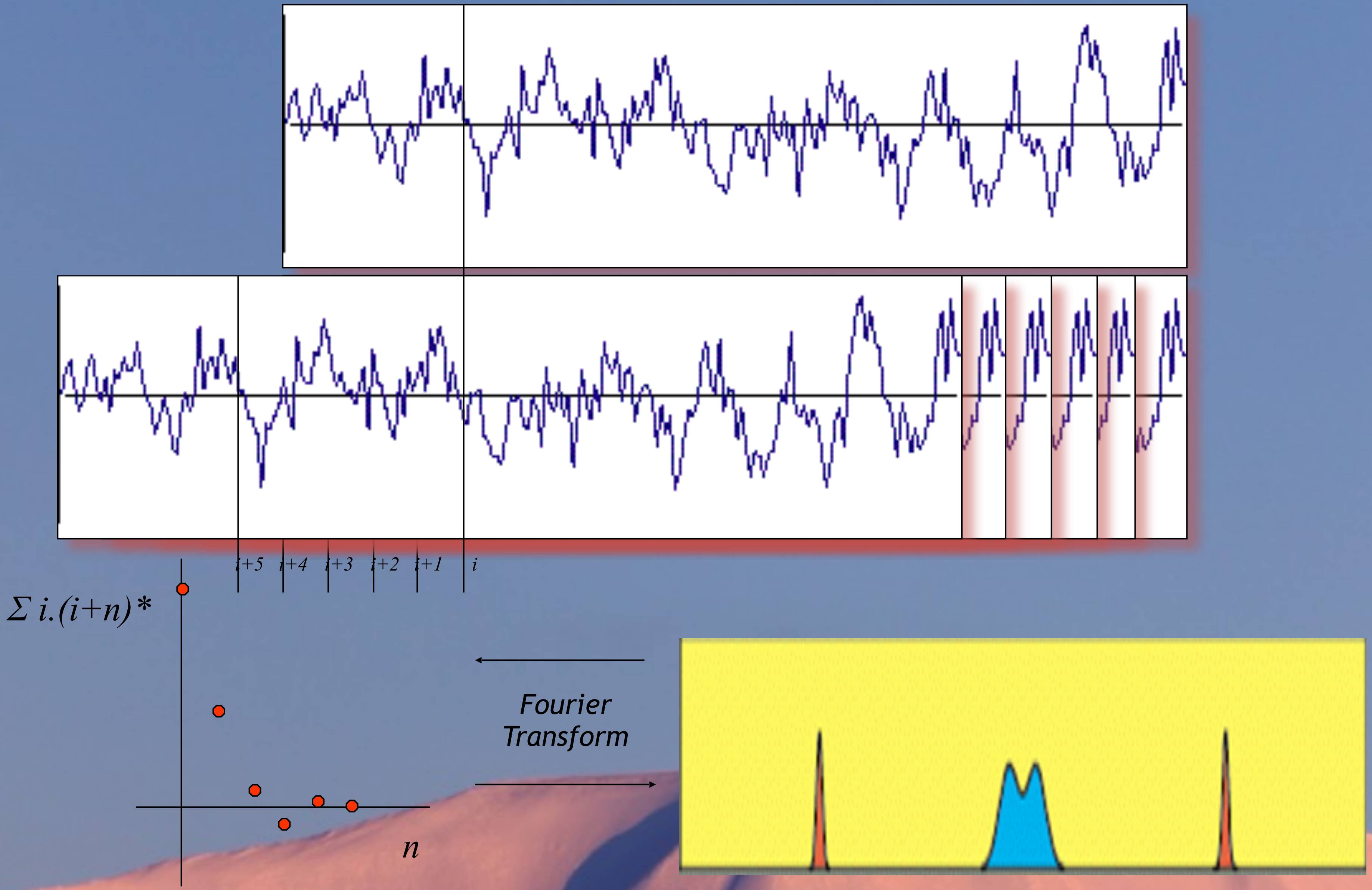
High power transmitter

Very sensitive receiver

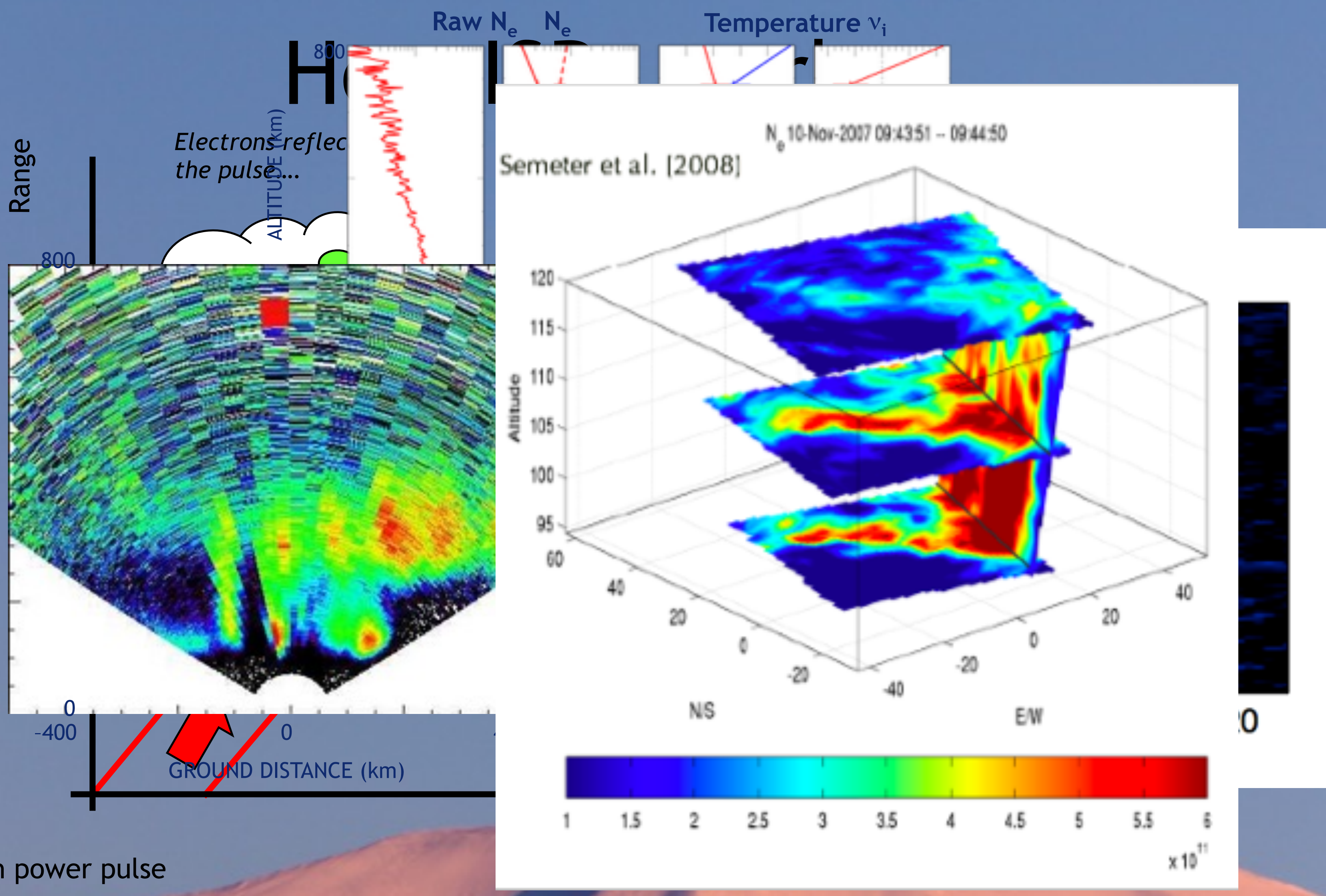
Only ~0.000000000000000001% of the transmitted power is returned!




electron
k
E sin(ω t – k x)
ion
$\frac{\sigma_{ion}}{\sigma_{ele}} = \left(\frac{m_{ele}}{m_{ion}}\right)^2$



i+5
i+4
i+3
i+2
i+1
i
Σ i.(i+n)*
n
Fourier
Transform

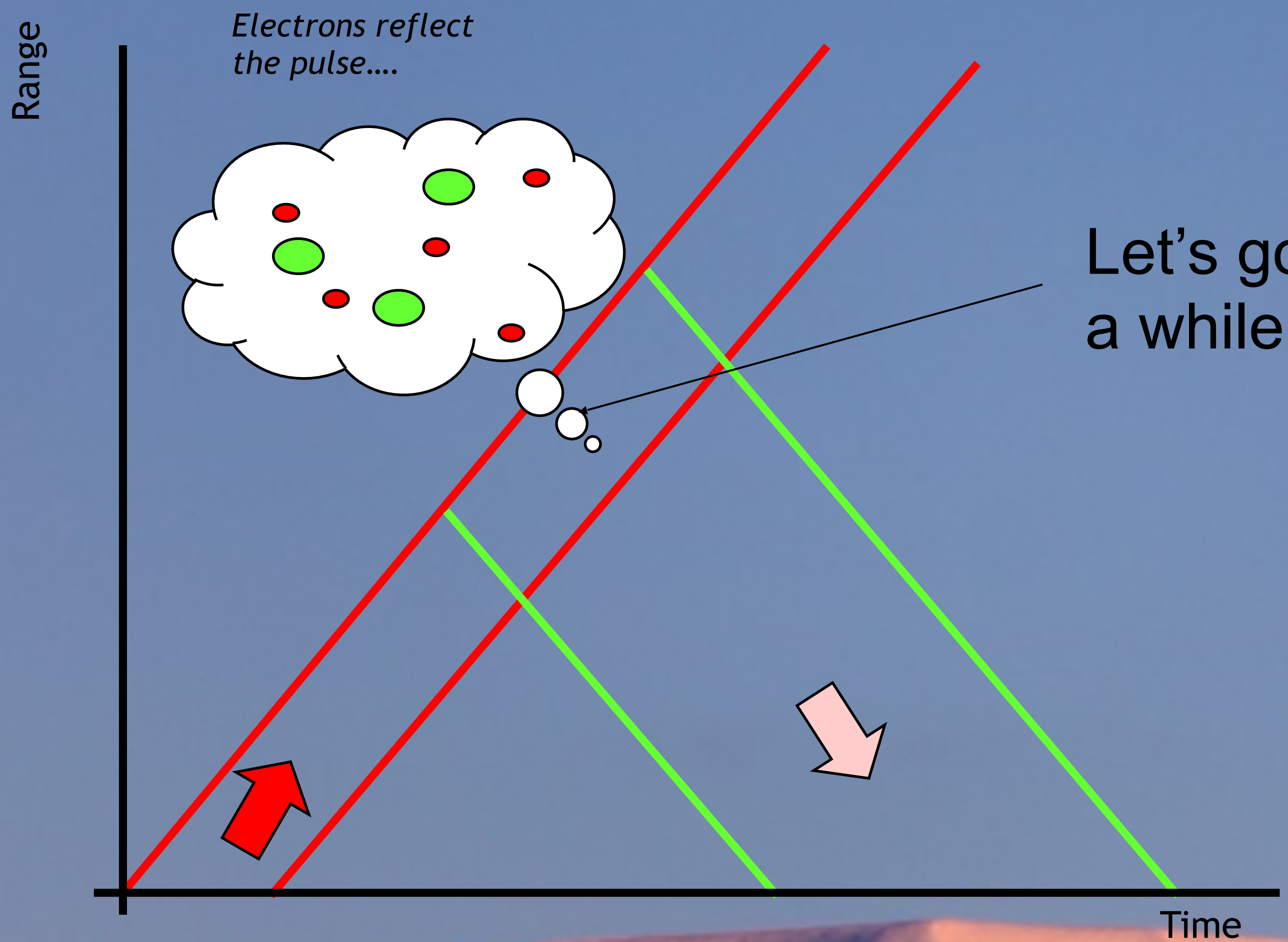


High power pulse

Only ~0.0000000000000000001% of the transmitted power is returned!

With incoherent Scatter Radars we can determine statistical properties of the charged particle distributions

How ISRs work…



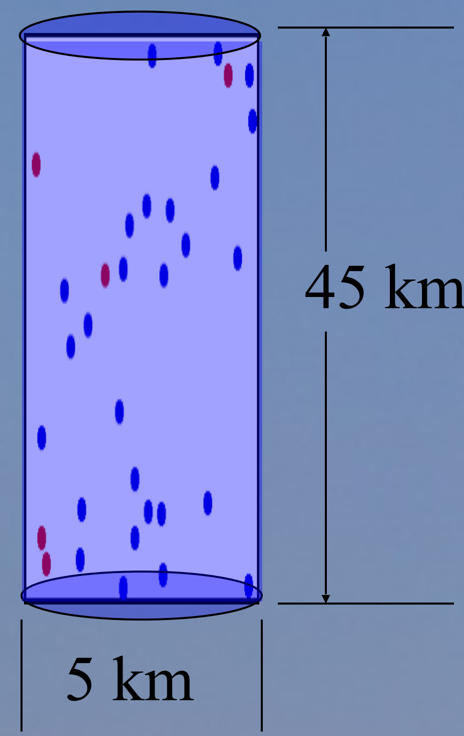
Let's go back here for a while…

High power pulse

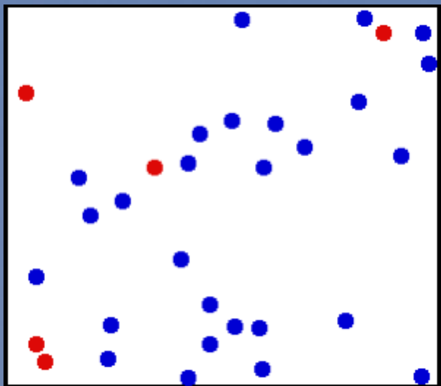
Very sensitive receiver

Total cross section estimate:

Consider an antenna with a 1-degree beam measuring the ionospheric plasma at 300 km range and using a 300 microsecond pulse.
If the electron density is 10^{12} m^{-3}, the total number of electrons scattering into a given measurement is $\sim 8.8 \times 10^{23}$. This yields a total cross-section of 88 mm^2 – we need a big radar!



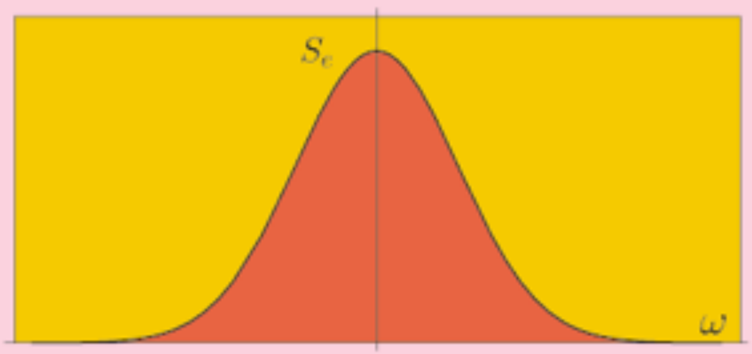
For TRUE incoherent scatter…



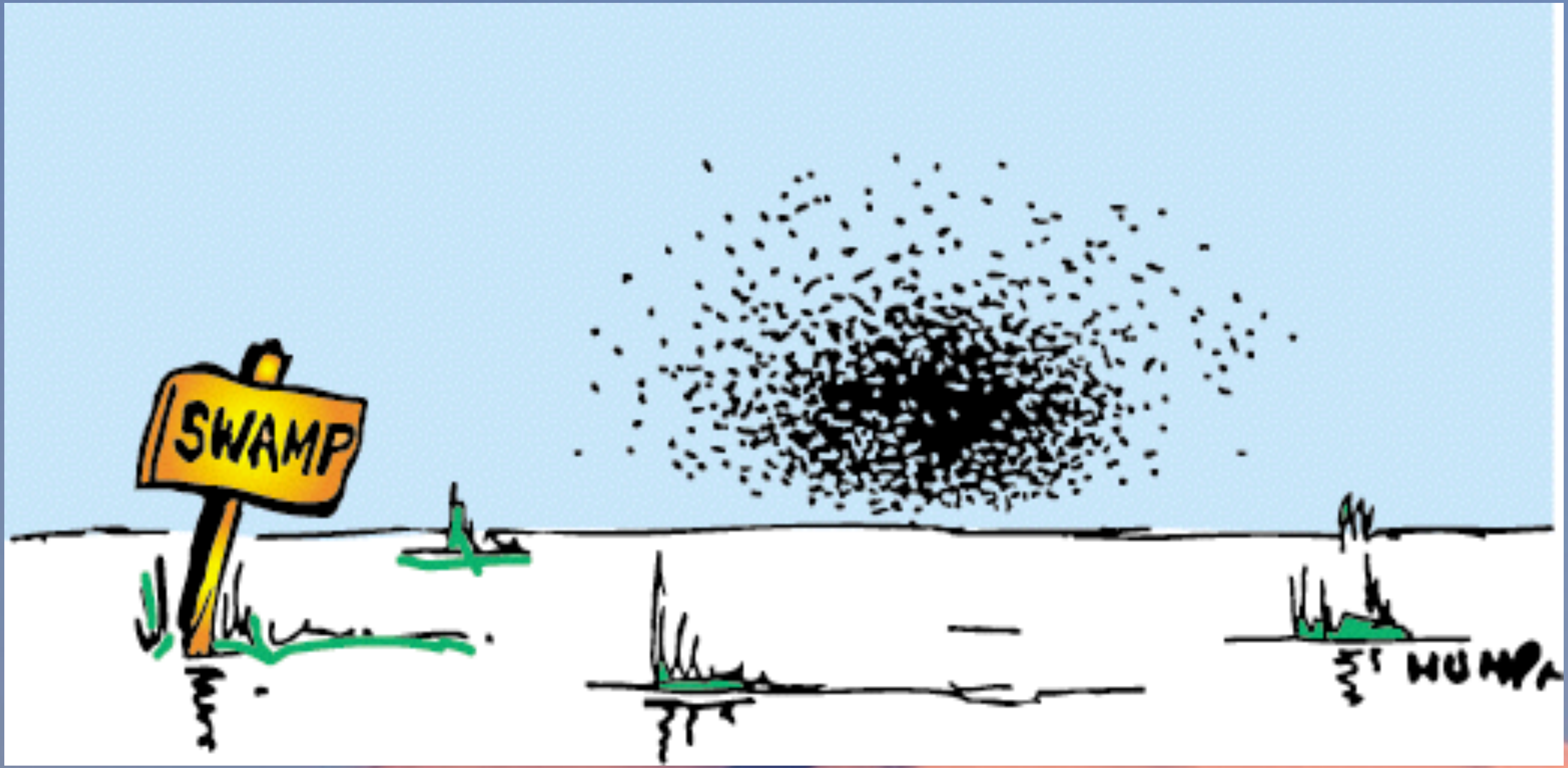
no collective interactions

$$S_e(\mathbf{k},\omega) = N_e \left|1 - \cancel{\frac{\chi_e(\mathbf{k},\omega)}{\epsilon(\mathbf{k},\omega)}}\right|^2 \int d\mathbf{v}\, f_e(\mathbf{v})\delta(\omega - \mathbf{k}\cdot\mathbf{v}) + N_i \left|\cancel{\frac{\chi_e(\mathbf{k},\omega)}{\epsilon(\mathbf{k},\omega)}}\right|^2 \int d\mathbf{v}\, f_i(\mathbf{v})\delta(\omega - \mathbf{k}\cdot\mathbf{v})$$

$$S_e(\mathbf{k},\omega) = N_e \int d\mathbf{v}\, f_e(\mathbf{v})\delta(\omega - \mathbf{k}\cdot\mathbf{v})$$



Incoherent scattering - the short story



Incoherent scattering - the short story

- We only see scattering from the electrons

  …but they also tell the story about the ion dynamics…

Collective behavior…

- There are a number of wave modes existing inherently in the ionospheric plasma…

Langmuir waves

- High frequency electrostatic waves
- Dispersion relation:

$$\omega_r = \left(\omega_{pe}^2 + 3k^2 v_{the}^2\right)^{1/2} = \omega_p \left(1 + 3k^2 \lambda_{De}^2\right)^{1/2}$$

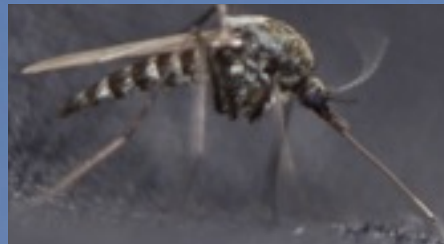
$$\omega_i = -C \frac{\omega_{pe}}{(k\lambda_{De})^3} \exp\left(-\tfrac{1}{2} k^2 \lambda_{De}^2\right), \quad C = \sqrt{\frac{\pi}{8}}\, e^{-3/2}$$

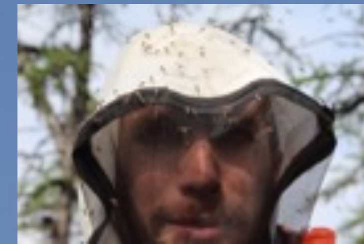
Ion acoustic waves

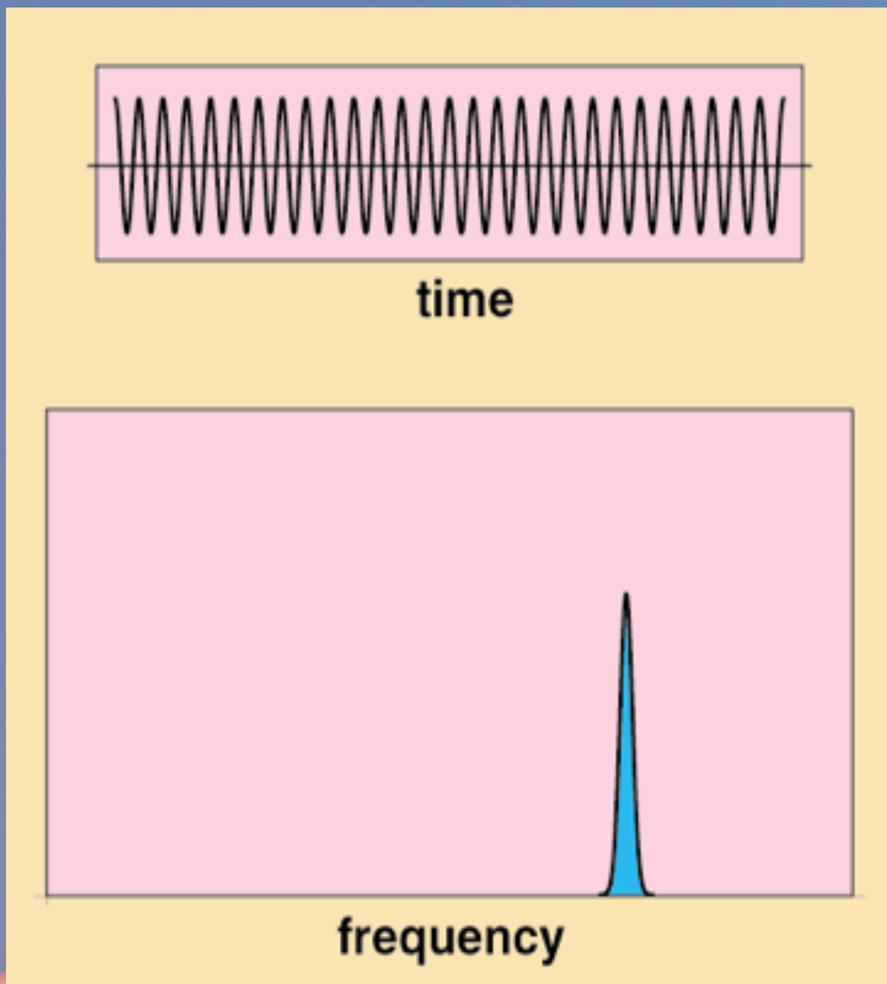
* Ion acoustic waves:

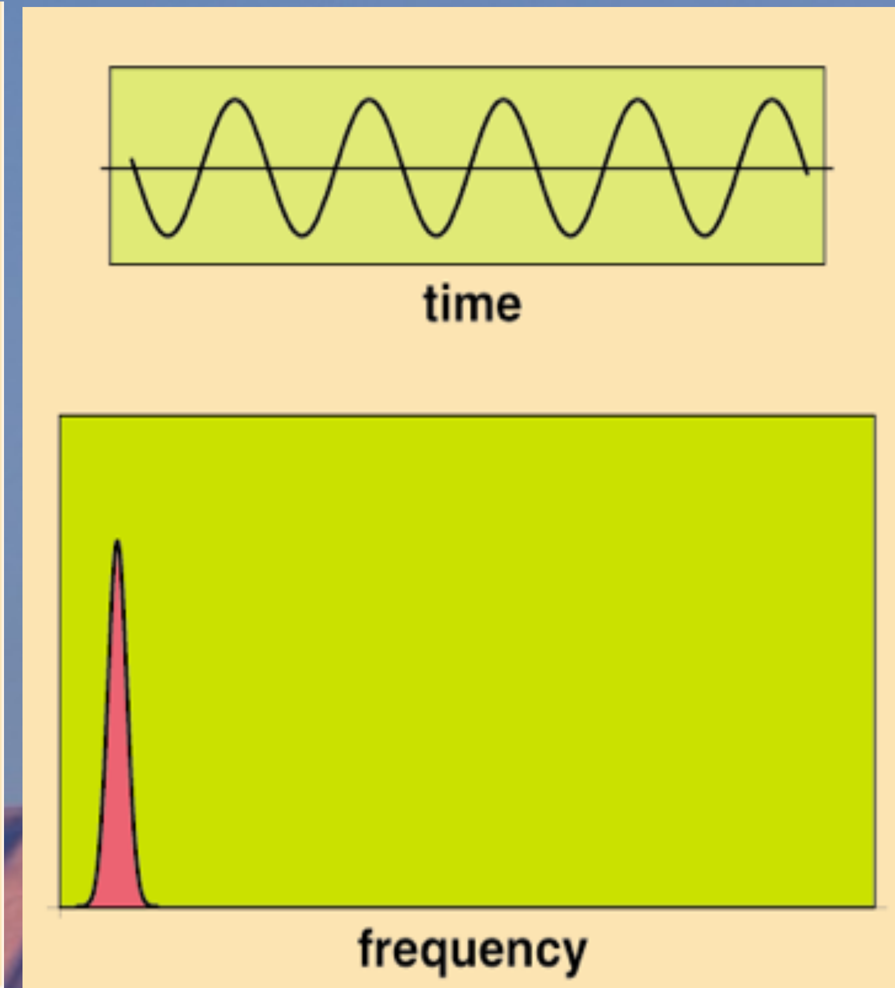
(63) $$\omega_r = \frac{k c_s}{1 + k^2 \lambda_{De}^2}, \qquad c_s = \left(\frac{k_B T_e + 3 k_B T_i}{m_i} \right)^{1/2}$$

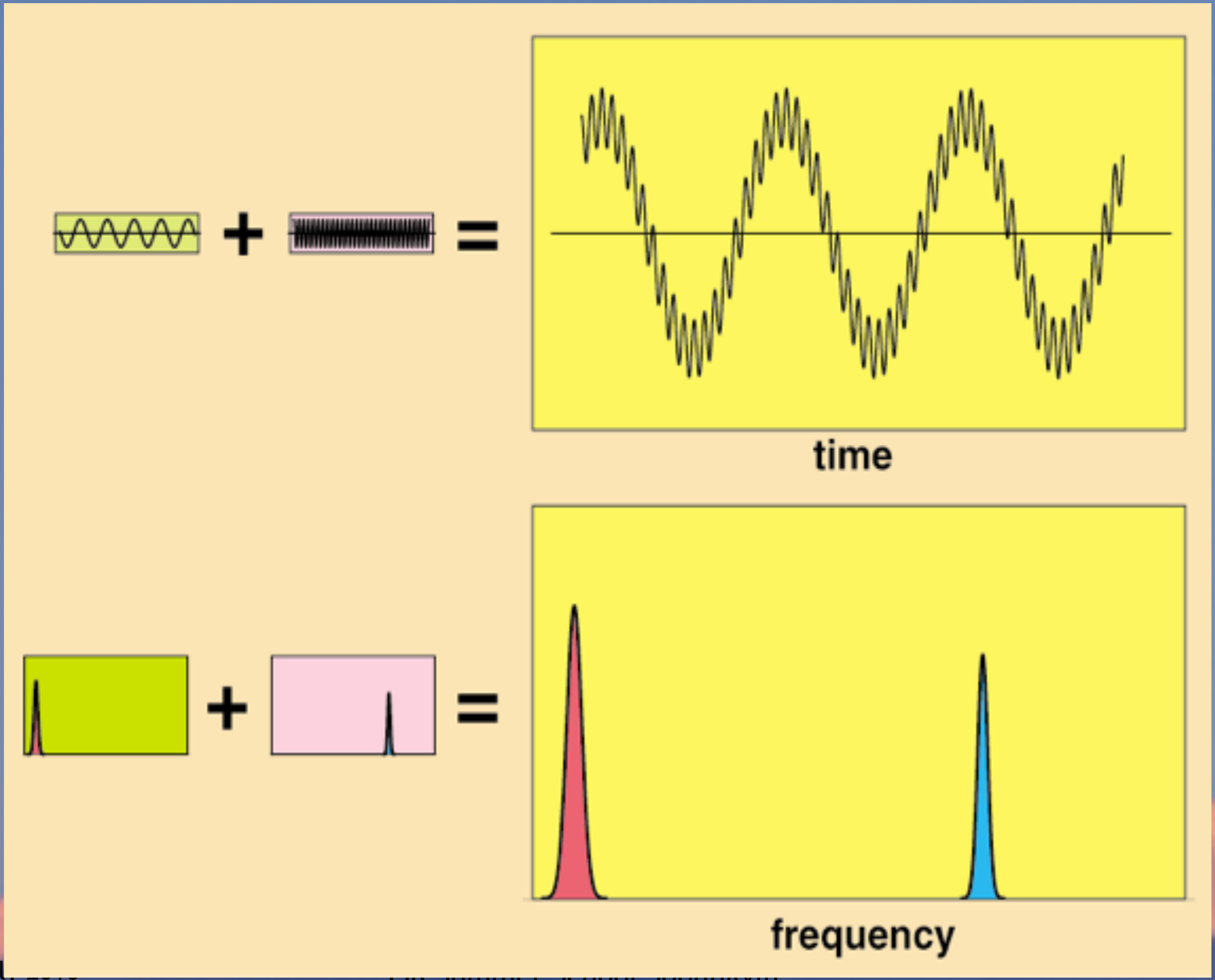
(64) $$\omega_i = -\sqrt{\frac{\pi}{8}} \frac{\omega_r}{(1 + k^2 \lambda_{De}^2)^{3/2}} \left[\left(\frac{T_e}{T_i} \right)^{3/2} \exp\left(- \frac{T_e/T_i}{2(1 + k^2 \lambda_{De}^2)} \right) + \sqrt{\frac{m_e}{m_i}} \right]$$



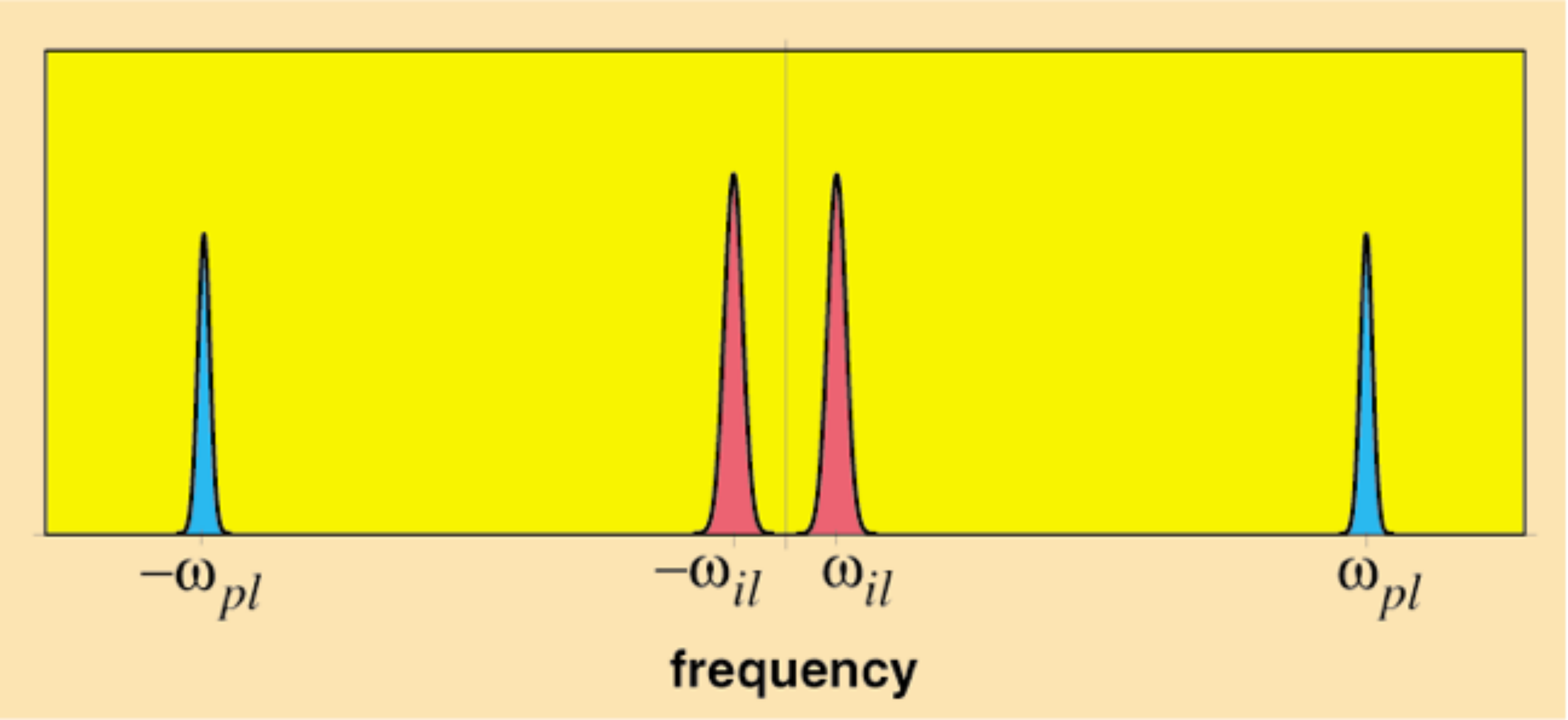



time
frequency

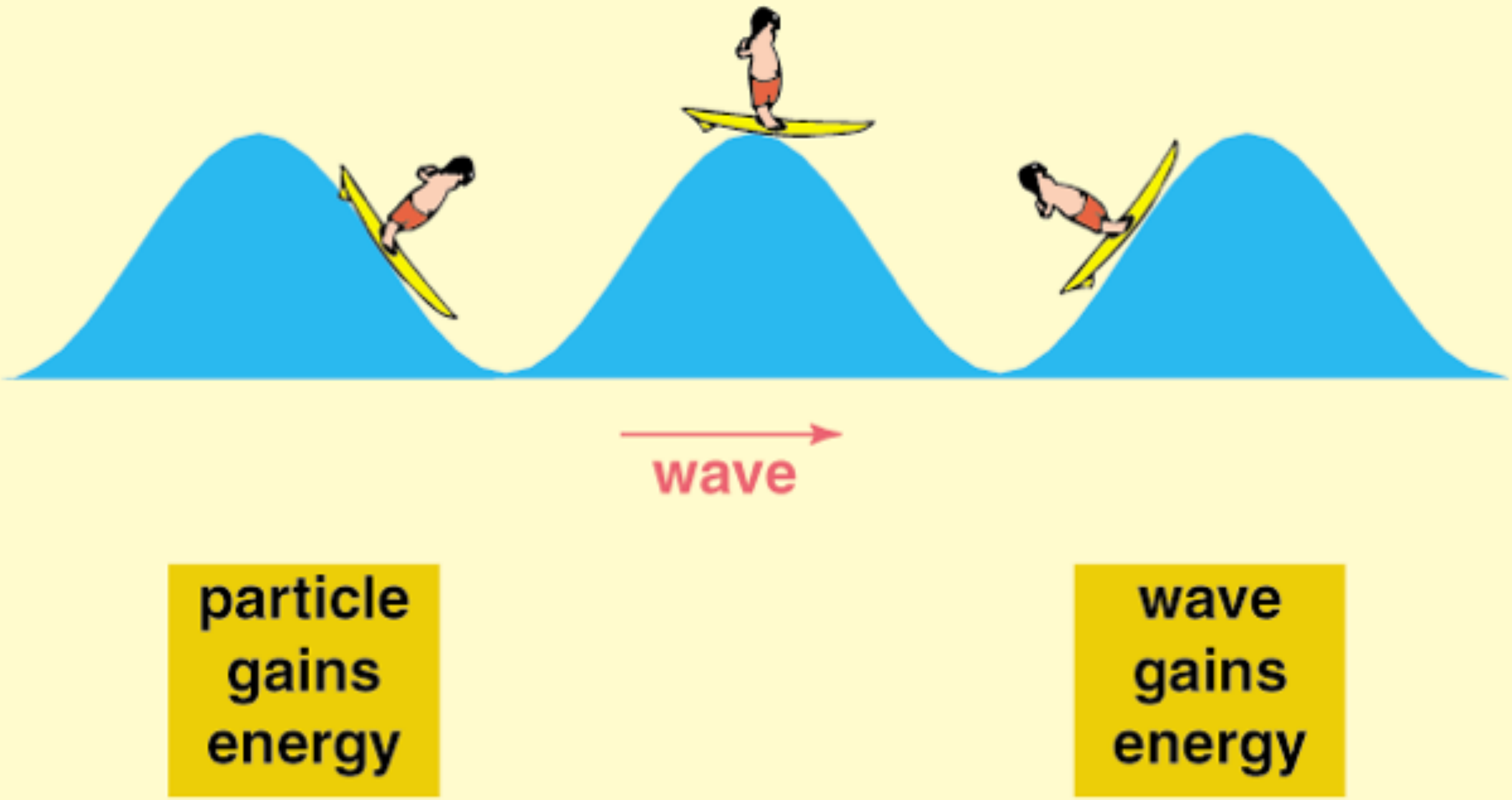

time
frequency

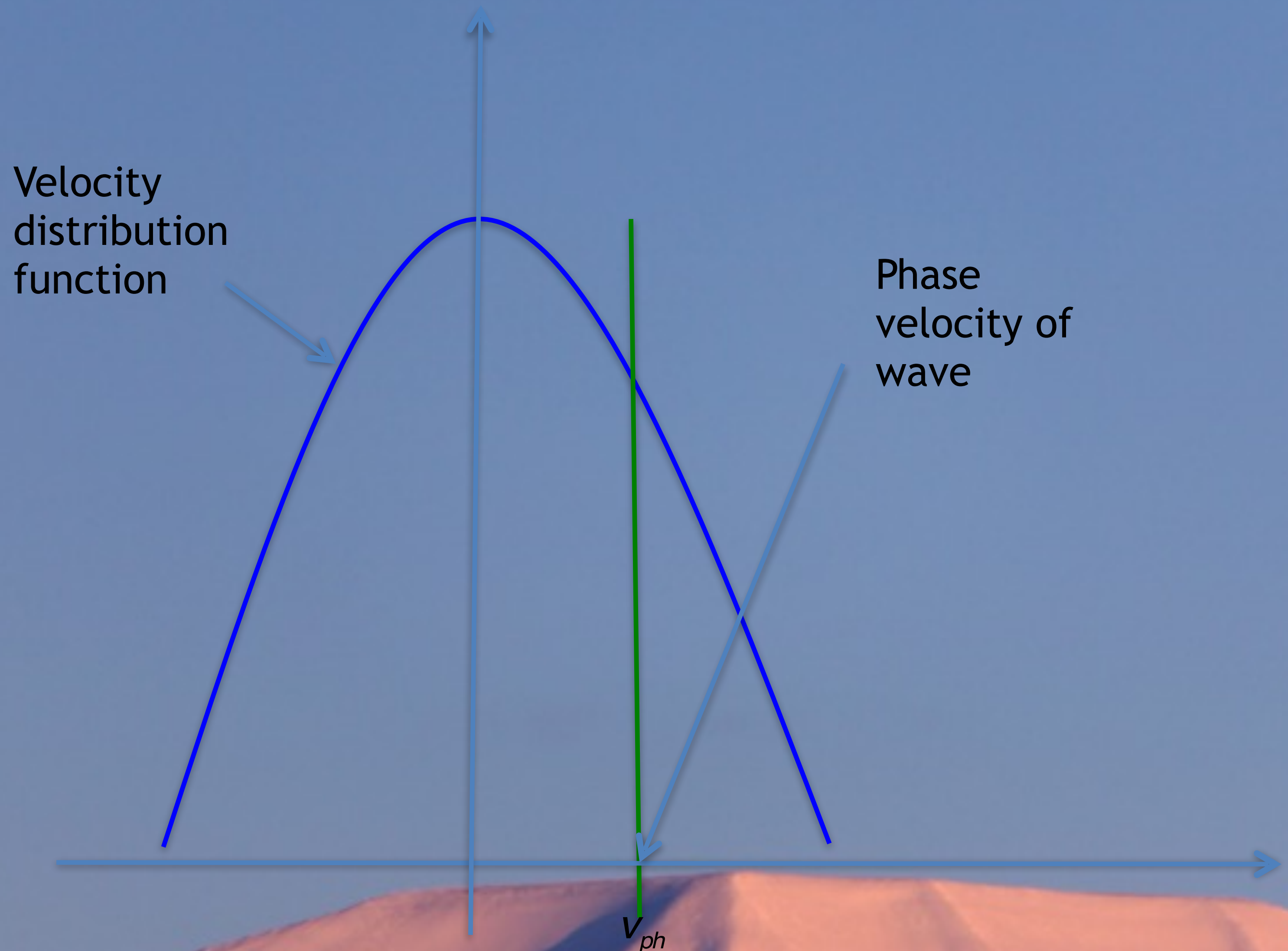

time
frequency

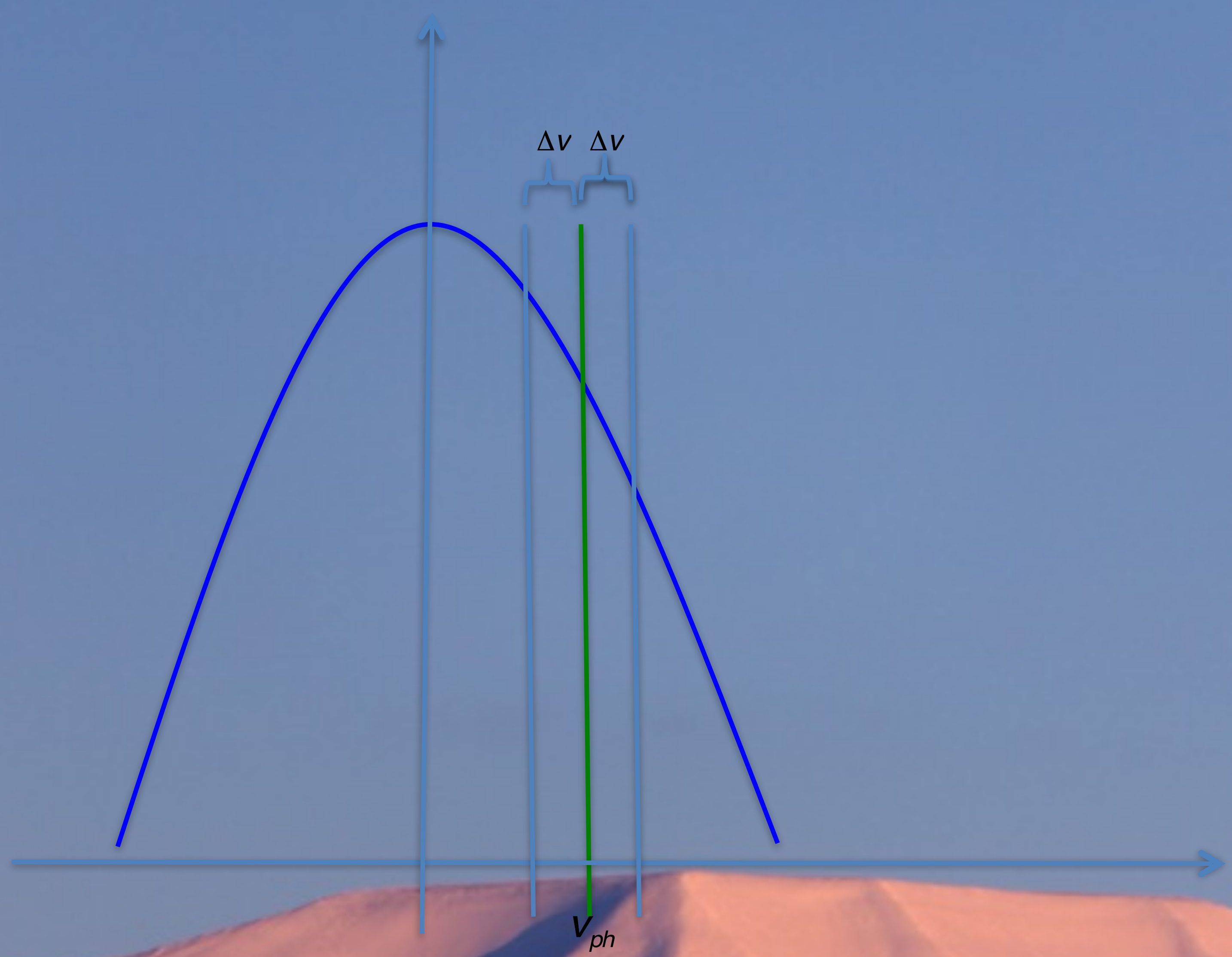
Plasma Wave Approach (cont'd)

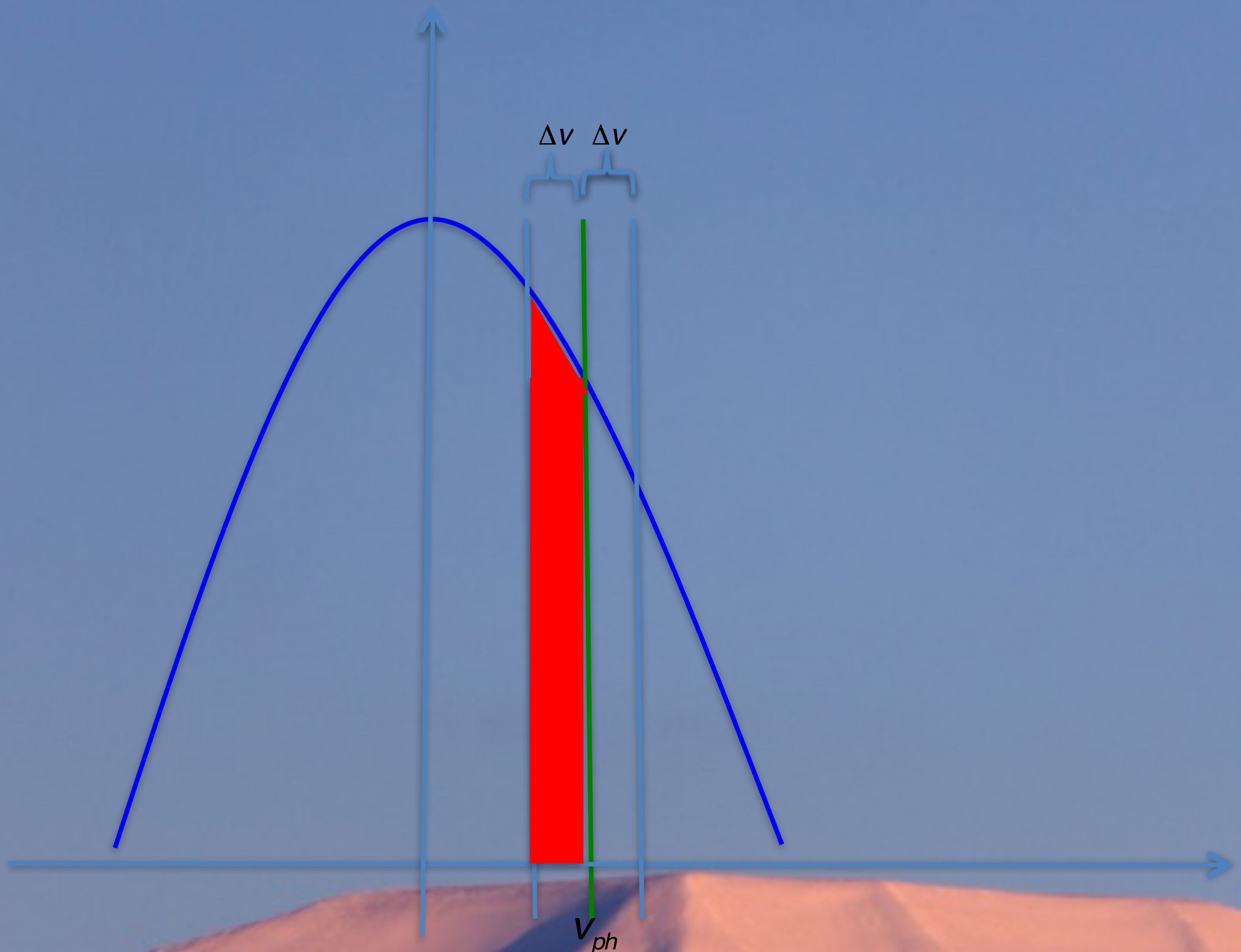


Landau wave-particle interactions

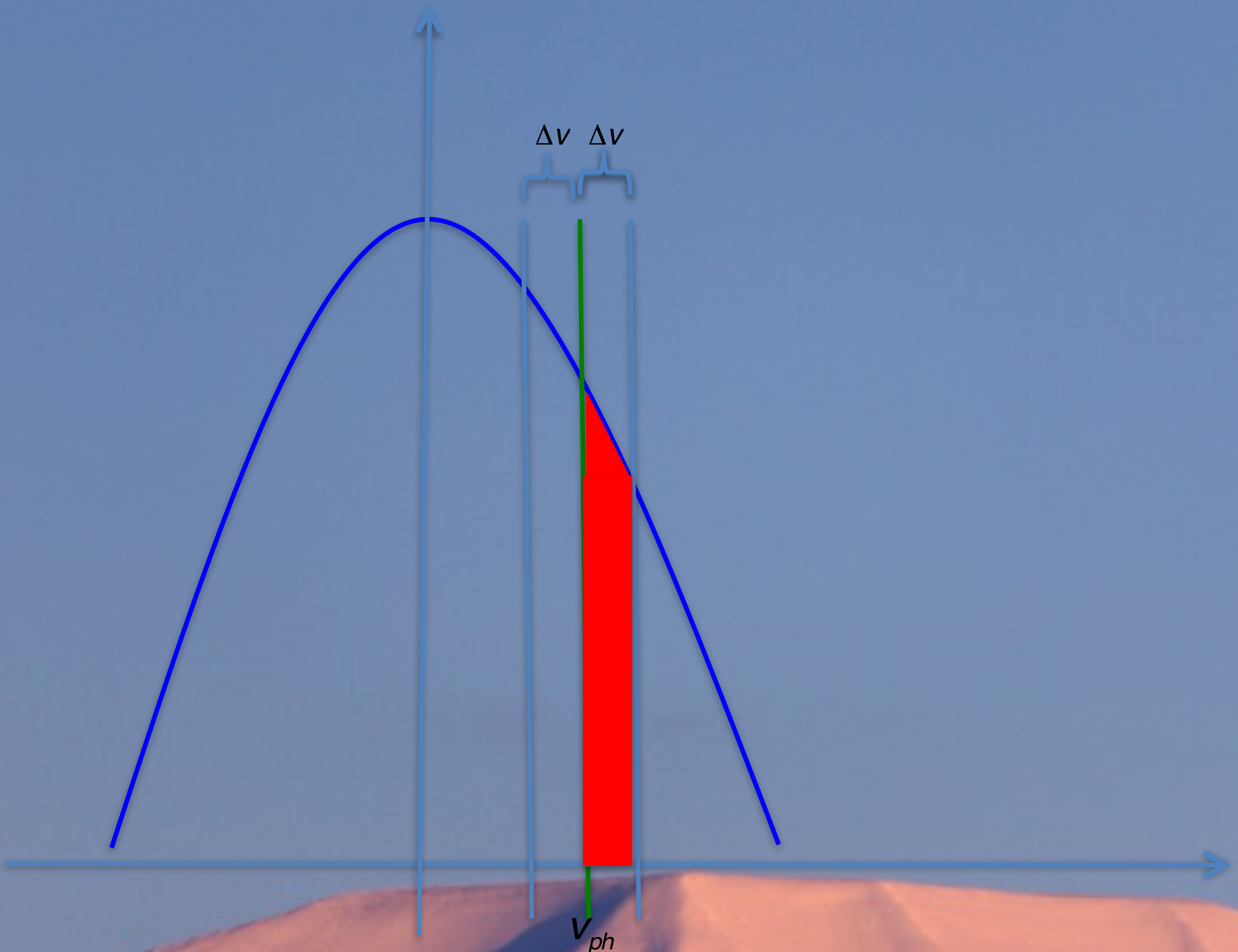



Velocity distribution function
Phase velocity of wave
v_{ph}

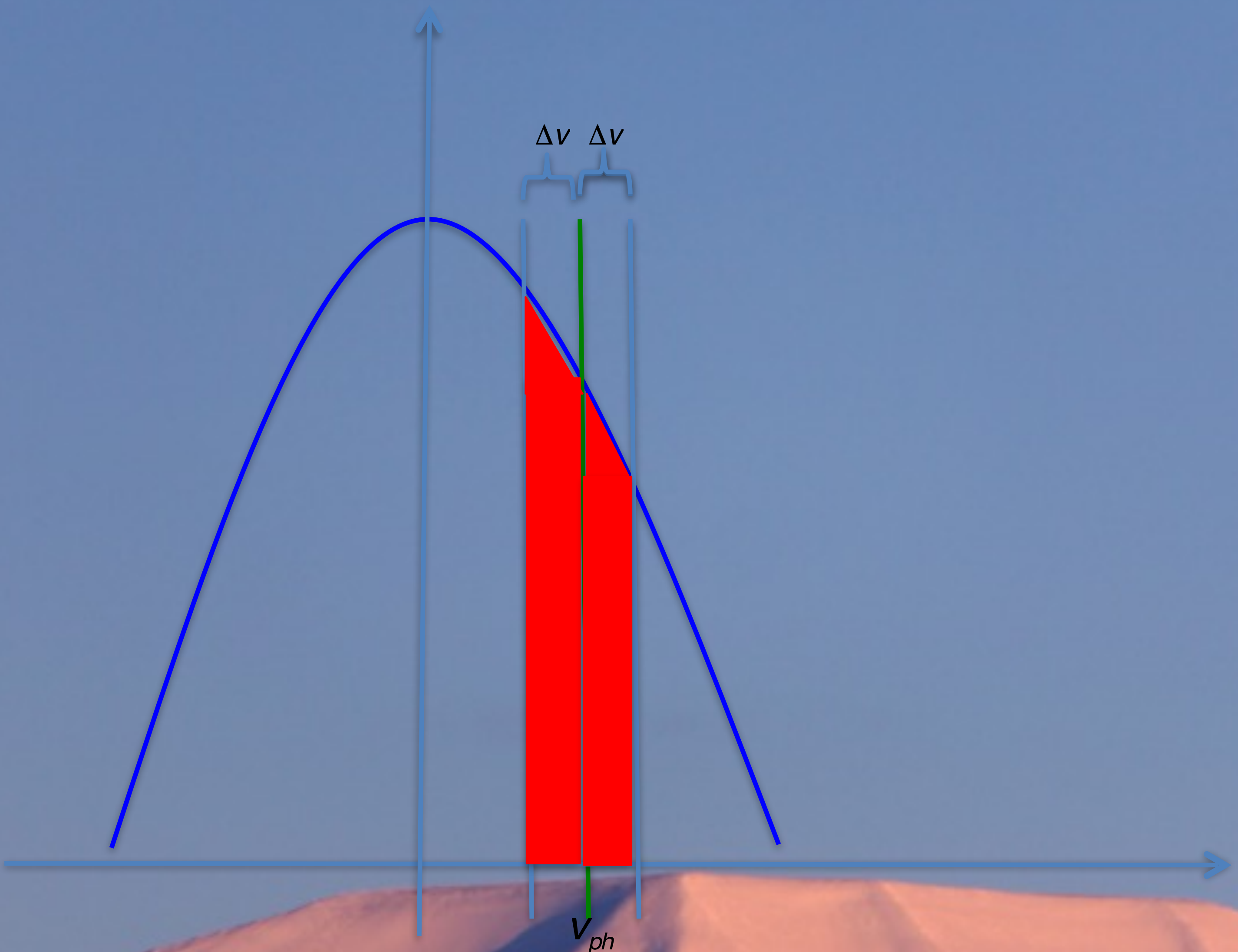

Δv Δv
v ph



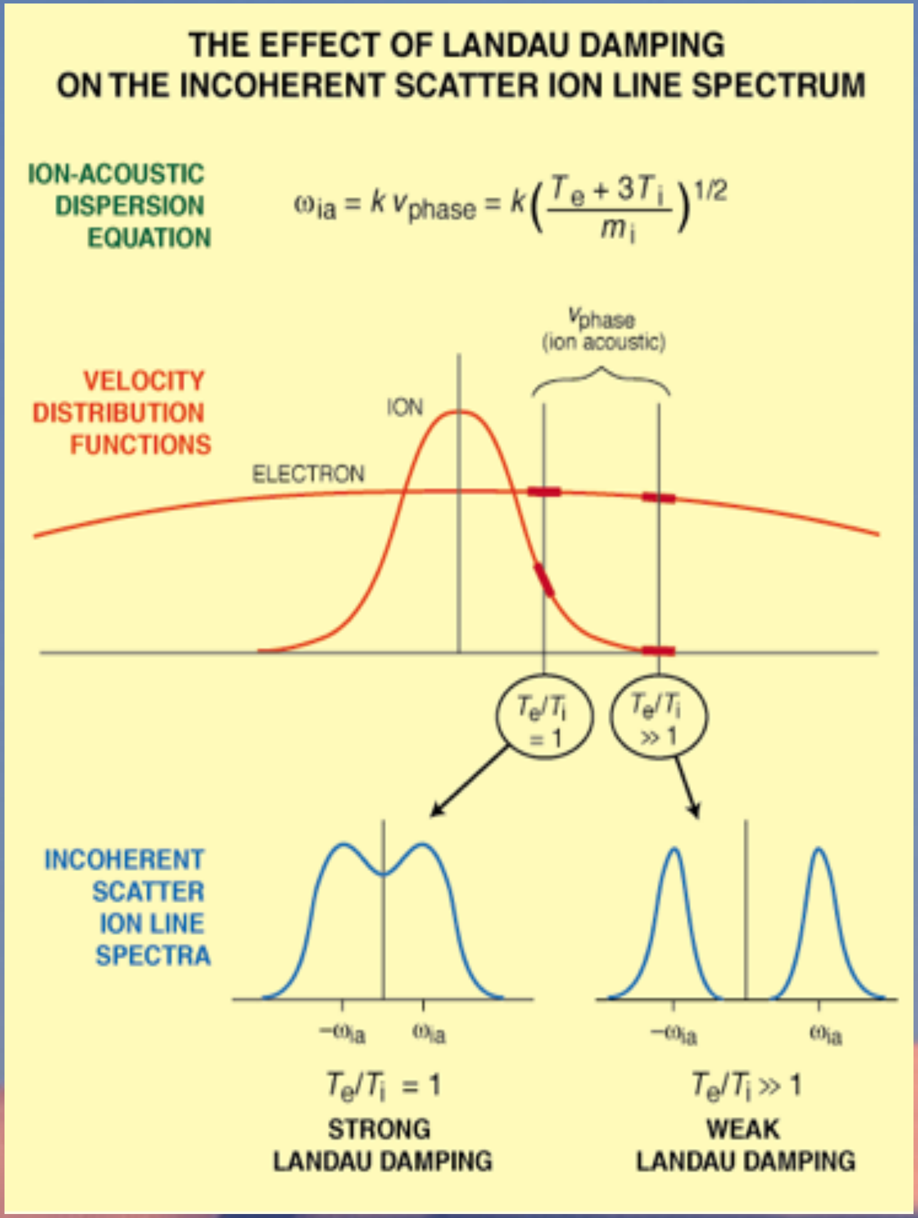
Δv Δv
v_{ph}



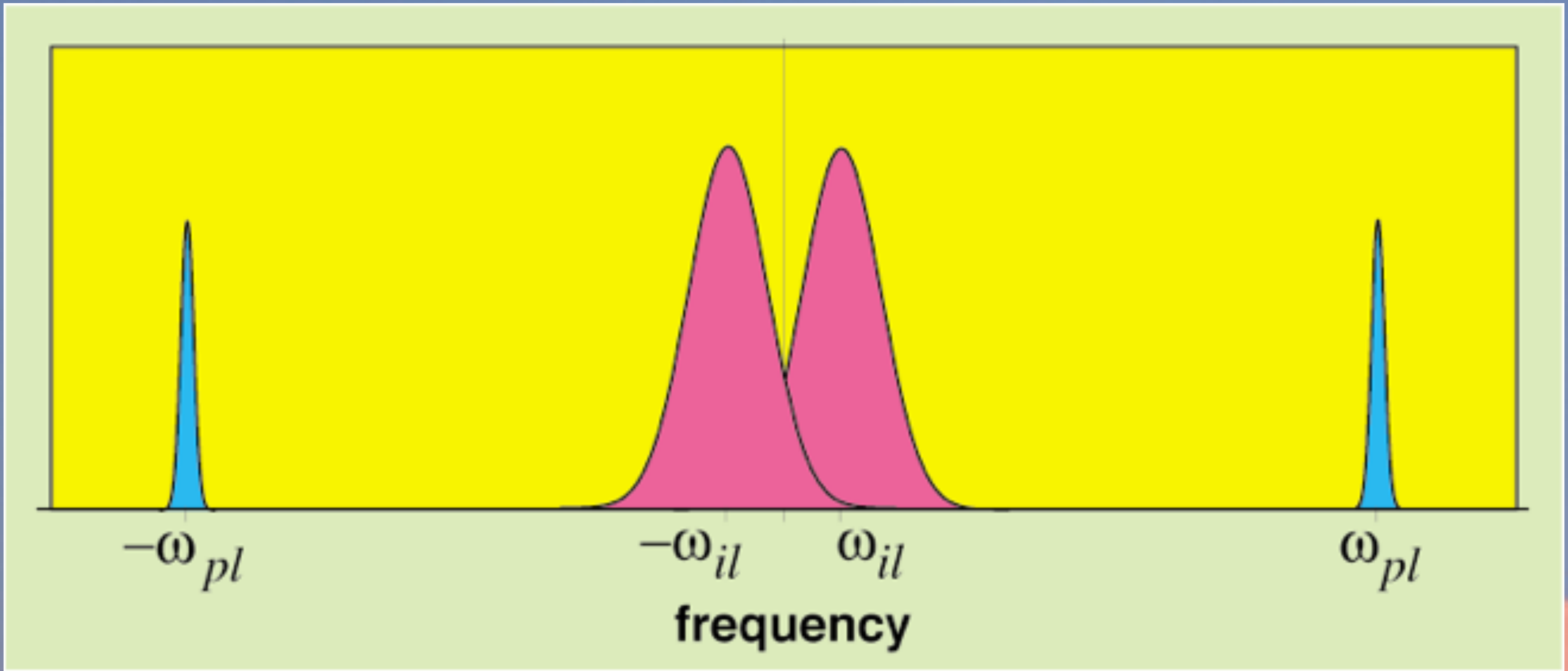
Δv Δv
v_{ph}



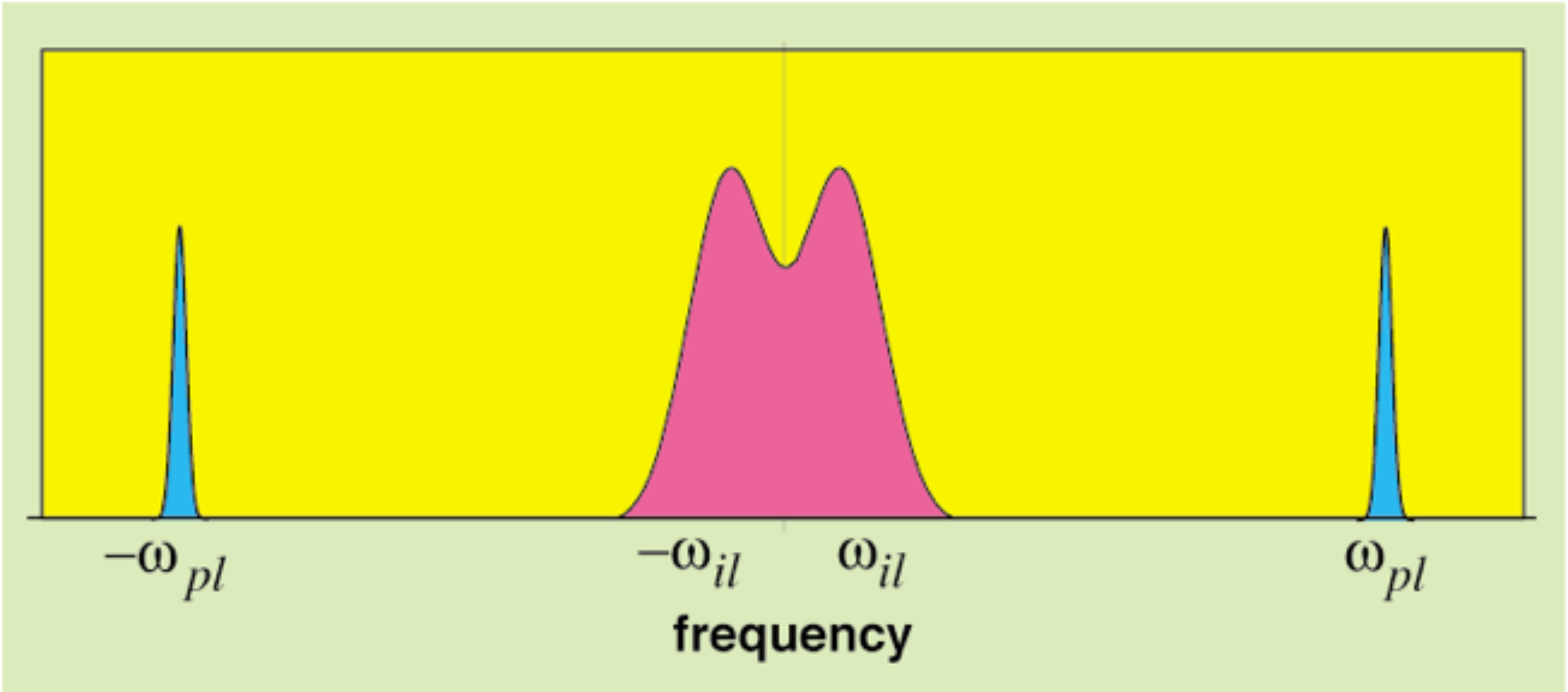
Δv Δv
v_{ph}



THE EFFECT OF LANDAU DAMPING
ON THE INCOHERENT SCATTER ION LINE SPECTRUM
ION-ACOUSTIC DISPERSION EQUATION
$\omega_{ia} = k\, v_{phase} = k\left(\frac{T_e + 3T_i}{m_i}\right)^{1/2}$
v_{phase} (ion acoustic)
VELOCITY DISTRIBUTION FUNCTIONS
ION
ELECTRON
$T_e/T_i = 1$
$T_e/T_i \gg 1$
INCOHERENT SCATTER ION LINE SPECTRA
$-\omega_{ia}$
ω_{ia}
$-\omega_{ia}$
ω_{ia}
$T_e/T_i = 1$
STRONG LANDAU DAMPING
$T_e/T_i \gg 1$
WEAK LANDAU DAMPING



$-\omega_{pl}$
$-\omega_{il}$
ω_{il}
ω_{pl}
frequency



$-\omega_{pl}$
$-\omega_{il}$
ω_{il}
ω_{pl}
frequency

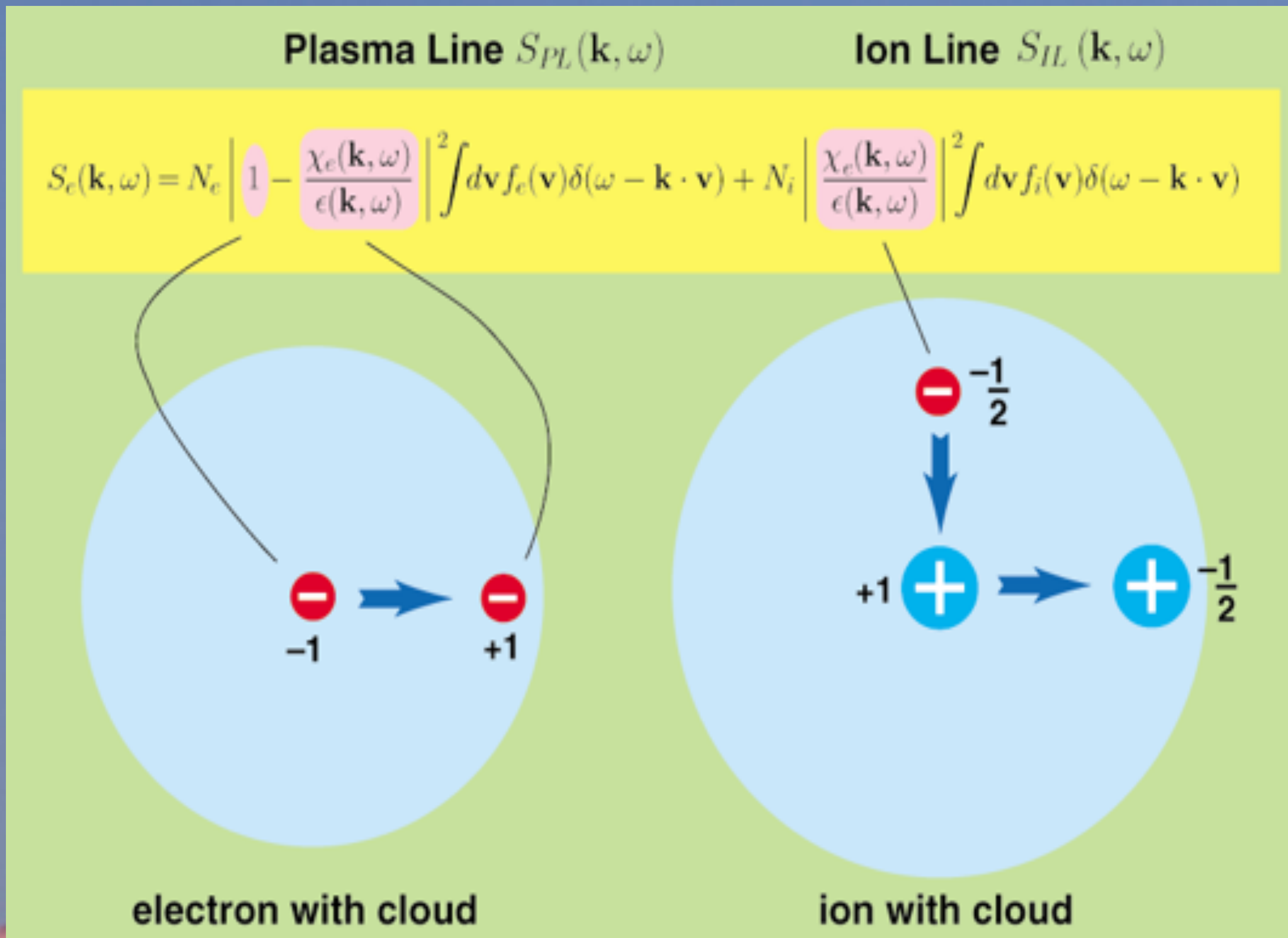
Incoherent Scattering Spectrum

$$S_e(\mathbf{k},\omega) = N_e \left|1 - \frac{\chi_e(\mathbf{k},\omega)}{\epsilon(\mathbf{k},\omega)}\right|^2 \int d\mathbf{v}\, f_e(\mathbf{v})\,\delta(\omega - \mathbf{k}\cdot\mathbf{v}) + N_i \left|\frac{\chi_e(\mathbf{k},\omega)}{\epsilon(\mathbf{k},\omega)}\right|^2 \int d\mathbf{v}\, f_i(\mathbf{v})\,\delta(\omega - \mathbf{k}\cdot\mathbf{v})$$

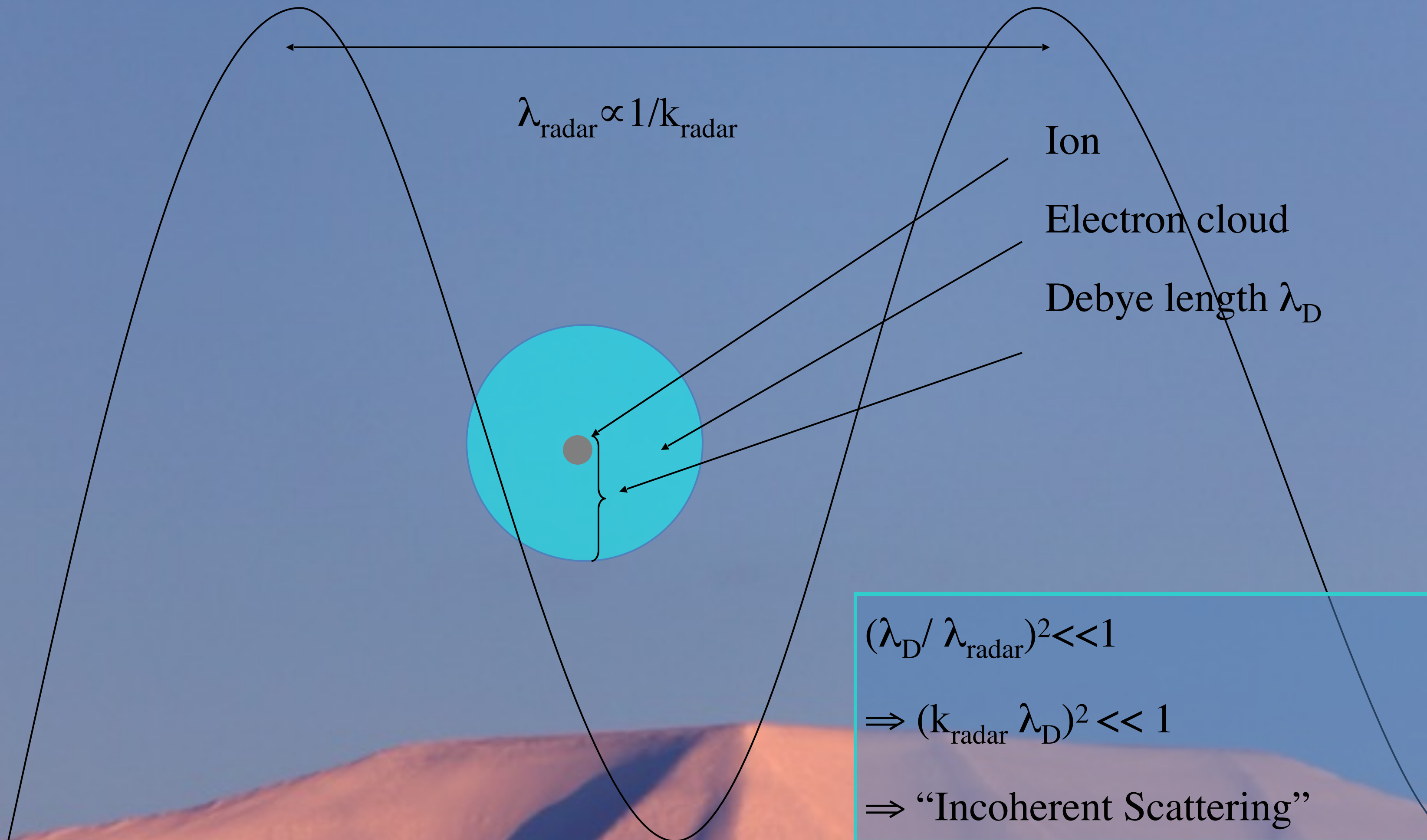
Plasma line

Ion line

electric susceptibility	$\chi_{e,i}(\mathbf{k},\omega)$
dielectric constant function	$\epsilon(\mathbf{k},\omega)$
velocity distribution function	$f_{e,i}(\mathbf{v})$


Plasma Line $S_{PL}(\mathbf{k},\omega)$
Ion Line $S_{IL}(\mathbf{k},\omega)$
$S_e(\mathbf{k},\omega) = N_e \left| 1 - \frac{\chi_e(\mathbf{k},\omega)}{\epsilon(\mathbf{k},\omega)} \right|^2 \int d\mathbf{v} f_e(\mathbf{v}) \delta(\omega - \mathbf{k}\cdot\mathbf{v}) + N_i \left| \frac{\chi_e(\mathbf{k},\omega)}{\epsilon(\mathbf{k},\omega)} \right|^2 \int d\mathbf{v} f_i(\mathbf{v}) \delta(\omega - \mathbf{k}\cdot\mathbf{v})$
−1
+1
$-\frac{1}{2}$
+1
$-\frac{1}{2}$
electron with cloud
ion with cloud

Debye length dependence



$(\lambda_D / \lambda_{radar})^2 << 1$

$\Rightarrow (k_{radar}\, \lambda_D)^2 << 1$

$\Rightarrow$ “Incoherent Scattering”

Plasma Line $S_{PL}(\mathbf{k},\omega)$ **Ion Line** $S_{IL}(\mathbf{k},\omega)$

$$S_e(\mathbf{k},\omega) = N_e \left|1 - \frac{\chi_e(\mathbf{k},\omega)}{\epsilon(\mathbf{k},\omega)}\right|^2 \int d\mathbf{v}\, f_e(\mathbf{v})\delta(\omega - \mathbf{k}\cdot\mathbf{v}) + N_i \left|\frac{\chi_e(\mathbf{k},\omega)}{\epsilon(\mathbf{k},\omega)}\right|^2 \int d\mathbf{v}\, f_i(\mathbf{v})\delta(\omega - \mathbf{k}\cdot\mathbf{v})$$

Plasma Line $S_{PL}(\mathbf{k},\omega)$ **Ion Line** $S_{IL}(\mathbf{k},\omega)$

$$S_e(\mathbf{k},\omega) = N_e \left|1 - \frac{\chi_e(\mathbf{k},\omega)}{\epsilon(\mathbf{k},\omega)}\right|^2 \int d\mathbf{v}\, f_e(\mathbf{v})\delta(\omega - \mathbf{k}\cdot\mathbf{v}) + N_i \left|\frac{\chi_e(\mathbf{k},\omega)}{\epsilon(\mathbf{k},\omega)}\right|^2 \int d\mathbf{v}\, f_i(\mathbf{v})\delta(\omega - \mathbf{k}\cdot\mathbf{v})$$

$$\epsilon(\mathbf{k},\omega) = \mathbf{0}$$

$$\omega_{pl}(k) \approx \omega_{pe}(1 + 3\lambda_D^2 k^2)$$

$$\omega_{ia}(k) \approx k\sqrt{\frac{T_e + 3T_i}{m_i}}$$

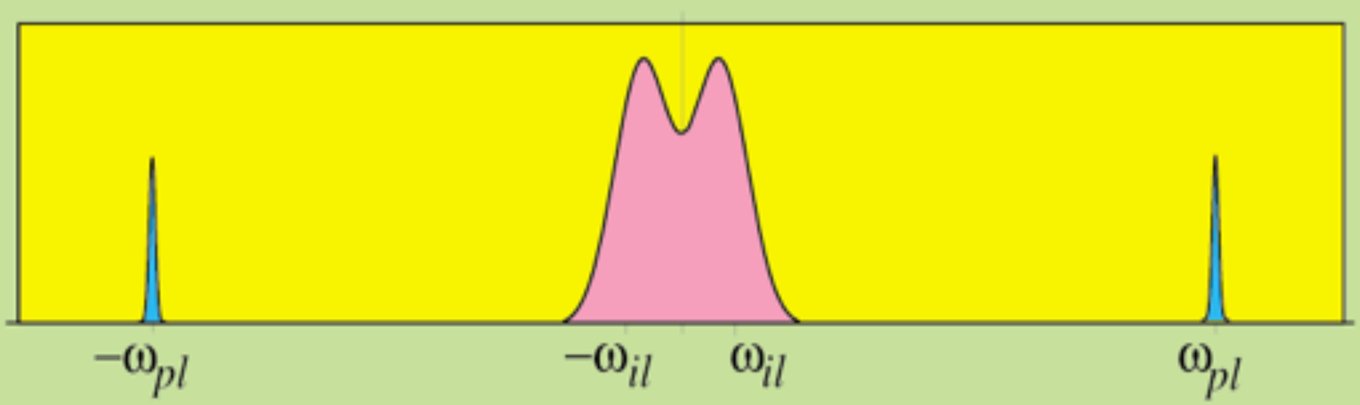
Plasma Line $S_{PL}(\mathbf{k},\omega)$ **Ion Line** $S_{IL}(\mathbf{k},\omega)$

$$S_e(\mathbf{k},\omega) = N_e\left|1-\frac{\chi_e(\mathbf{k},\omega)}{\epsilon(\mathbf{k},\omega)}\right|^2\int d\mathbf{v}\, f_e(\mathbf{v})\delta(\omega-\mathbf{k}\cdot\mathbf{v}) + N_i\left|\frac{\chi_e(\mathbf{k},\omega)}{\epsilon(\mathbf{k},\omega)}\right|^2\int d\mathbf{v}\, f_i(\mathbf{v})\delta(\omega-\mathbf{k}\cdot\mathbf{v})$$

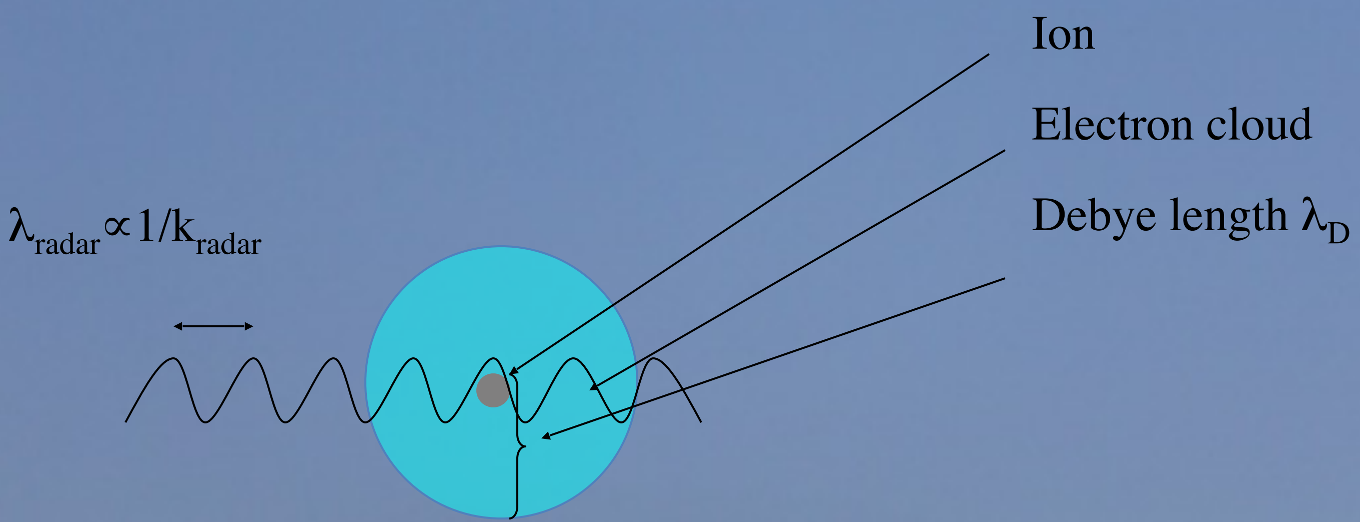
$$\epsilon(\mathbf{k},\omega) = \mathbf{0}$$

$$\omega_{pl}(k) \approx \omega_{pe}(1+3\lambda_D^2k^2)$$

$$\omega_{ia}(k) \approx k\sqrt{\frac{T_e+3T_i}{m_i}}$$



Debye length dependence



$(\lambda_D / \lambda_{radar})^2 > 1$

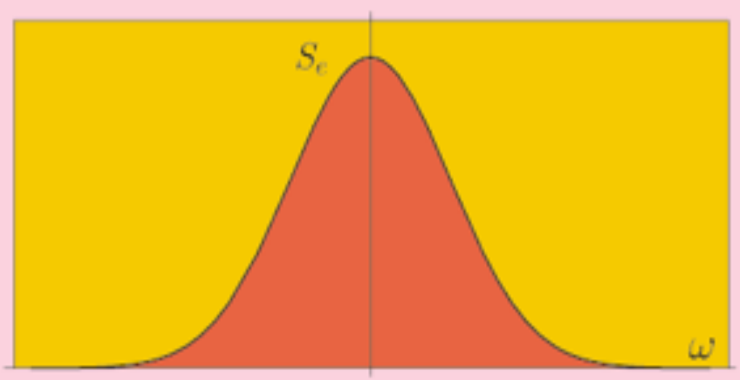
$\Rightarrow (k_{radar} \lambda_D)^2 > 1$

$\Rightarrow$ No collective interactions

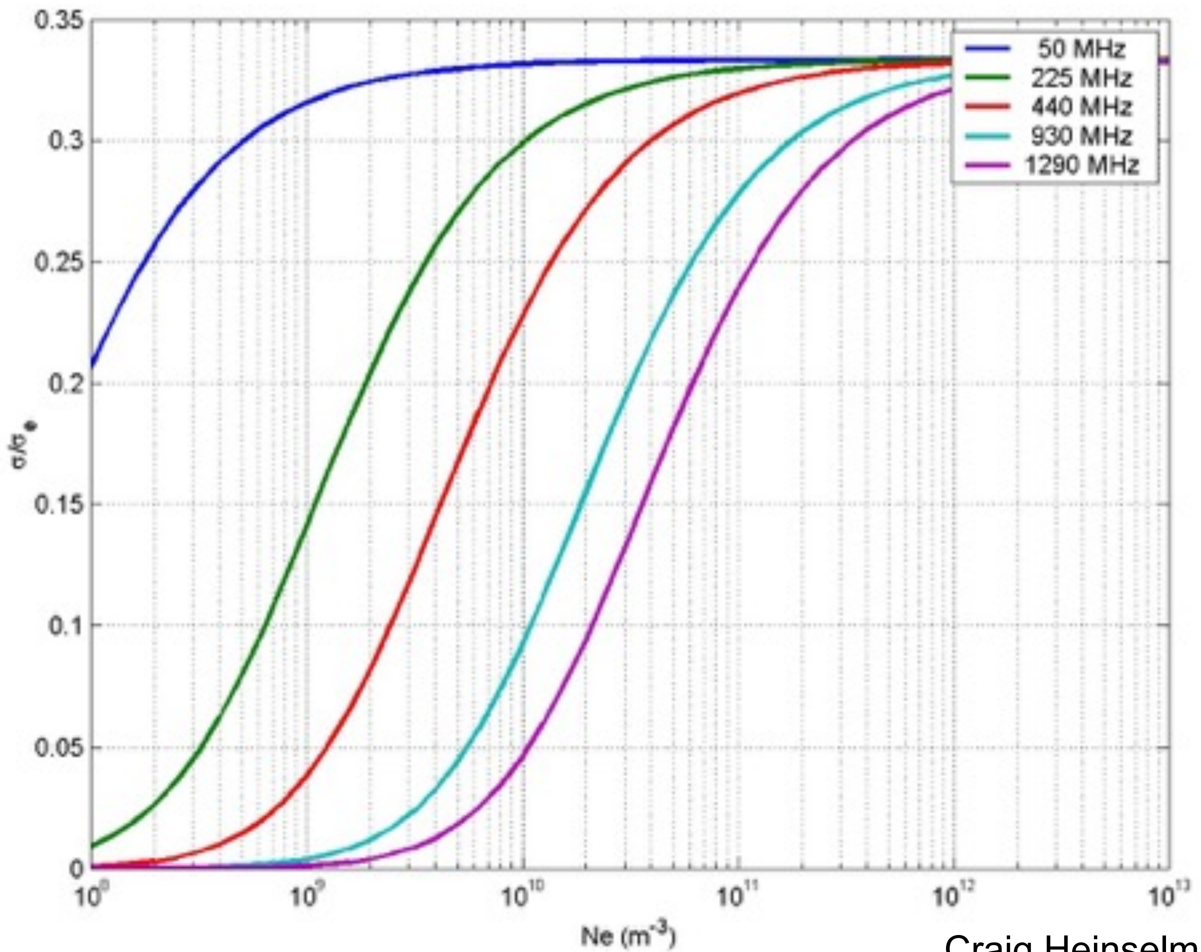
no collective interactions

$$S_e(\mathbf{k},\omega) = N_e\left|1 - \frac{\chi_e(\mathbf{k},\omega)}{\epsilon(\mathbf{k},\omega)}\right|^2 \int d\mathbf{v}\, f_e(\mathbf{v})\delta(\omega - \mathbf{k}\cdot\mathbf{v}) + N_i\left|\frac{\chi_e(\mathbf{k},\omega)}{\epsilon(\mathbf{k},\omega)}\right|^2 \int d\mathbf{v}\, f_i(\mathbf{v})\delta(\omega - \mathbf{k}\cdot\mathbf{v})$$

$$S_e(\mathbf{k},\omega) = N_e \int d\mathbf{v}\, f_e(\mathbf{v})\delta(\omega - \mathbf{k}\cdot\mathbf{v})$$



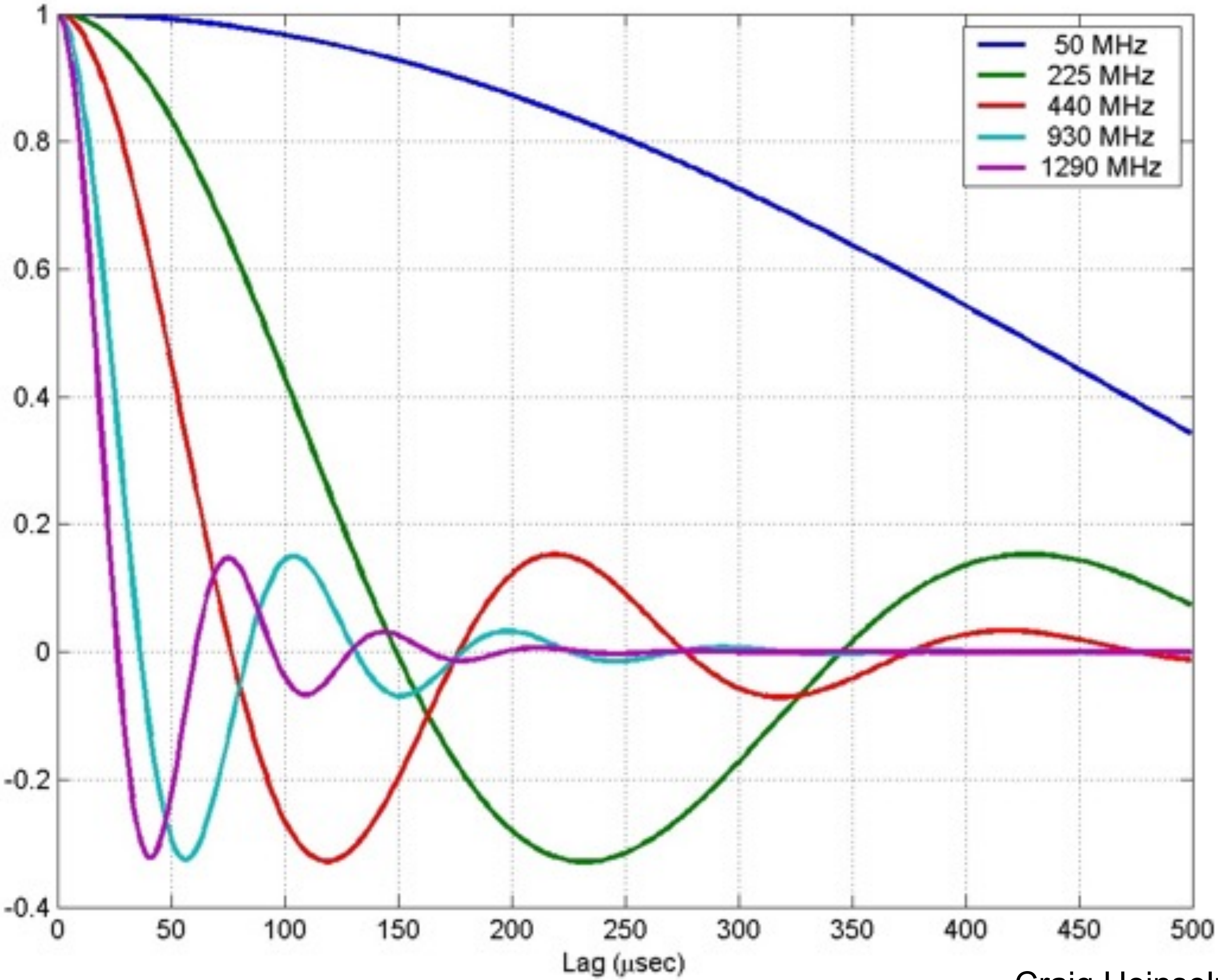
Debye Length Dependencies



Craig Heinselman

Parameters
Ti: 1000 K
Te: 2000 K

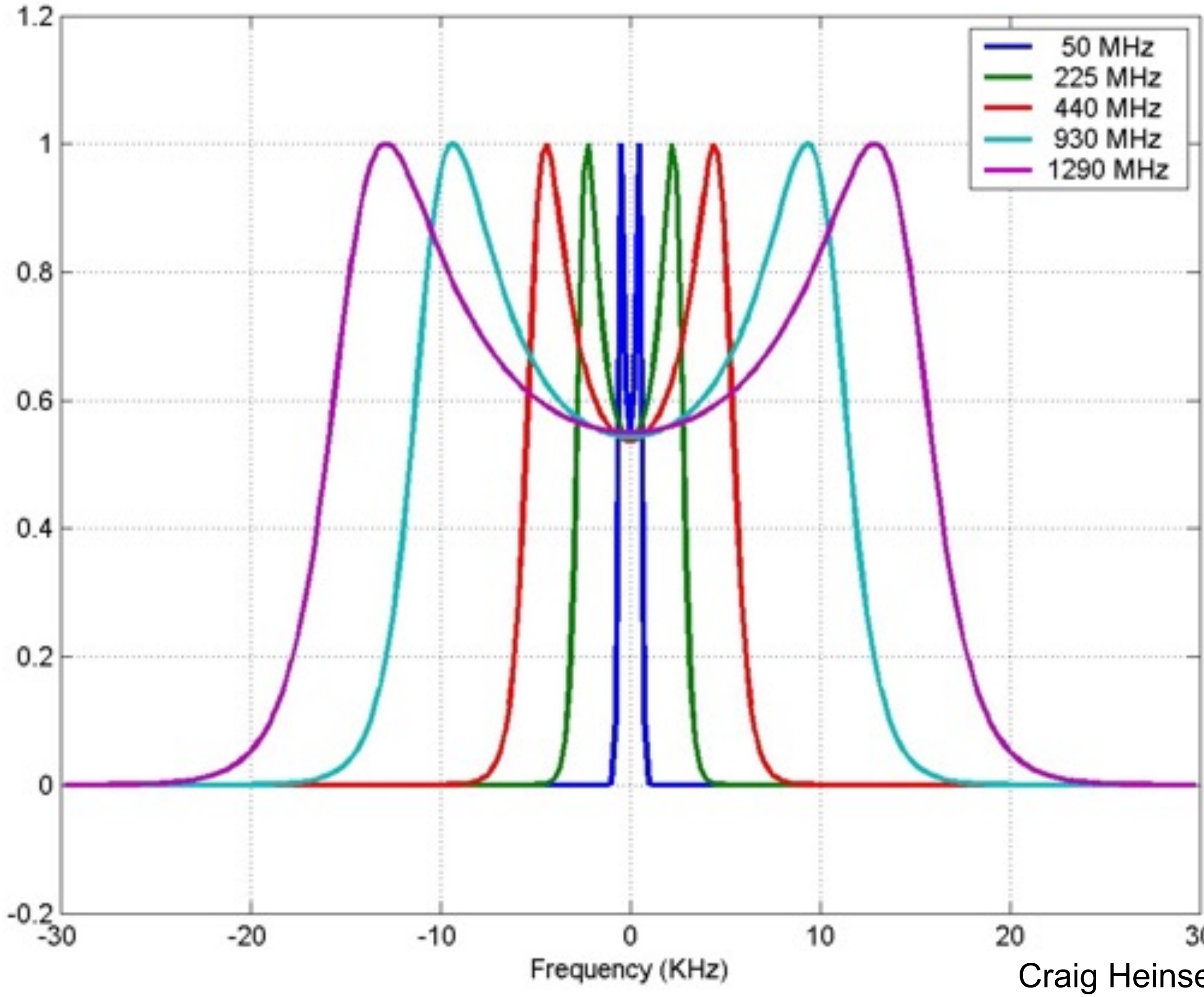
Radar Frequency Dependencies



Parameters
Ne: 10^{12} m^{-3}
Ti: 1000 K
Te: 2000 K
Comp: 100% O^+
ν_{in}: 10^{-6} KHz

Craig Heinselman

Radar Frequency Dependencies

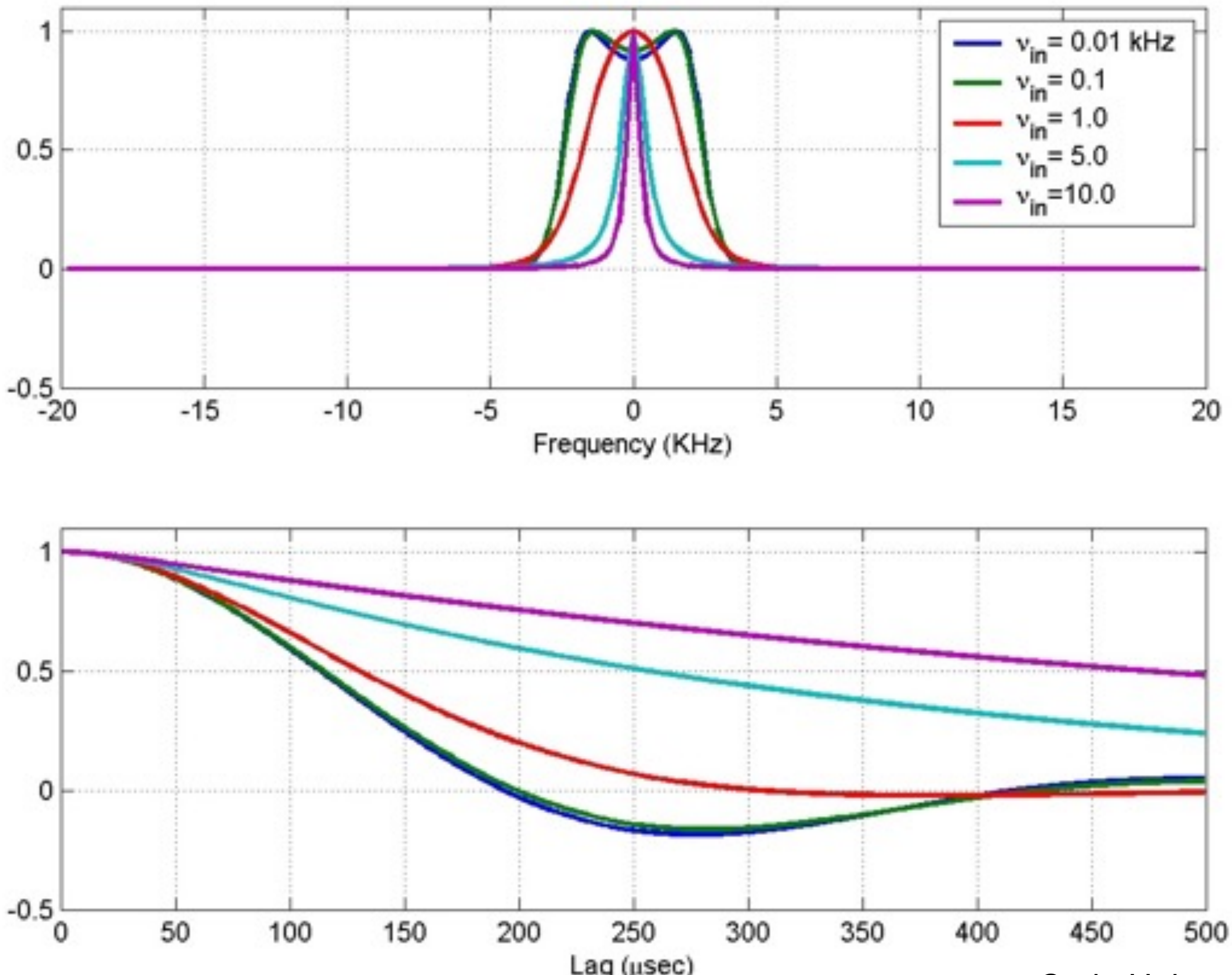


Parameters
Ne: 10^{12} m^{-3}
Ti: 1000 K
Te: 2000 K
Comp: 100% O^+
ν_{in}: 10^{-6} KHz

Craig Heinselman

With the frequency of the radar chosen (which is a one time thing!), how does the spectra depend on geophysical parameters?

Ion-Neutral Collision Frequency



Parameters
Freq: 449 MHz
Ne: 10^{12} m^{-3}
Ti: 500 K
Te: 500 K
Comp: 100% NO^{+}

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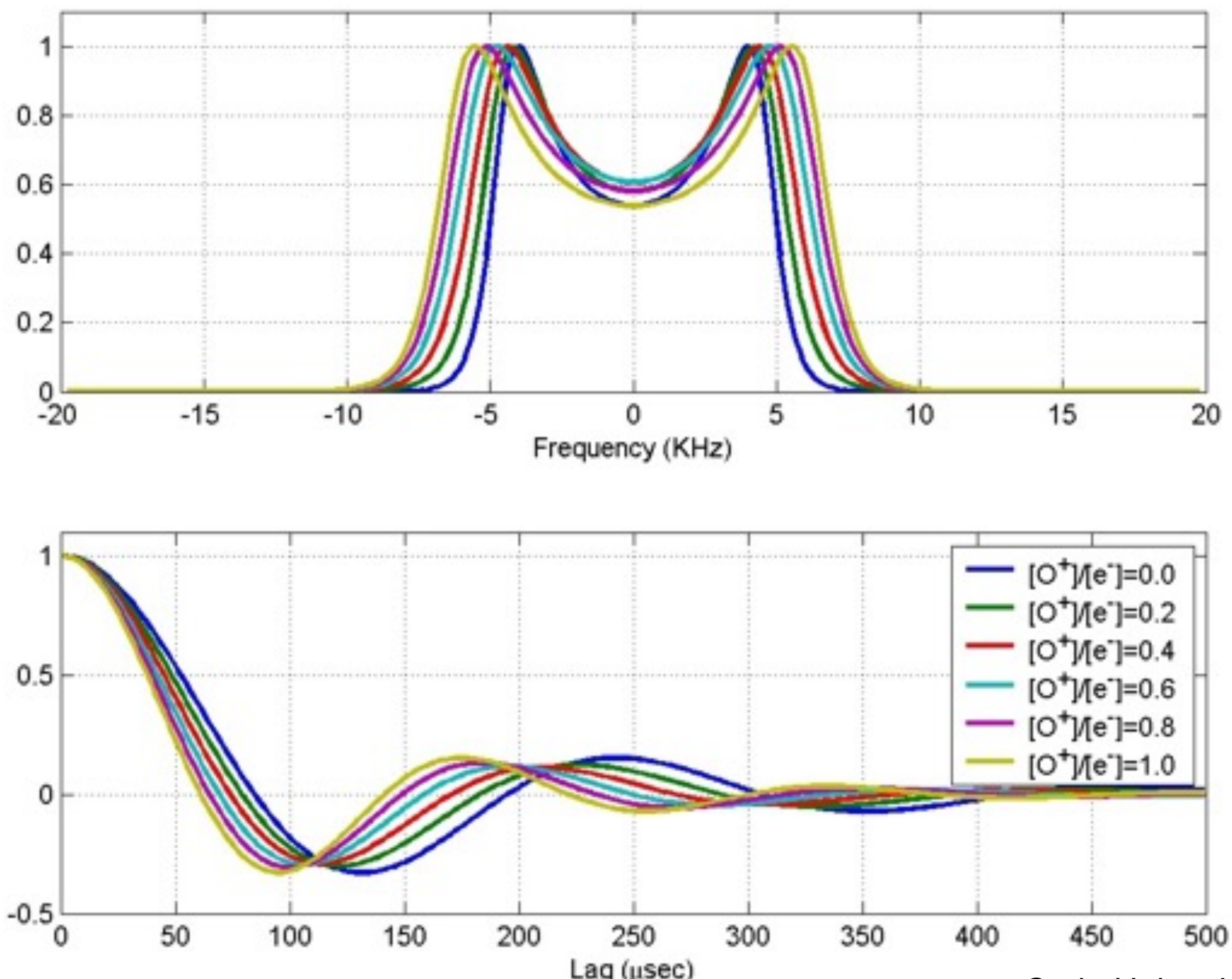
Ion Composition (O^+ vs. NO^+)



Parameters

Freq: 449 MHz

Ne: 10^{12} m^{-3}

Ti: 1500 K

Te: 3000 K

ν_{in}: 10^{-6} KHz

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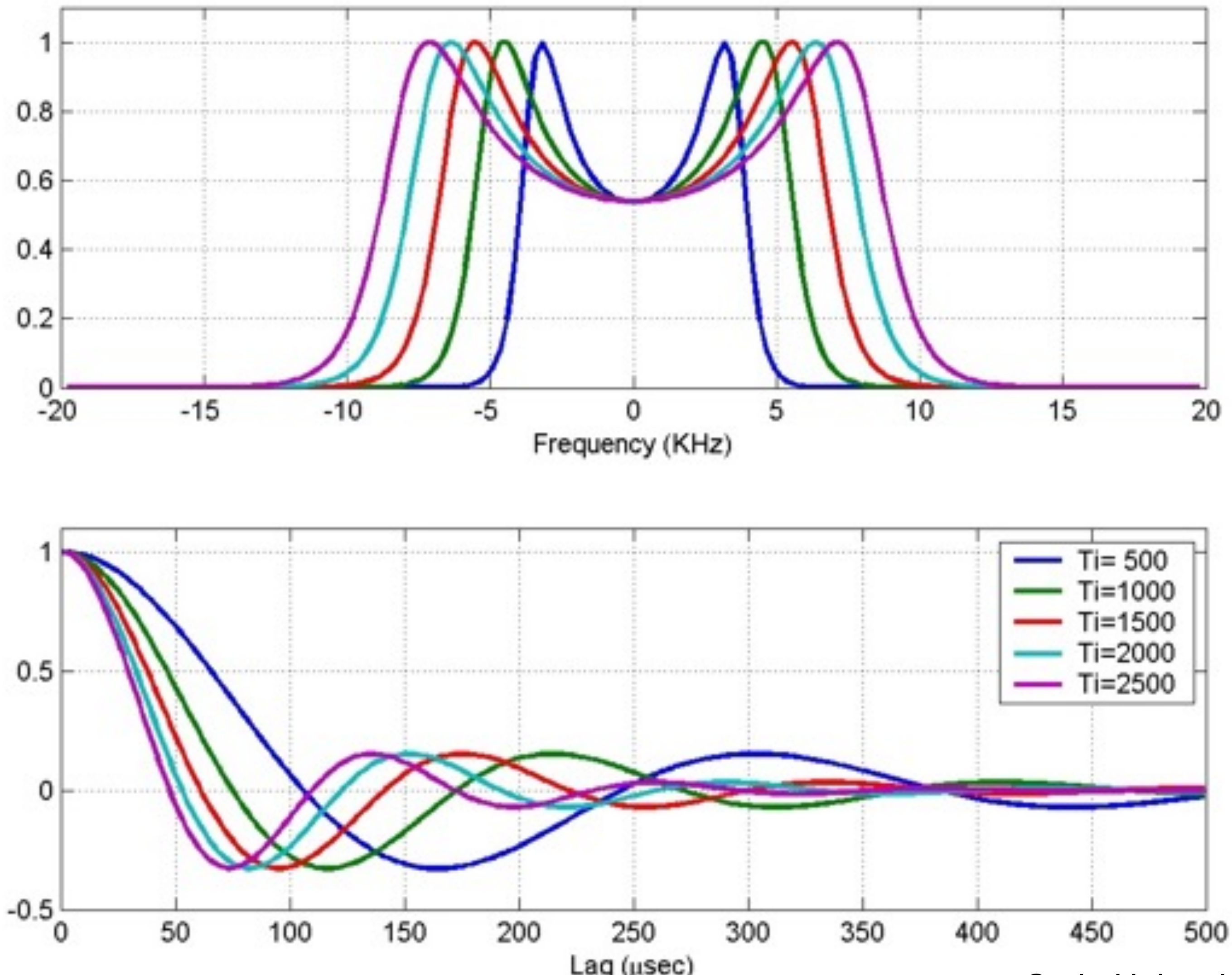
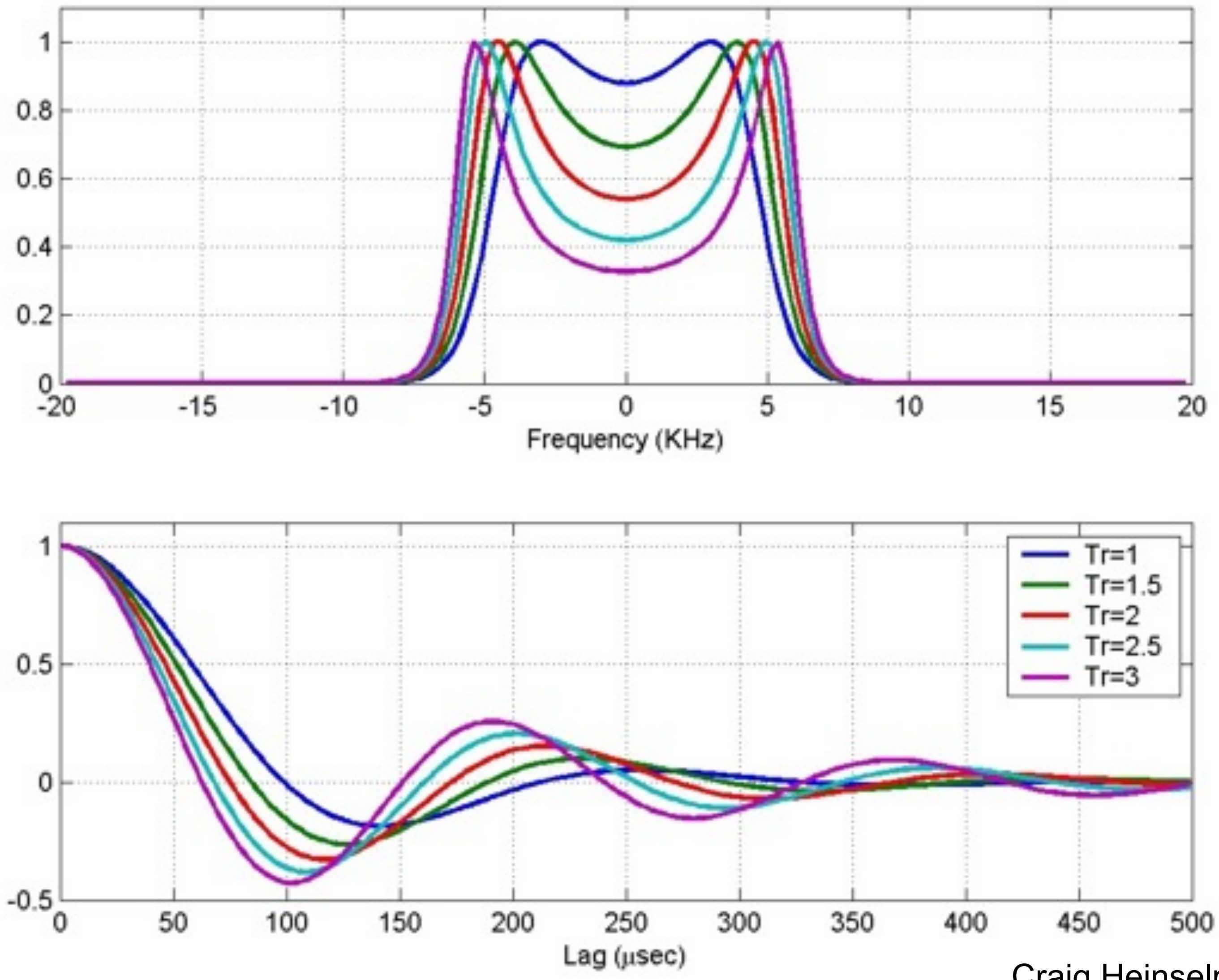
Ion Temperature



Parameters
Freq: 449 MHz
Ne: 10^{12} m^{-3}
Te: 2*Ti
Comp: 100% O^{+}
ν_{in}: 10^{-6} KHz

Craig Heinselman

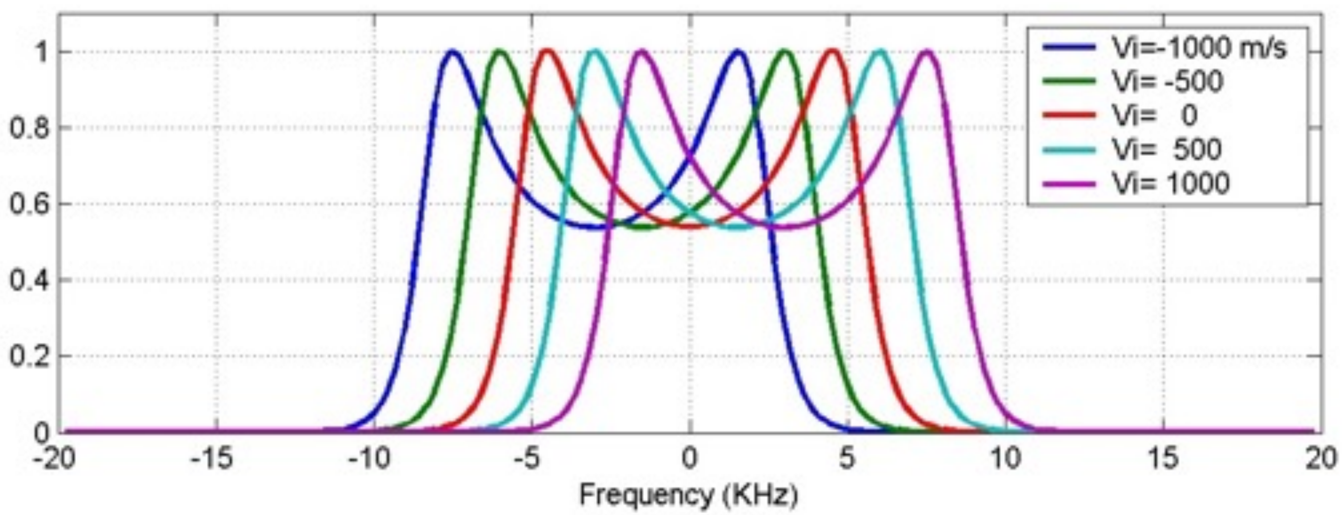
Electron/Ion Temperature Ratio

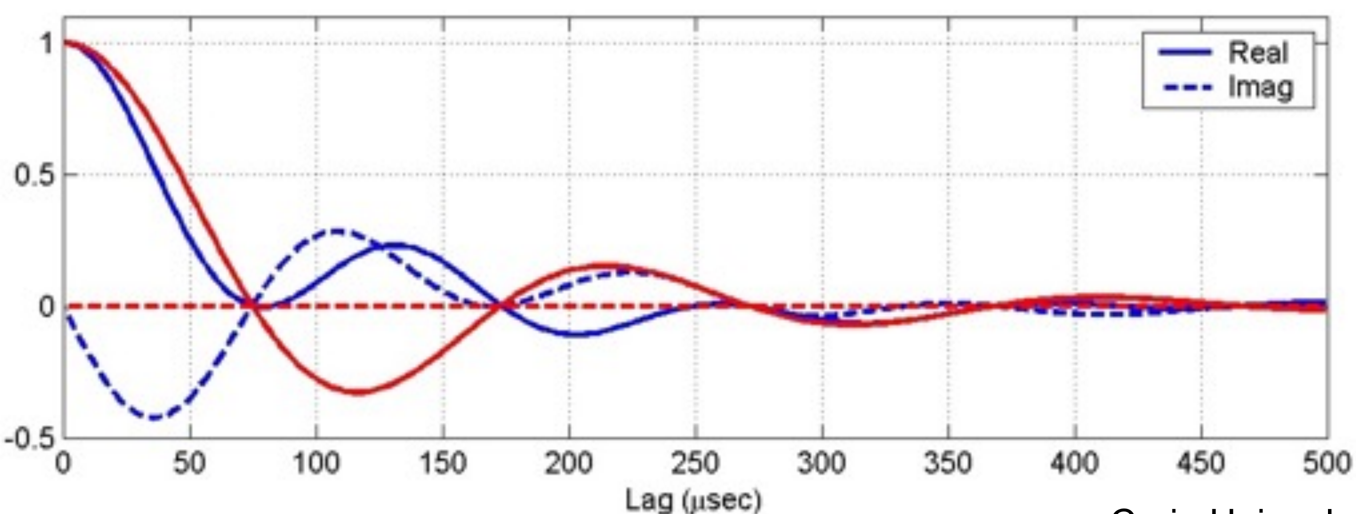


Parameters
Freq: 449 MHz
Ne: 10^{12} m^{-3}
Ti: 1000 K
Comp: 100% O^{+}
ν_{in}: 10^{-6} KHz

Craig Heinselman

Ion Velocity

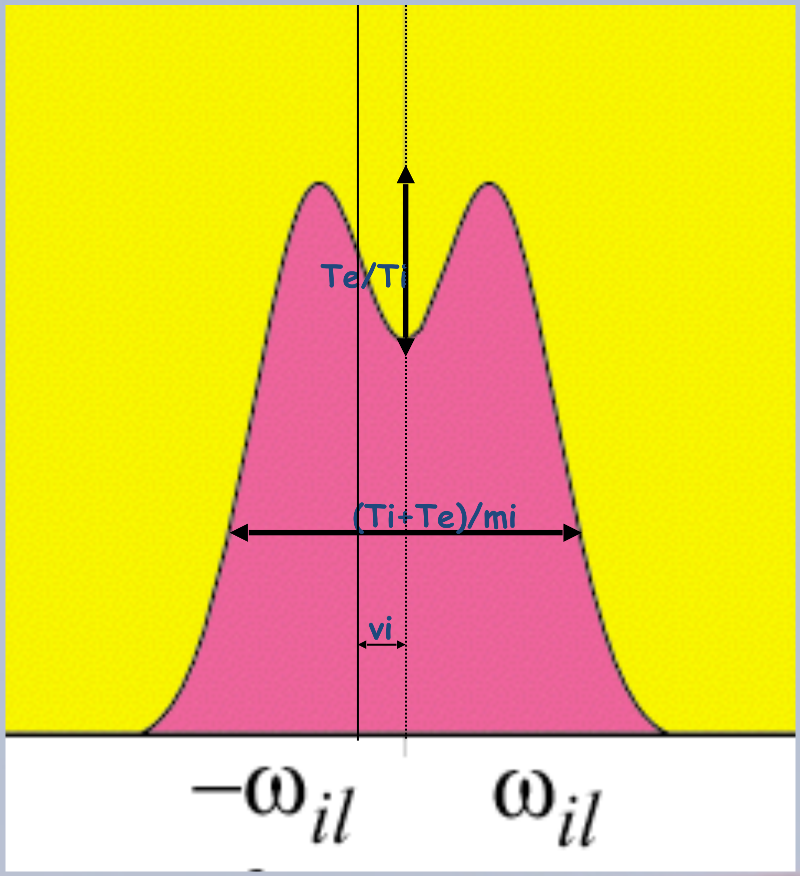




Parameters
Freq: 449 MHz
Ne: 10^{12} m^{-3}
Ti: 1000 K
Te: 2000 K
Comp: 100% O^+
ν_{in}: 10^{-6} KHz

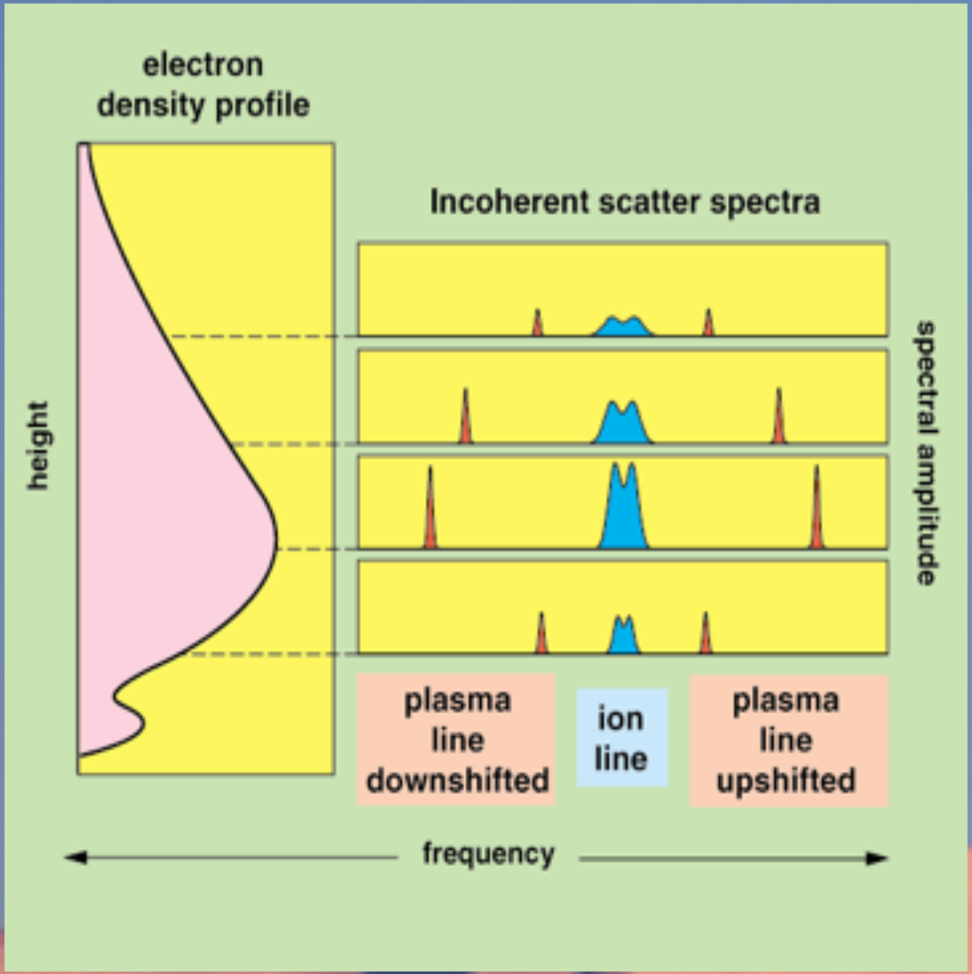
Craig Heinselman

…or to sum up…

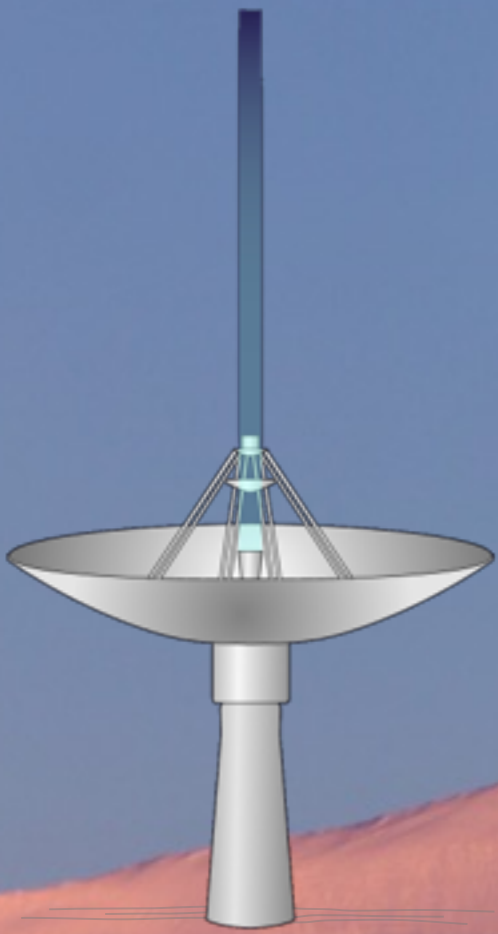


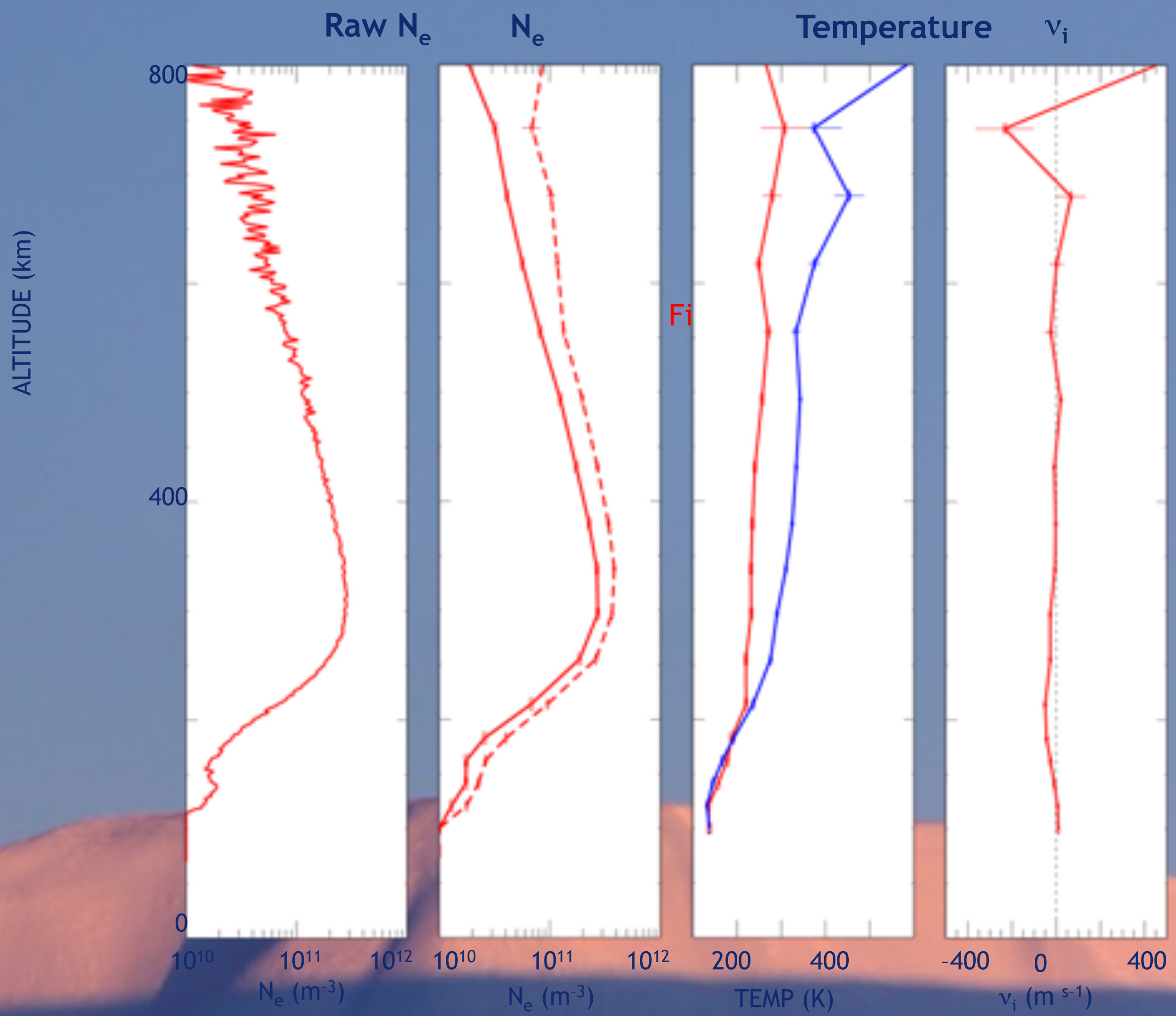
- Ion (and electron) temperature (Ti and Te) to ion mass (mi) ratio from the width of the spectra
- Electron to ion temperature ratio (Te/Ti) from "peak_to_valley" ratio
- Electron (= ion) density from total area (corrected for temperatures)
- Ion velocity (vi) from the Doppler shift

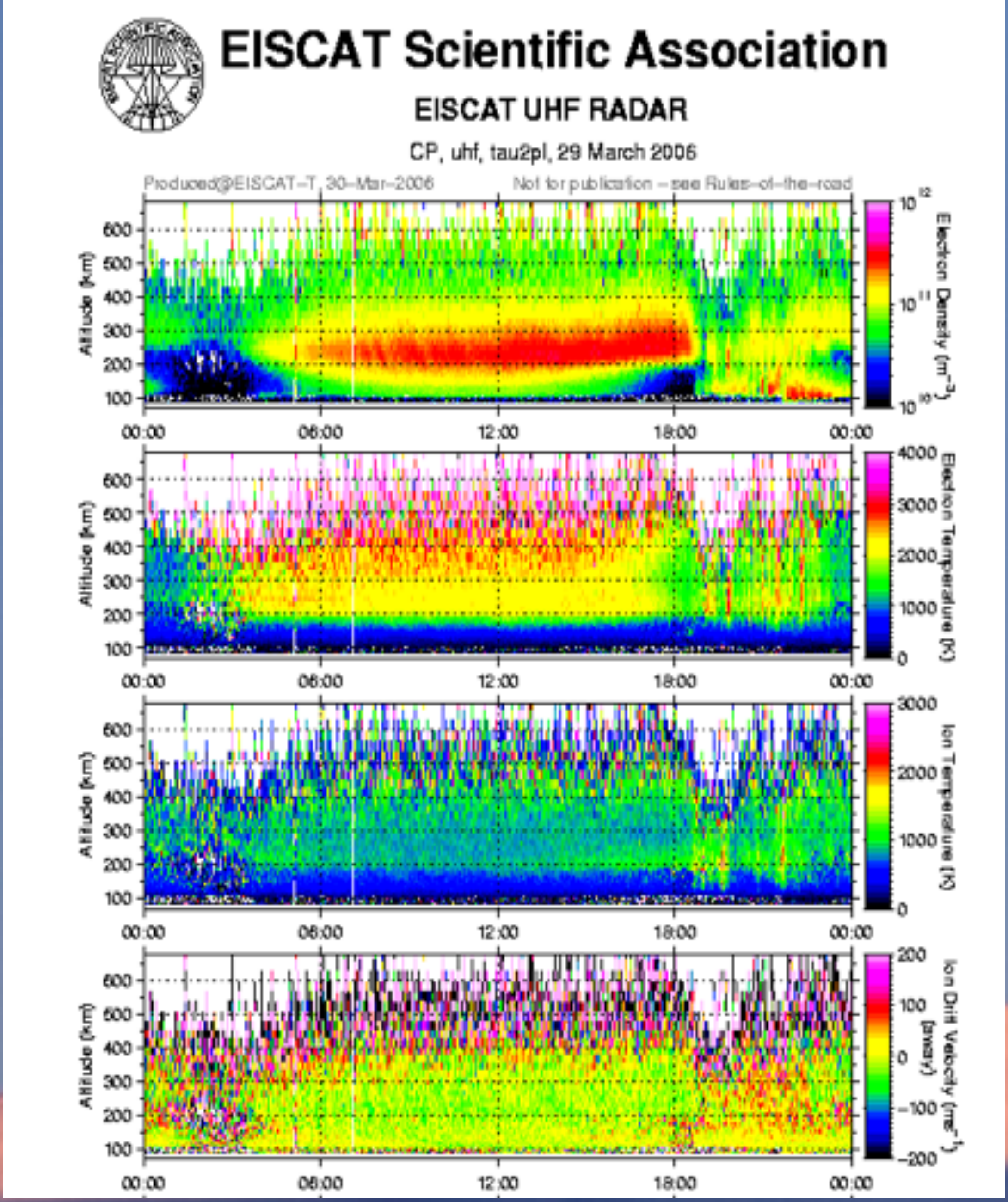
Spectral space as a function of altitude



Plasma Parameter Profile

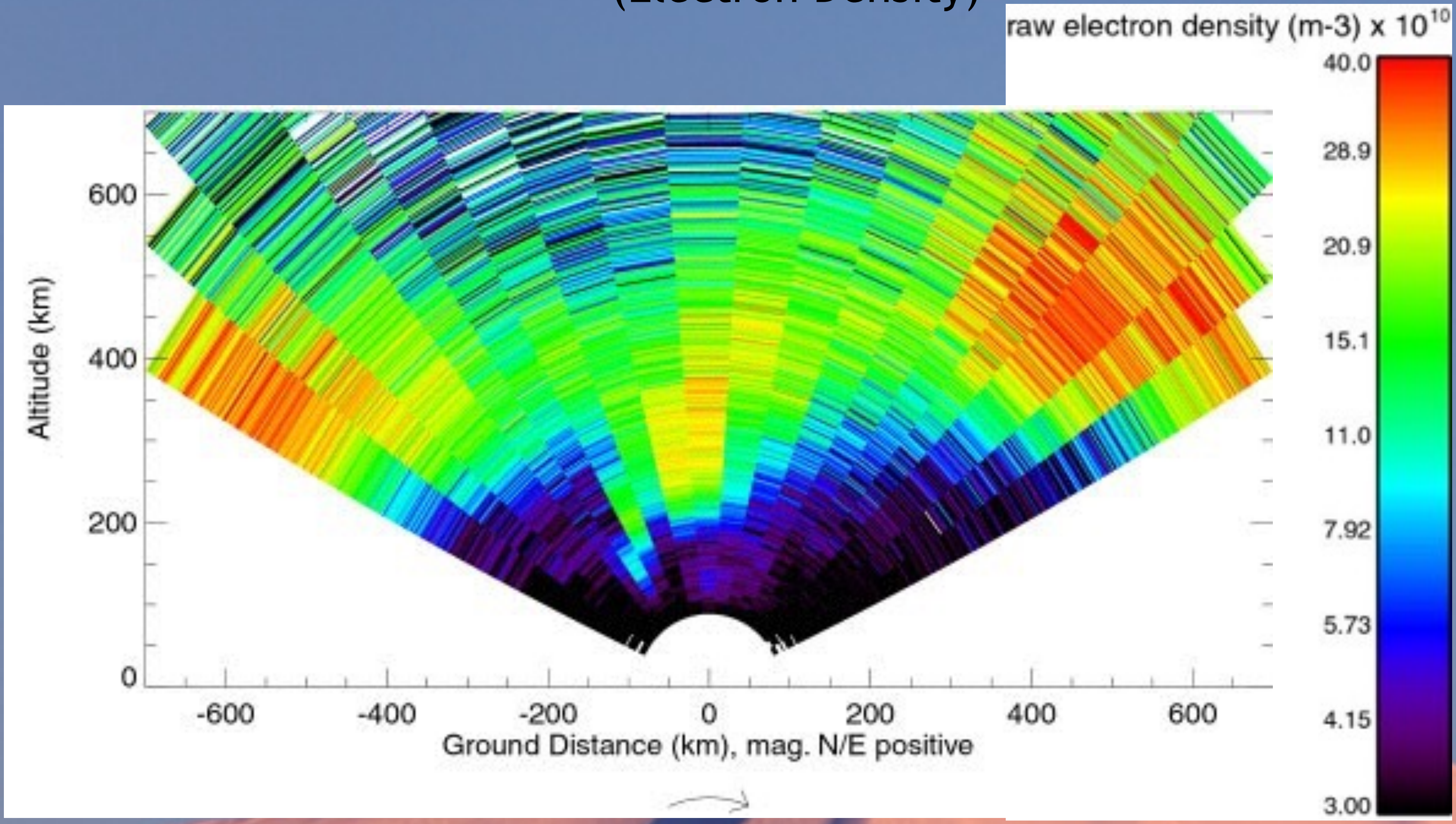




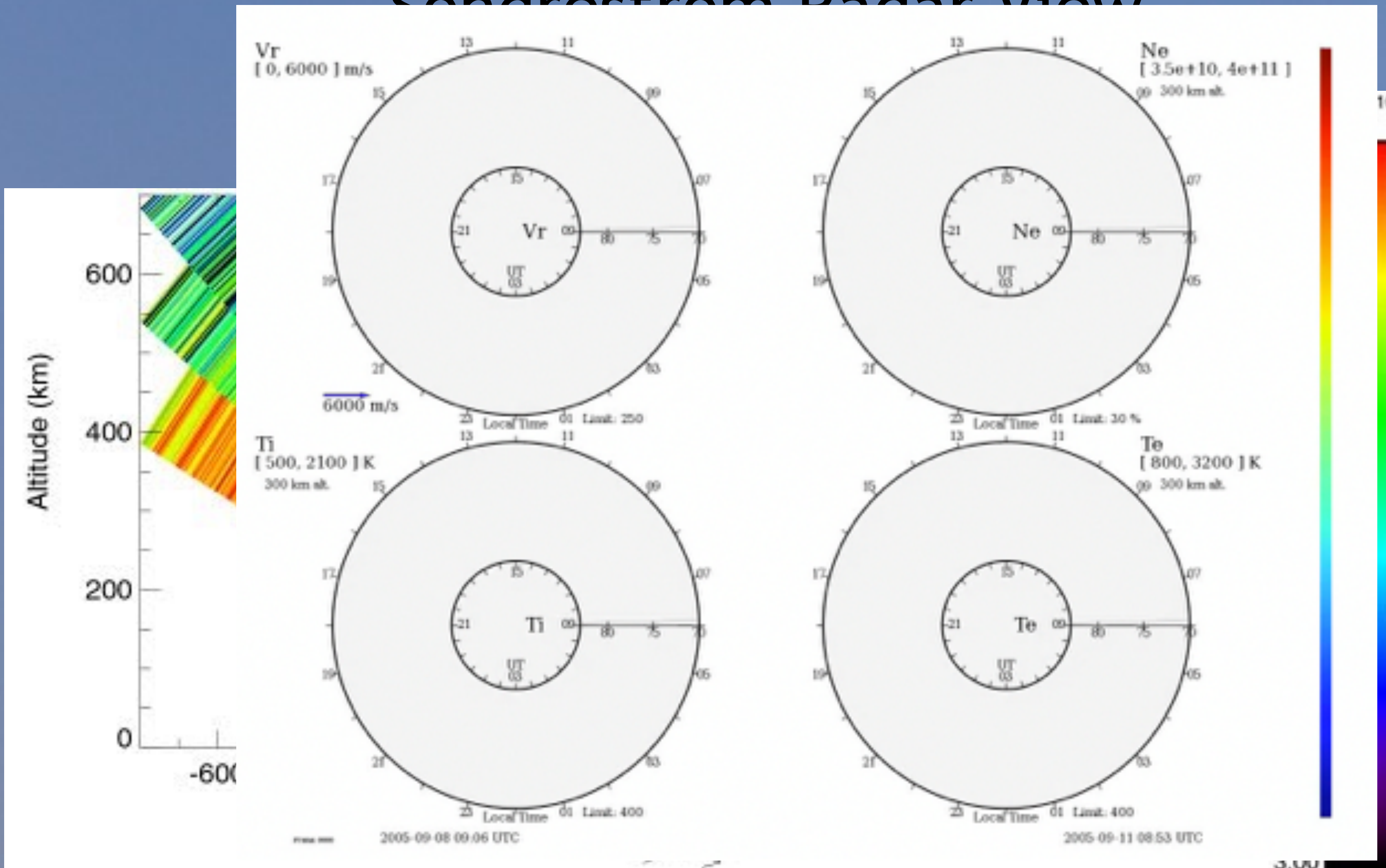

EISCAT Scientific Association
EISCAT UHF RADAR
CP, uhf, tau2pl, 29 March 2006
Produced@EISCAT-T, 30-Mar-2006
Not for publication – see Rules-of-the-road
Altitude (km)
Electron Density (m⁻³)
Electron Temperature (K)
Ion Temperature (K)
Ion Drift Velocity (ms⁻¹) [away]
00:00
06:00
12:00
18:00

Sondrestrom Radar View

(Electron Density)



Sondrestrom Radar View



And this is the level data we will work on in the MADRIGAL session…