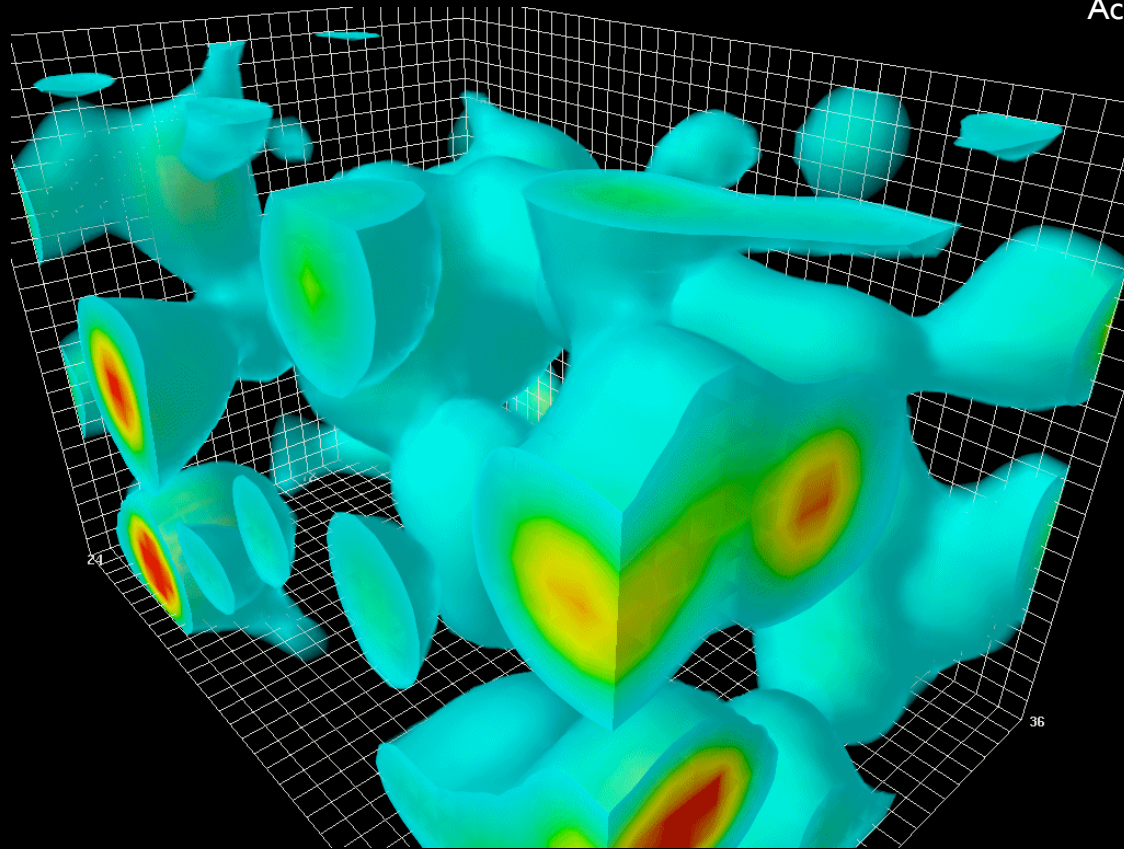


8.712  
Introduction to Heavy-Ion Physics  
Part II

3/18/2008

Gunther Roland

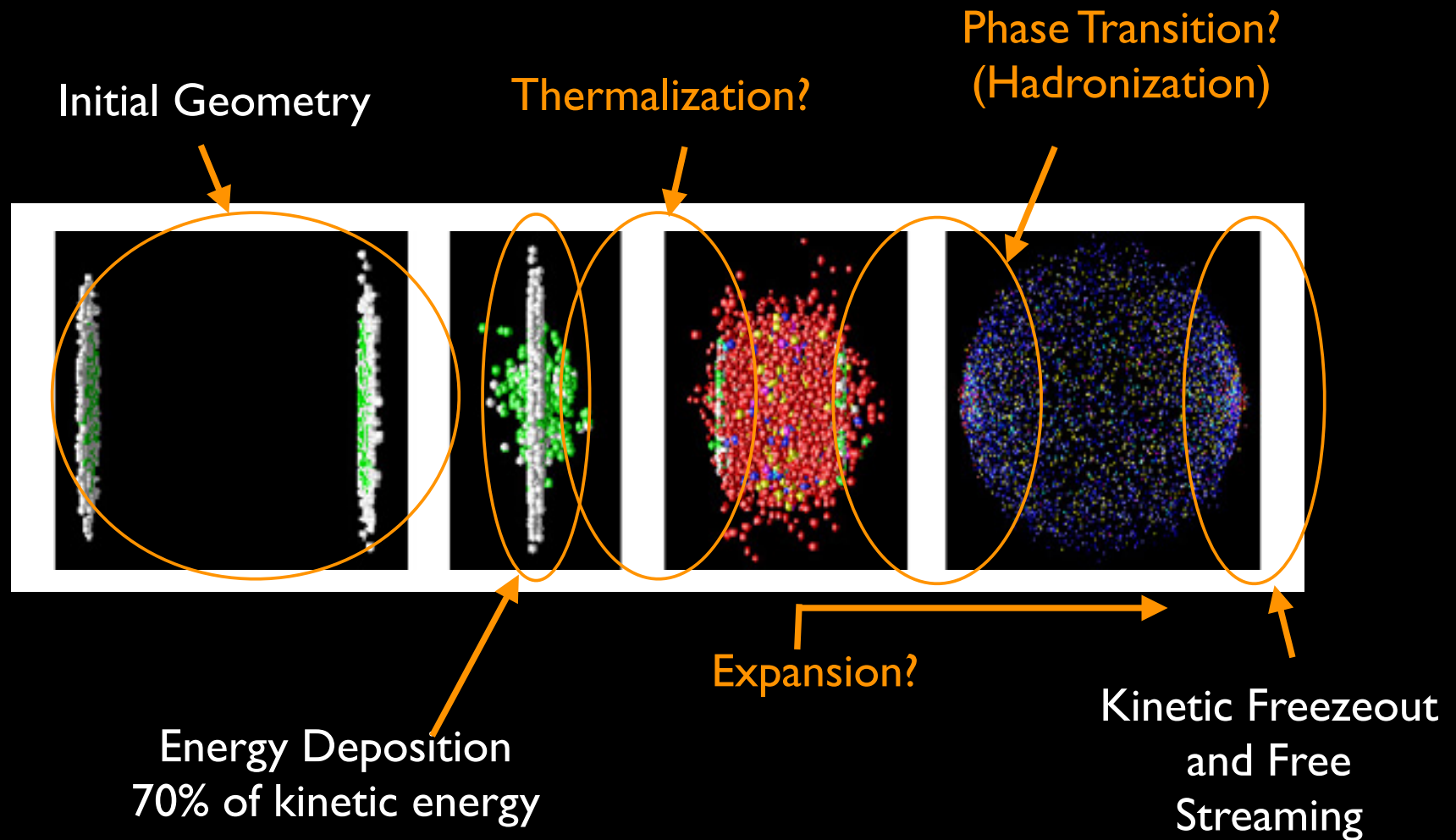
What are the properties of the  
vacuum at  $10^{12}\text{K}$ ?  $\rightarrow$  CMS@LHC



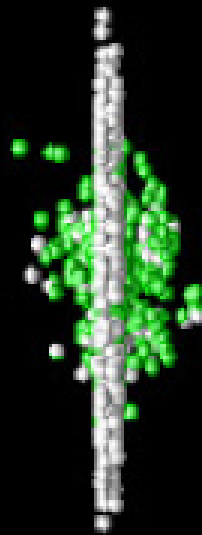
Action density on  $\approx 15\text{fm}^3$  lattice  
D. Leinweber, Adelaide

QCD vacuum

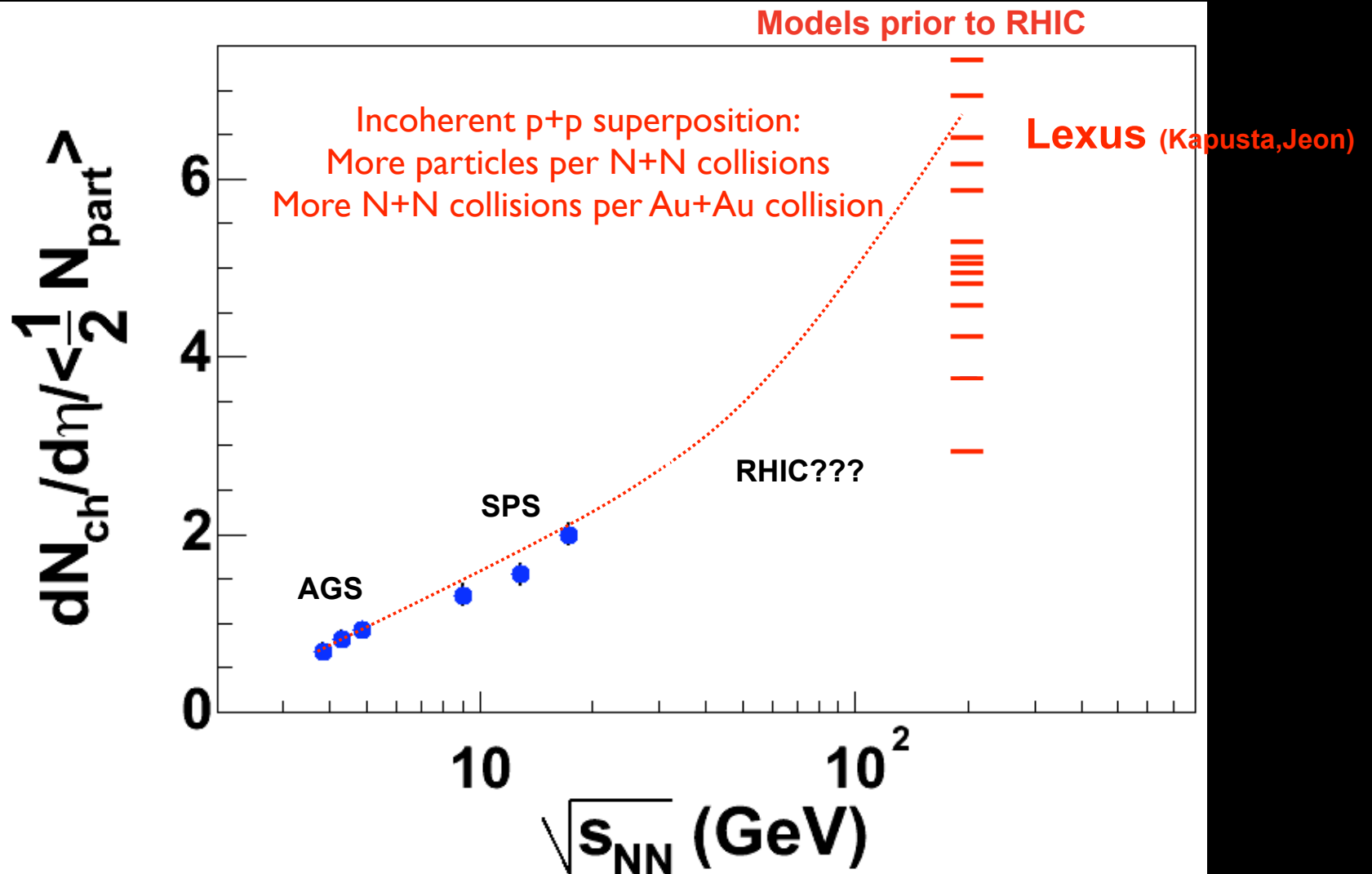
# Time Evolution of a Heavy Ion Collision



# Hadron Multiplicities



# Particle Density near Mid-Rapidity in Au+Au



# First RHIC Physics Paper

## Charged particle multiplicity near mid-rapidity in central Au+Au collisions at $\sqrt{s} = 56$ and 130 AGeV

B.B.Baek<sup>1</sup>, M.D.Baker<sup>2</sup>, D.S.Barton<sup>2</sup>, S.Bastlov<sup>3</sup>, R.Baum<sup>4</sup>, R.R.Betts<sup>1,7</sup>, A.Bialas<sup>4</sup>, R.Bindel<sup>8</sup>, W.Bogucki<sup>1</sup>, A.Budzanowski<sup>1</sup>, W.Busza<sup>2</sup>, A.Carroll<sup>2</sup>, M.Ceglia<sup>2</sup>, Y.-H.Chang<sup>8</sup>, A.E.Chen<sup>8</sup>, T.Coghlan<sup>2</sup>, C.Conner<sup>7</sup>, W.Cryz<sup>4</sup>, B.Dabrowski<sup>1</sup>, M.P.Decowski<sup>2</sup>, M.Despot<sup>2</sup>, P.Fita<sup>2</sup>, J.Fitch<sup>2</sup>, M.Fried<sup>2</sup>, K.Galuska<sup>2</sup>, R.Ganz<sup>7</sup>, E.Garcia-Solis<sup>8</sup>, N.George<sup>2</sup>, J.Godlewski<sup>2</sup>, C.Gomes<sup>2</sup>, E.Griesmayer<sup>2</sup>, K.Gulbrandsen<sup>2</sup>, S.Gushue<sup>2</sup>, J.Hahk<sup>2</sup>, C.Hallwell<sup>7</sup>, P.Haridas<sup>2</sup>, A.Hayes<sup>2</sup>, G.A.Heintzelman<sup>2</sup>, C.Henderson<sup>2</sup>, R.Hollis<sup>7</sup>, R.Holynski<sup>2</sup>, B.Holzman<sup>7</sup>, E.Johnson<sup>8</sup>, J.Kane<sup>2</sup>, J.Katzy<sup>2,4</sup>, W.Kita<sup>2</sup>, J.Kotula<sup>2</sup>, H.Kramer<sup>2</sup>, W.Kucowicz<sup>7</sup>, P.Kulinich<sup>2</sup>, C.Law<sup>2</sup>, M.Lemler<sup>2</sup>, J.Ligocki<sup>2</sup>, W.T.Lin<sup>8</sup>, S.Manly<sup>2,10</sup>, D.McLeod<sup>7</sup>, J.Michalowski<sup>2</sup>, A.Mignerey<sup>8</sup>, J.Mulmenstadt<sup>2</sup>, M.Neal<sup>2</sup>, R.Noulon<sup>7</sup>, A.Olszewski<sup>2,3</sup>, R.Pak<sup>2</sup>, I.C.Park<sup>2</sup>, M.Patel<sup>2</sup>, H.Pernegger<sup>2</sup>, M.Plisko<sup>2</sup>, C.Reed<sup>2</sup>, L.P.Remsberg<sup>2</sup>, M.Reuter<sup>7</sup>, C.Roland<sup>2</sup>, G.Rokand<sup>2</sup>, D.Rose<sup>2</sup>, L.Rosenberg<sup>2</sup>, J.Ryan<sup>2</sup>, A.Samagiri<sup>10</sup>, P.Sarin<sup>2</sup>, P.Sawicki<sup>2</sup>, J.Scaduto<sup>2</sup>, J.Shea<sup>8</sup>, J.Sinacore<sup>2</sup>, W.Skulski<sup>2</sup>, S.G.Steadman<sup>2</sup>, G.S.F.Stephans<sup>2</sup>, P.Steinberg<sup>2</sup>, A.Straszek<sup>2</sup>, M.Stodulski<sup>2</sup>, M.Strz<sup>2</sup>, Z.Stopa<sup>2</sup>, A.Sukhanov<sup>2</sup>, K.Surowicka<sup>2</sup>, J.-L.Tang<sup>2</sup>, R.Teng<sup>2</sup>, A.Trzupak<sup>2</sup>, C.Vale<sup>2</sup>, G.J.van Nieuwenhuizen<sup>2</sup>, R.Verderi<sup>2</sup>, B.Wadsworth<sup>2</sup>, F.L.H.Wold<sup>2</sup>, B.Wosiek<sup>2</sup>, K.Wozniak<sup>2</sup>, A.H.Wuossma<sup>4</sup>, B.Wyslouch<sup>2</sup>, K.Zaleski<sup>2</sup>, P.Zychowski<sup>2</sup> (PHOBOS collaboration)

<sup>1</sup> Physics Division, Argonne National Laboratory, Argonne, IL 60439-4843

<sup>2</sup> Chemistry Department, Brookhaven National Laboratory, Upton, NY 11973-5000

<sup>3</sup> Institute of Nuclear Physics, Kraków, Poland

<sup>4</sup> Department of Physics, Jagellonian University, Kraków, Poland

<sup>5</sup> Laboratory for Nuclear Science, Massachusetts Institute of Technology, Cambridge, MA 02139-4307

<sup>6</sup> Department of Physics, National Central University, Chung-Li, Taiwan

<sup>7</sup> Department of Physics, University of Illinois at Chicago, Chicago, IL 60607-7039

<sup>8</sup> Department of Chemistry, University of Maryland, College Park, MD 20742

<sup>9</sup> Department of Physics and Astronomy, University of Rochester, Rochester, NY 14627

<sup>10</sup> Department of Physics, Yale University, New Haven, CT 06520

(July 21, 2000)

We present the first measurement of pseudorapidity densities of primary charged particles near mid-rapidity in Au+Au collisions at  $\sqrt{s} = 56$  and 130 AGeV. For the most central collisions, we find the charged particle pseudorapidity density to be  $dN/d\eta|_{|\eta|<1} = 408 \pm 12(\text{stat}) \pm 30(\text{sys})$  at 56 AGeV and  $588 \pm 12(\text{stat}) \pm 38(\text{sys})$  at 130 AGeV, values that are higher than any previously observed in nuclear collisions. Compared to proton-antiproton collisions, our data show an increase in the pseudorapidity density per participant by more than 40% at the higher energy.

PACS numbers: 25.75.-q

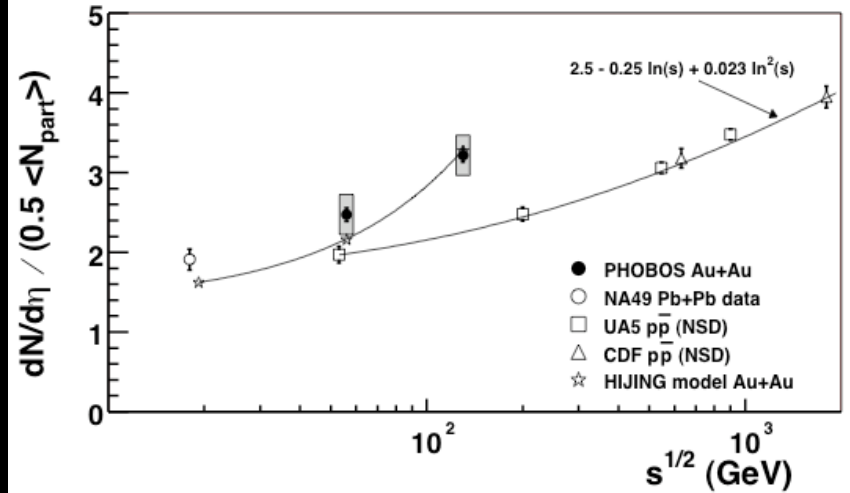
In June 2000, the Relativistic Heavy-Ion Collider (RHIC) at Brookhaven National Laboratory delivered the first collisions between Au nuclei at the highest center of mass energies achieved in the laboratory to date. In this paper we present data taken with the PHOBOS detector during the first collider run at energies of  $\sqrt{s} = 56$  and 130 AGeV. The ultimate goal of our work is to understand the behavior of strongly interacting matter at conditions of extreme density and temperature. Quantum chromodynamics (QCD), the fundamental theory of strong interactions, predicts that for sufficiently high energy density a new state of matter will be formed, the so-called quark-gluon plasma (QGP) [1]. The measurements shown here represent the first step toward the development of a full picture of the dynamical evolution of nucleus-nucleus collisions at RHIC energies.

Studying the dependence of charged particle densities

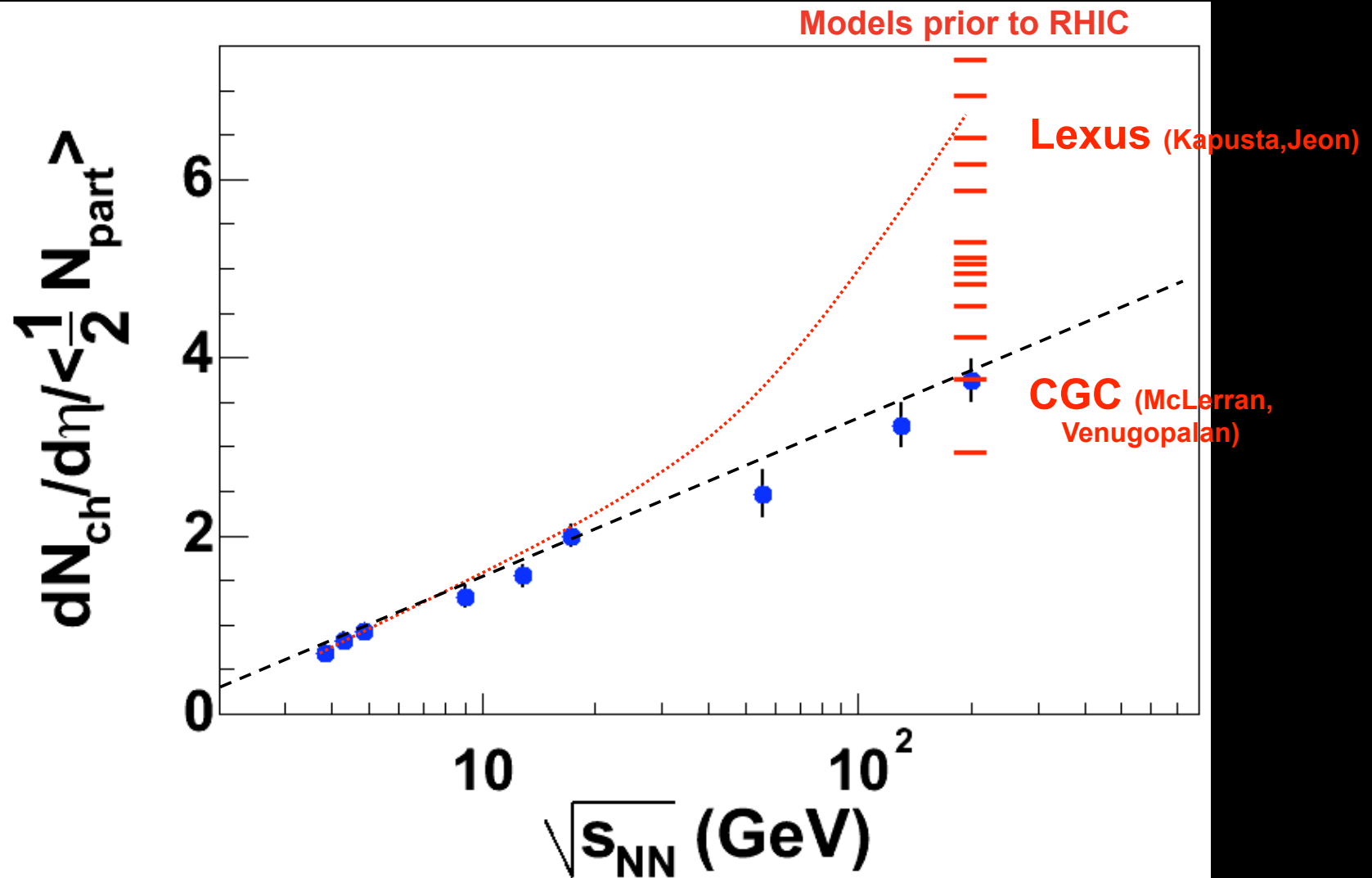
on energy and system size provides information on the interplay between hard parton-parton scattering processes, which can be calculated using perturbative QCD, and soft processes, which are treated by phenomenological models that describe the non-perturbative sector of QCD. Predictions for multi-particle production in high-energy heavy-ion collisions, obtained from a variety of models, typically vary by up to a factor of two [2].

In this letter we report data for the most central Au+Au collisions detected in our apparatus. We have determined the energy dependence of the density of primary charged particles emitted near  $90^\circ$  to the beam axis, characterized by the pseudorapidity density  $dN/d\eta|_{|\eta|<1}$ , where  $\eta = -\ln \tan(\theta/2)$  and  $\theta$  is the polar angle from the beam axis. These data provide the first means to constrain models of heavy-ion collisions at RHIC energies. They will allow the extraction of basic information about the initial conditions in these collisions, in particular the energy density, and thus form an essential element for the proper prediction or description of other observables.

The PHOBOS detector employs silicon pad detectors to perform tracking, vertex detection and multiplicity measurements. Details of the setup and the layout of the silicon sensors can be found elsewhere [3,4]. For the initial running period of the accelerator only a small fraction of the full setup was installed. It included the first 6 layers of the silicon spectrometer (SPEC), part of the two-layer silicon vertex detector (VTX) and

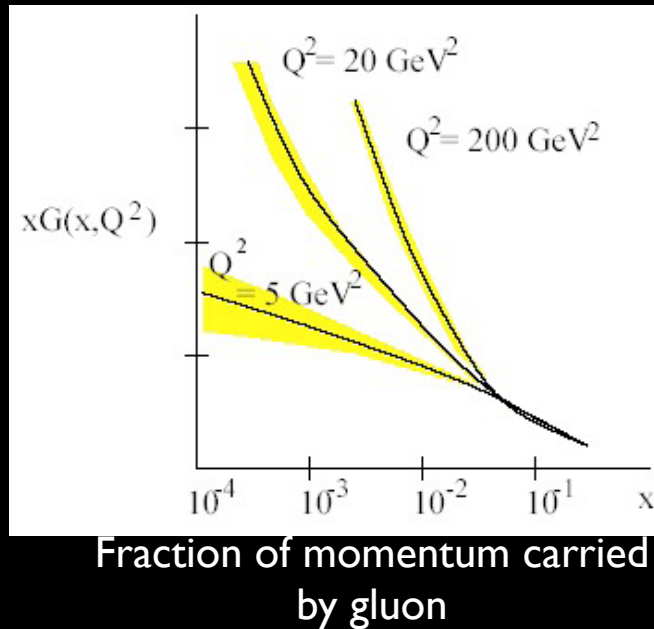


# Particle Density near Mid-Rapidity



# Parton Saturation

Gluon density in proton

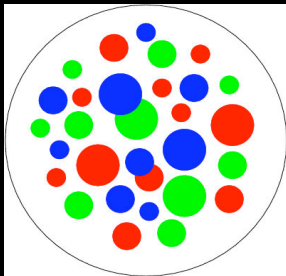


Idea(s):

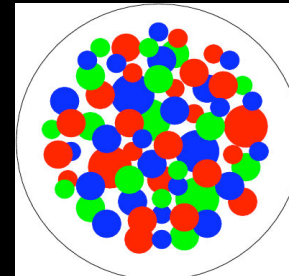
Entropy is not “created” but  
“liberated” from gluon distributions

Gluon density increases with  $x$ ,  $Q^2$

BUT: Gluons interact, limiting growth  
of gluon densities: Saturation



Low energy

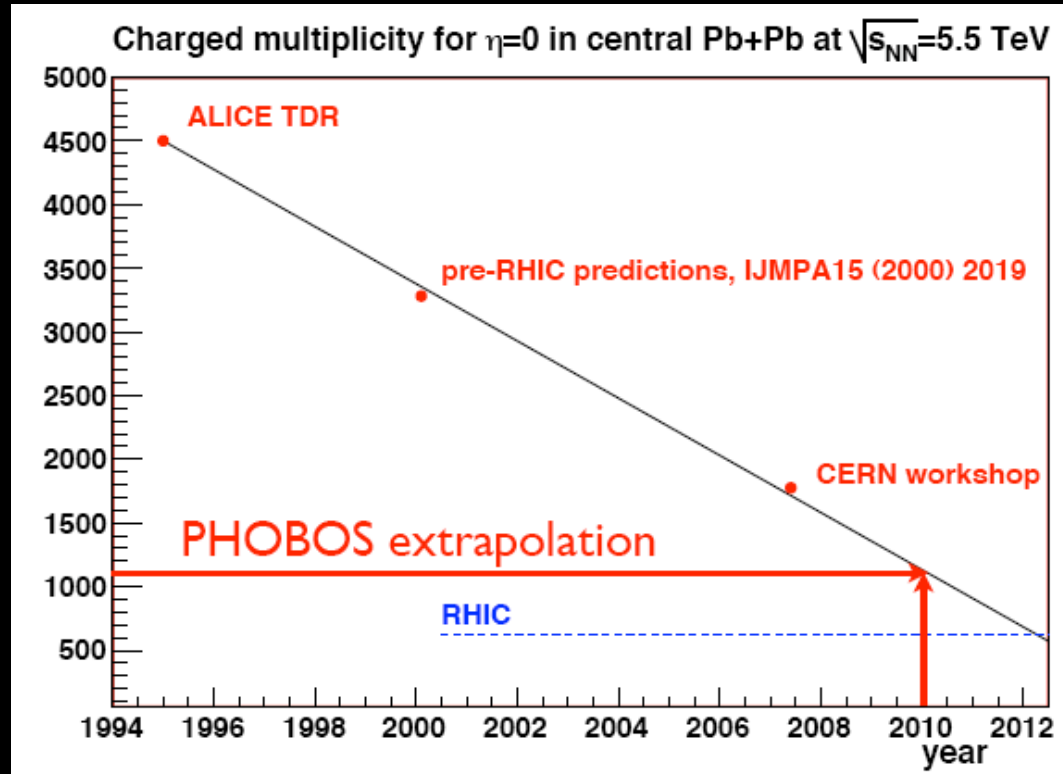


High energy



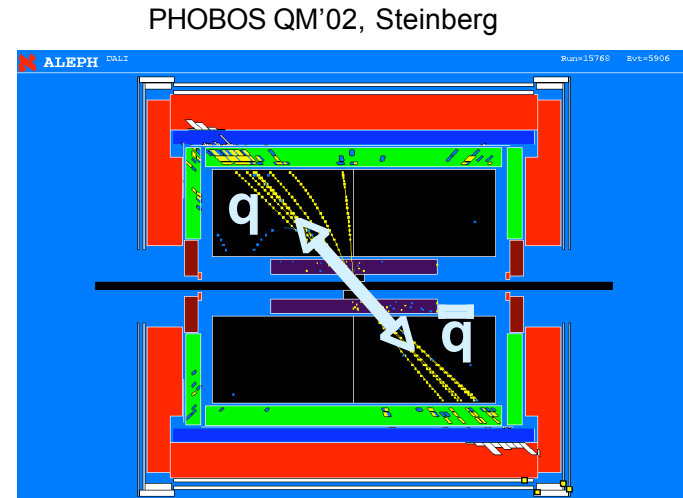
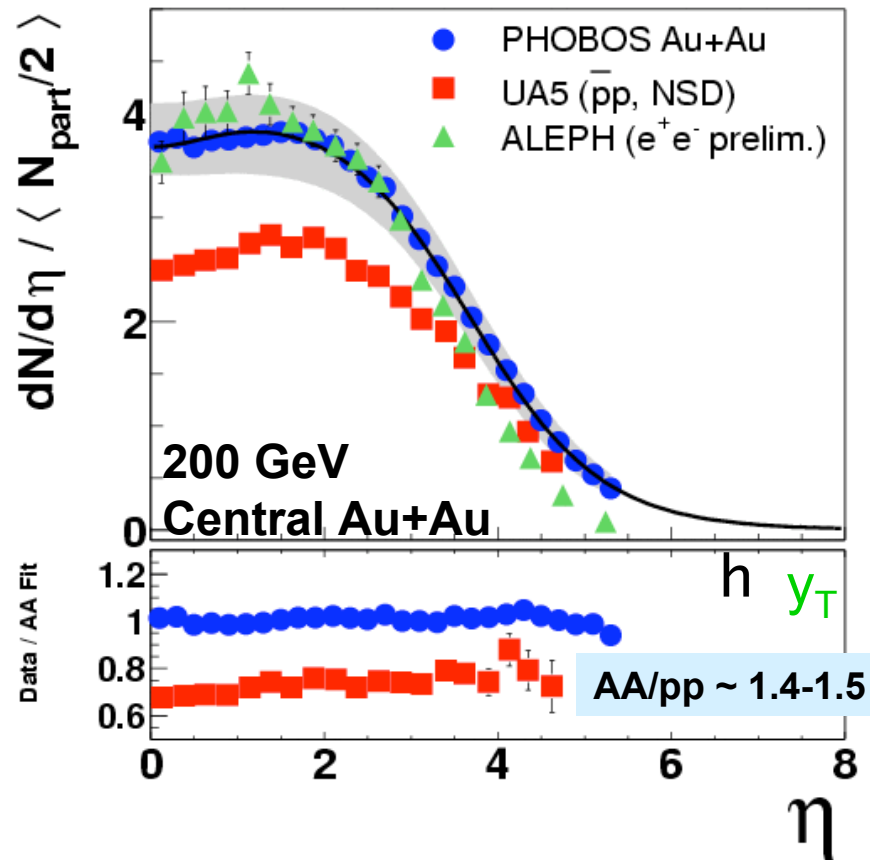
# One important lesson

from Nestor Armesto, QM 2008



Don't bet your experiment on predictions....

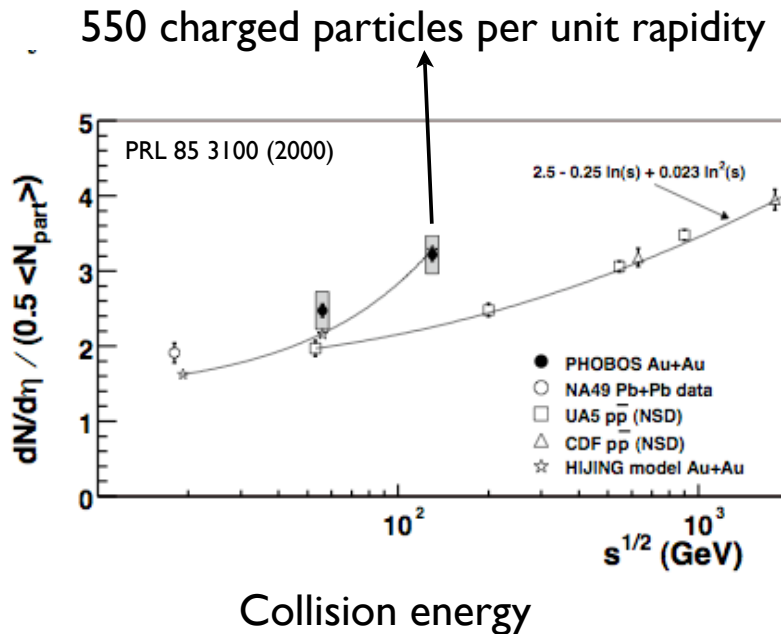
# $e^+e^-$ , p+p, A+A Correspondence ?



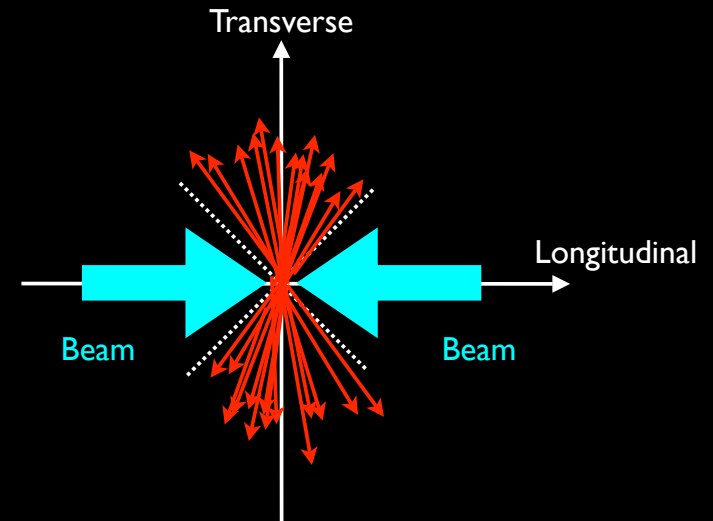
$e^+e^-$  measures  $dN/dy_T$   
(rapidity relative to  
“thrust” axis)

Surprising agreement in shape between AA/ $e^+e^-$ /pp

Angular particle density near  $90^\circ$   
normalized per participant



RHIC delivered first collisions on June 12th, 2000  
PHOBOS submitted first paper July 19th, 2000



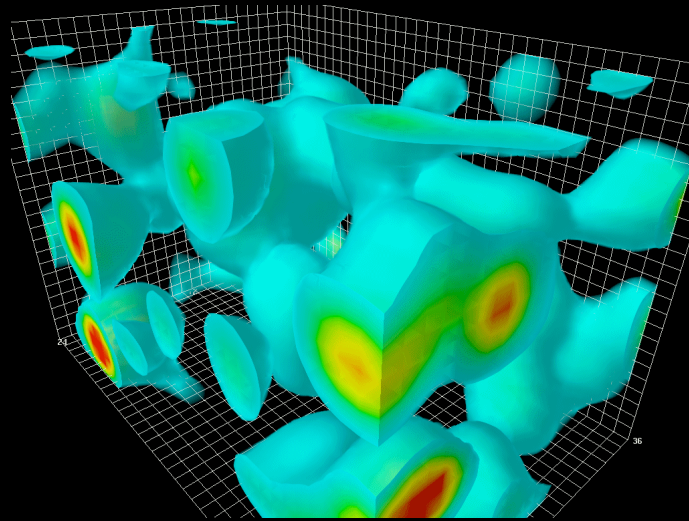
Use “energy flow” from  
longitudinal (=beam) to transverse direction  
to estimate energy/volume

$$\frac{1000 \text{ particles} \times 0.5 \text{ GeV/particle}}{\pi \times (7 \text{ fm})^2 \times 1 \text{ fm}} \approx 3 \text{ GeV/fm}^3$$

Much larger than  $\epsilon_{\text{crit}} \approx 0.7 \text{ GeV/fm}^3$

**But: Equilibration?**

What are the transport properties of  
the vacuum at  $10^{14}$  K?



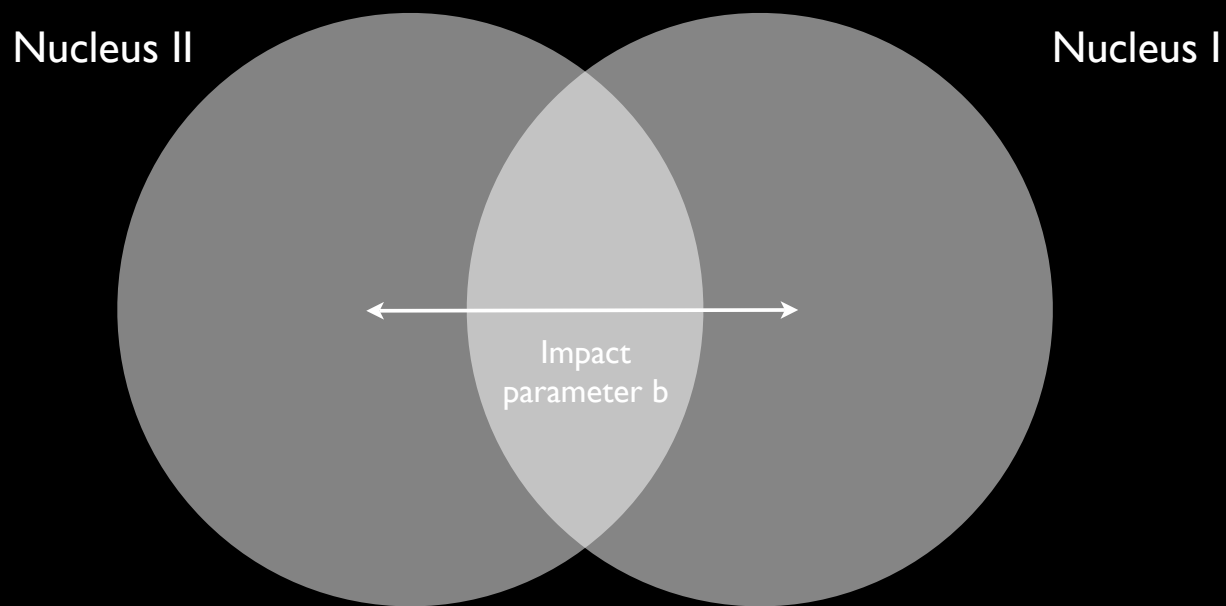
Action density on  $\approx 15\text{fm}^3$  lattice  
D. Leinweber, Adelaide

e.g. can one measure or calculate the  
viscosity of a system at  $T > 200\text{MeV}$ ?

first we need to show that we indeed  
produce an interacting system at such  
temperatures

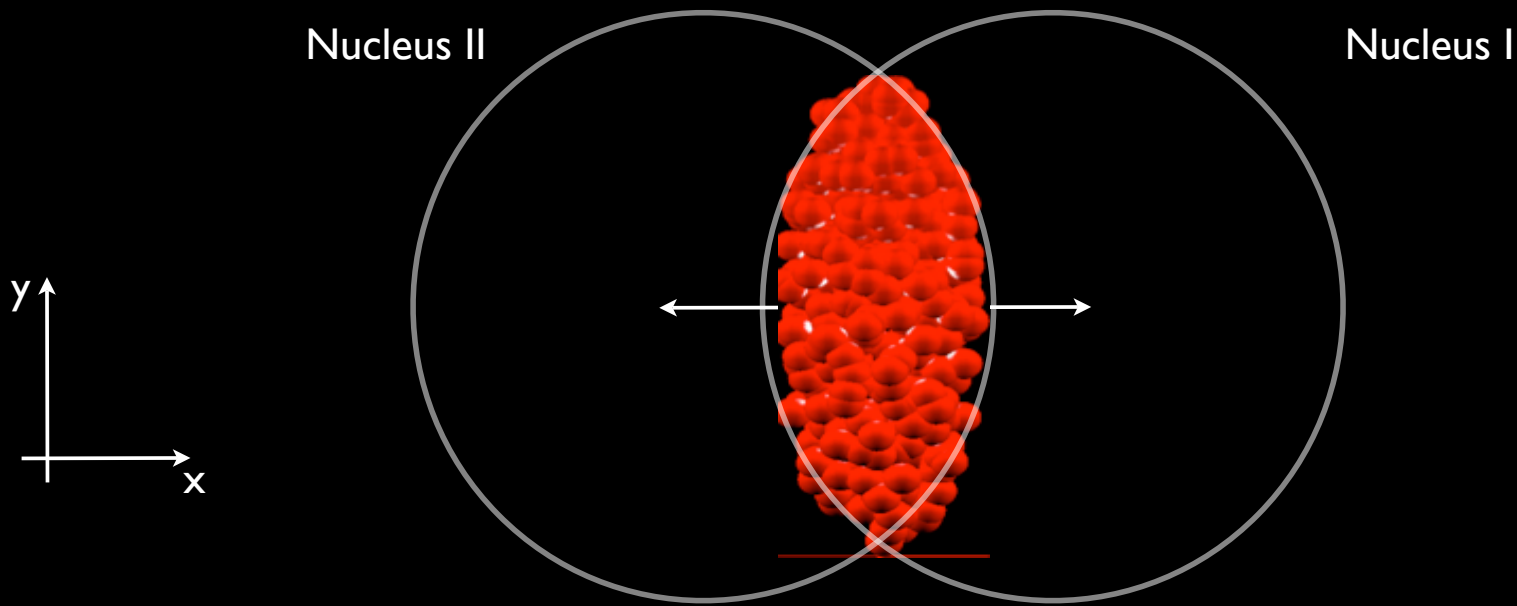
# How do we prove that we make “matter”?

Non-central collision (Transverse plane)



# How do we prove that we make “matter”?

Non-central collision (Transverse plane)



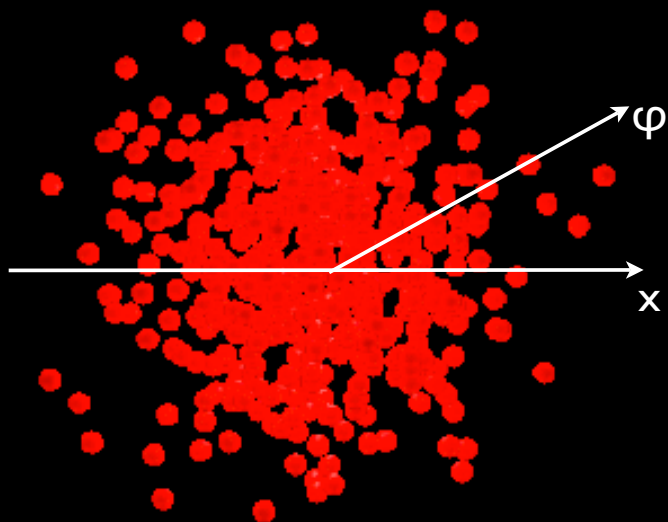
“Hot” overlap zone is asymmetric in azimuthal angle

Define:  $\epsilon_{std} = \frac{\sigma_y^2 - \sigma_x^2}{\sigma_y^2 + \sigma_x^2}$  “Initial State Eccentricity”

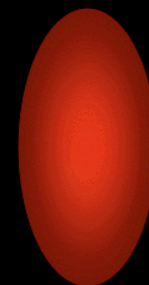
Non-interacting particles



Non-interacting particles

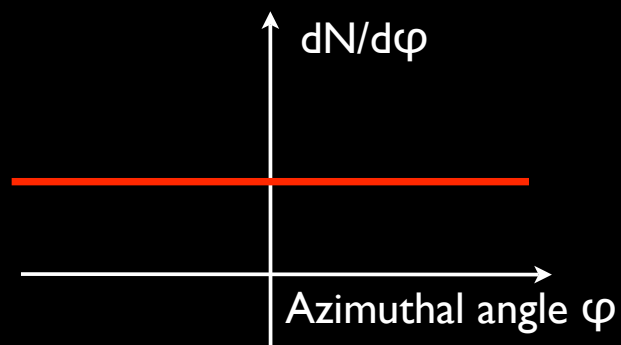


Collective expansion of Matter



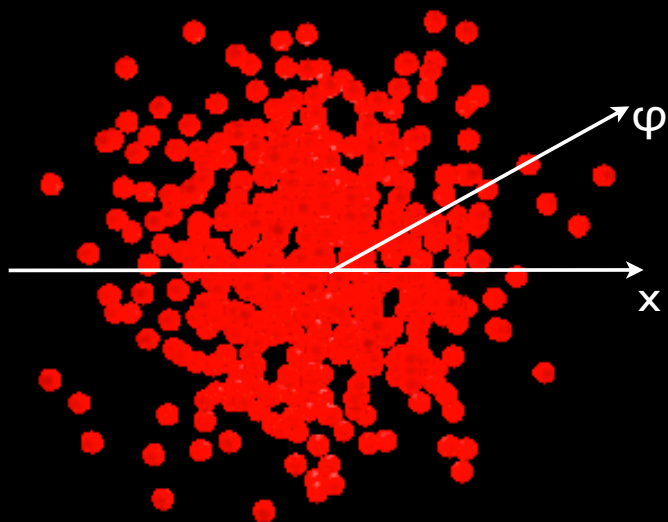
Shape information is not transferred to  
momentum space

Flat azimuthal distribution



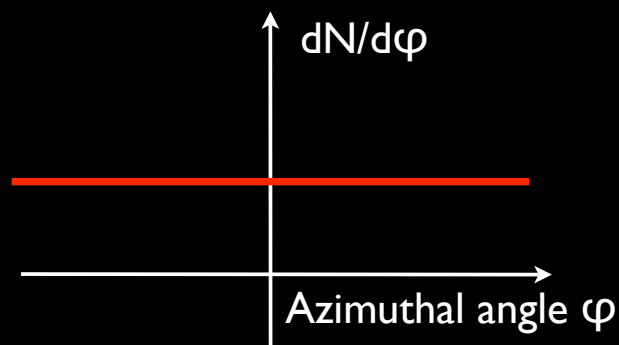


## Non-interacting particles

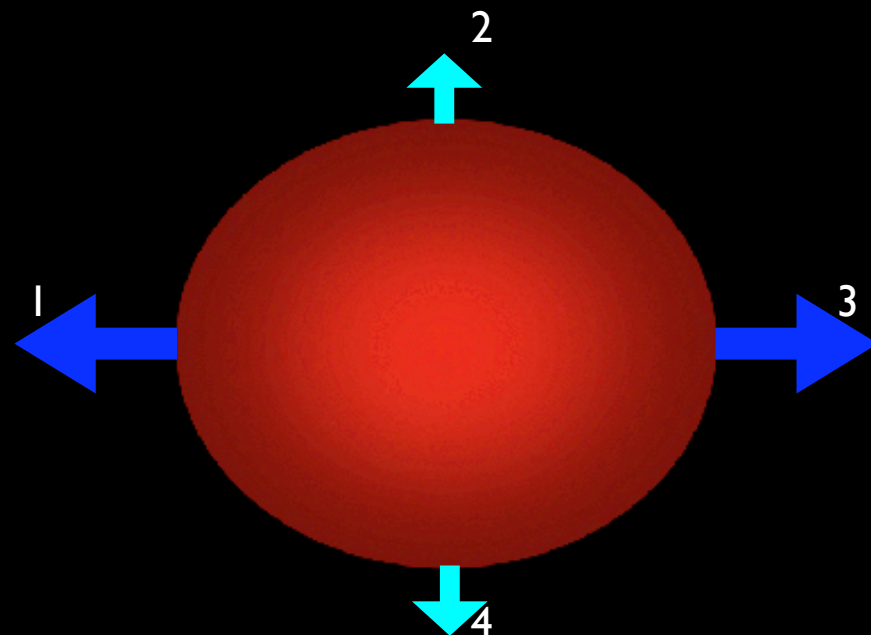


Shape information is not transferred to momentum space

Flat azimuthal distribution

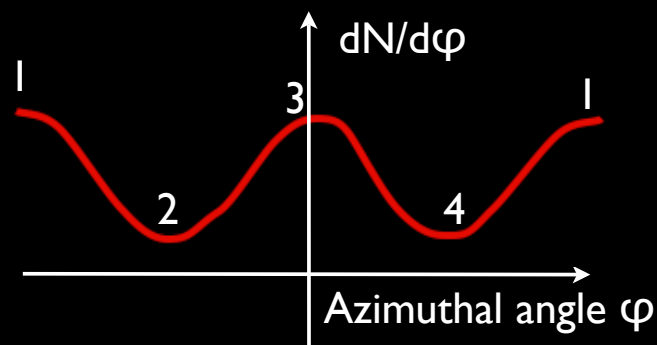


## Collective expansion of Matter



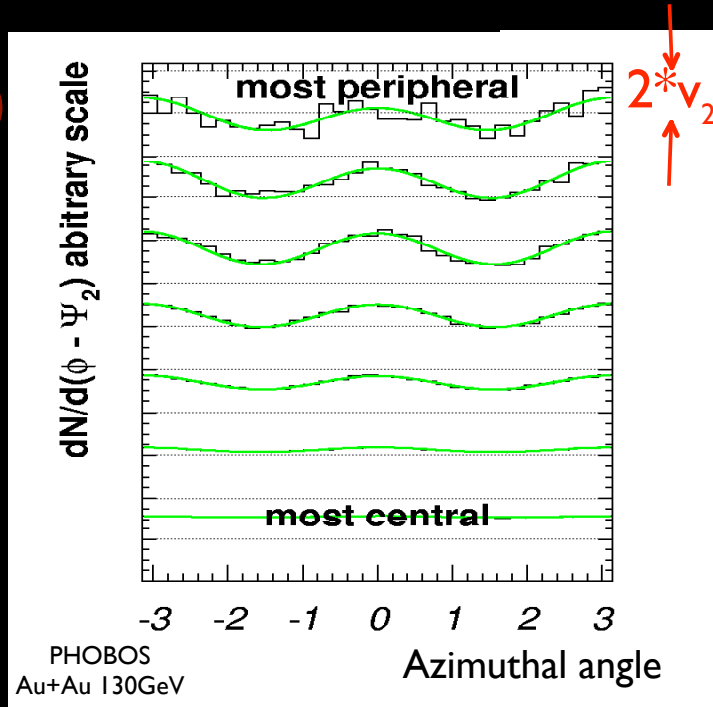
Shape information transformed into momentum space

$\cos(2\varphi)$  modulation of azimuthal distribution



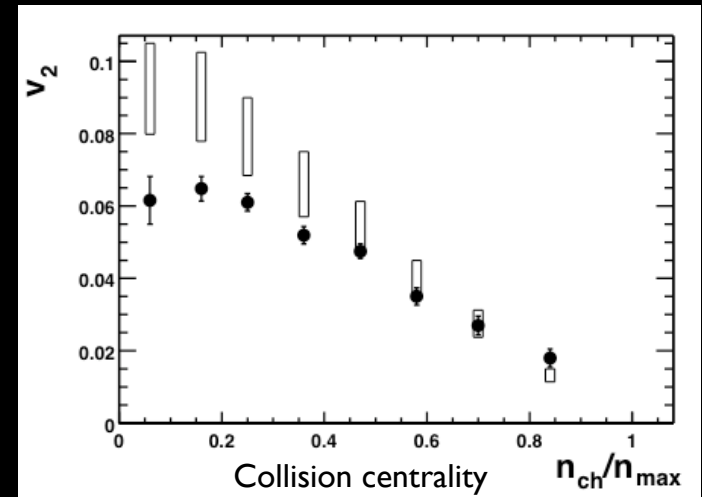
# How do we prove that we make “matter”?

Azimuthal distribution  
$$dN/d\varphi = 1 + 2 v_2 \cos(2(\varphi - \varphi_0))$$



## “Elliptic Flow”

STAR PRL 2000



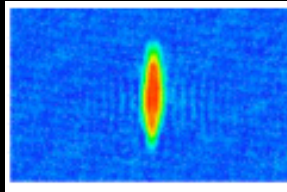
Peripheral collisions

central collisions

The initial anisotropy in coordinate space is translated into momentum space: Interactions → Equilibration (?)

# Hydrodynamic Evolution

Initial State



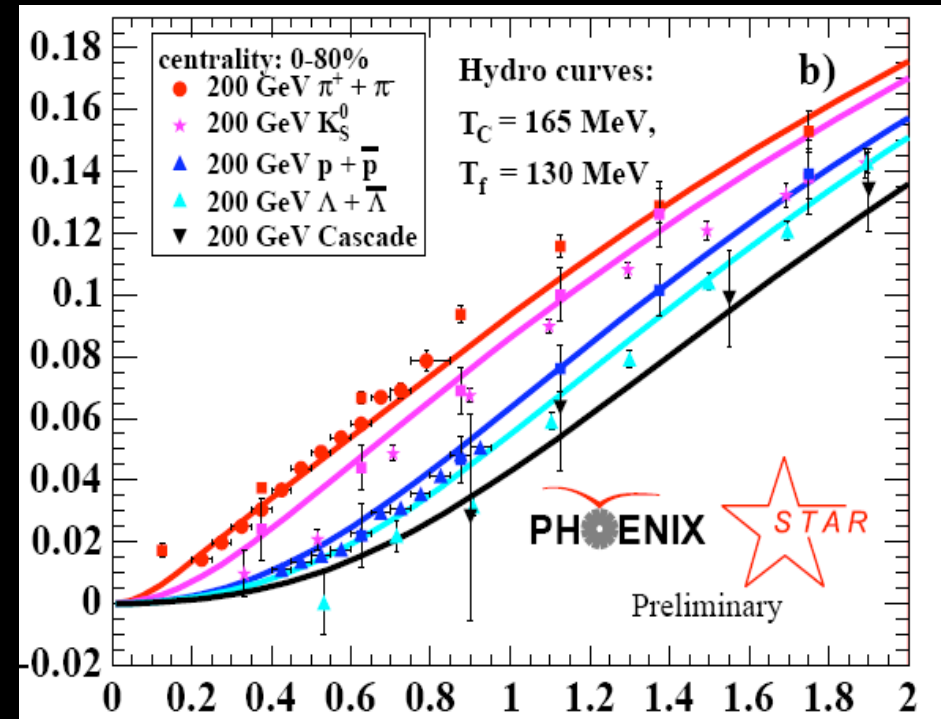
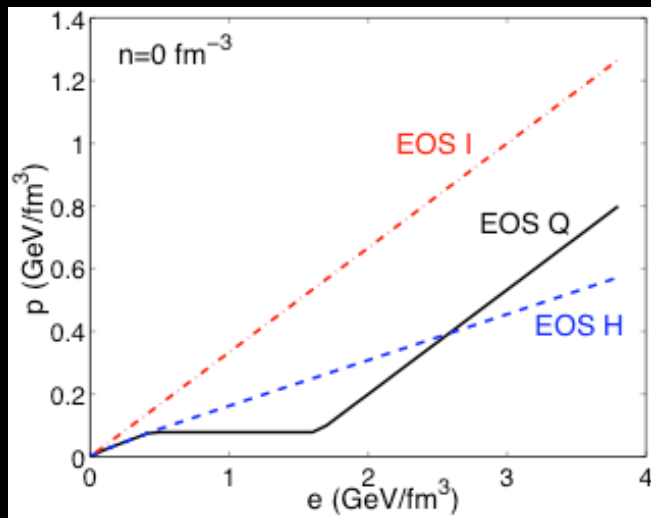
Energy/Momentum  
Conservation

$$\partial_\mu T^{\mu\nu} = 0$$

$$\partial_\mu j^\mu = 0$$

Baryon number  
Conservation

Equ. of State



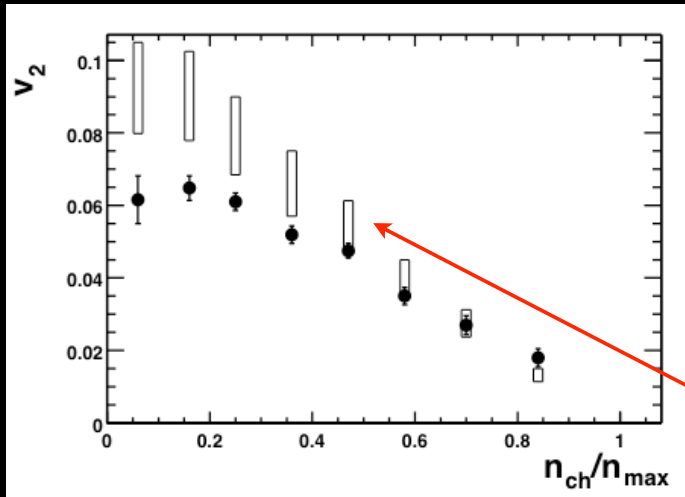
=

Ideal Hydro  
local equilibrium  
vanishing  $\lambda_{\text{MFP}}$   
no viscosity

# Hydrodynamics

“Ideal hydrodynamics”

STAR PRL 2000



Assumption:

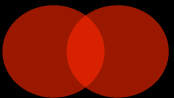
Shortly after initial collision ( $< 1-2 \text{ fm}/c$ ) a system in local equilibrium with very small mean free path is created

Local equilibrium  $\Leftrightarrow$  small  $\lambda_{mfp} \Leftrightarrow$  small shear viscosity

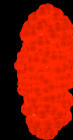
$v_2 \propto \epsilon$  (i.e. initial geometric eccentricity)

Mid-central data reach hydro prediction (calculated using  $\lambda_{mfp} = 0$ )

Peripheral collisions



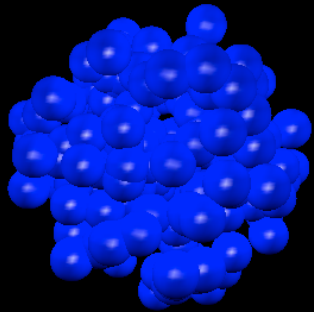
central collisions



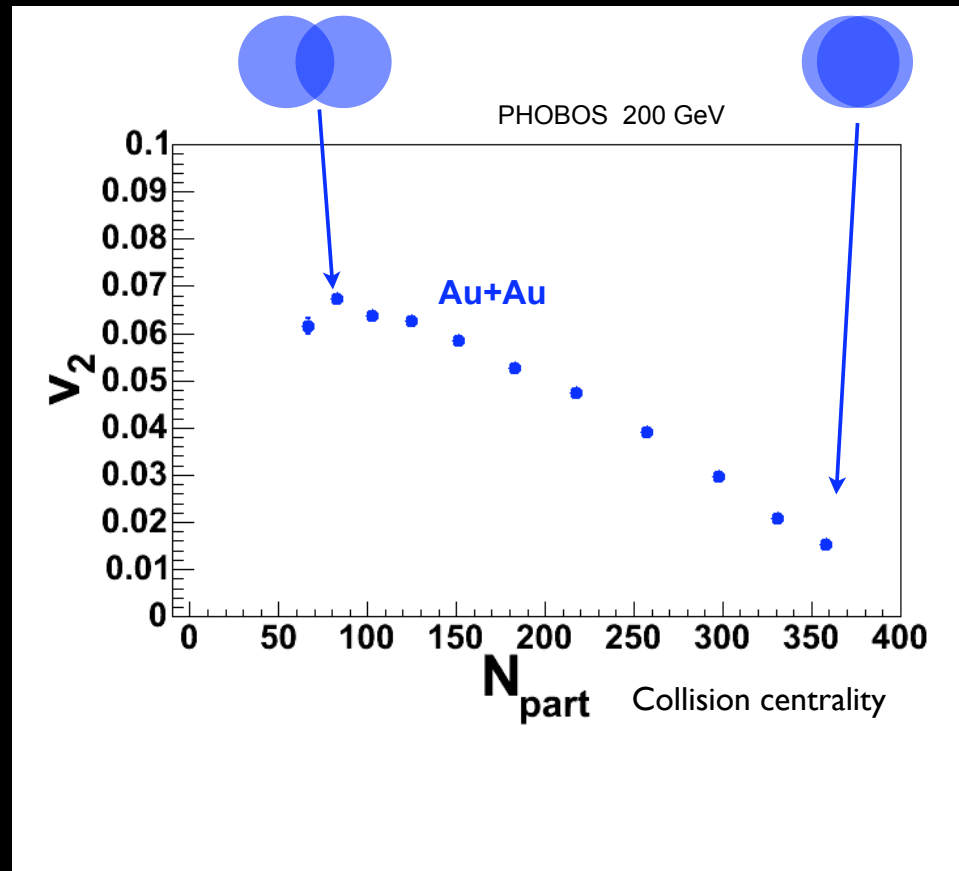
Once shape info is lost in free streaming, can't be recovered

# Elliptic Flow and Geometry, I

Test connection between geometry and elliptic flow  
by comparing Au+Au to Cu+Cu

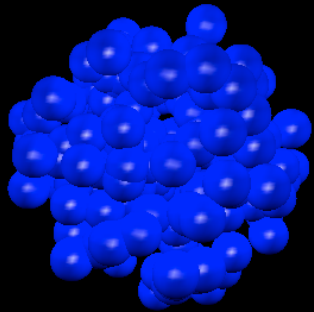


Gold  
 $A=197$

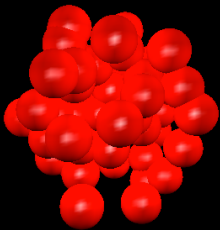


# Elliptic Flow and Geometry, I

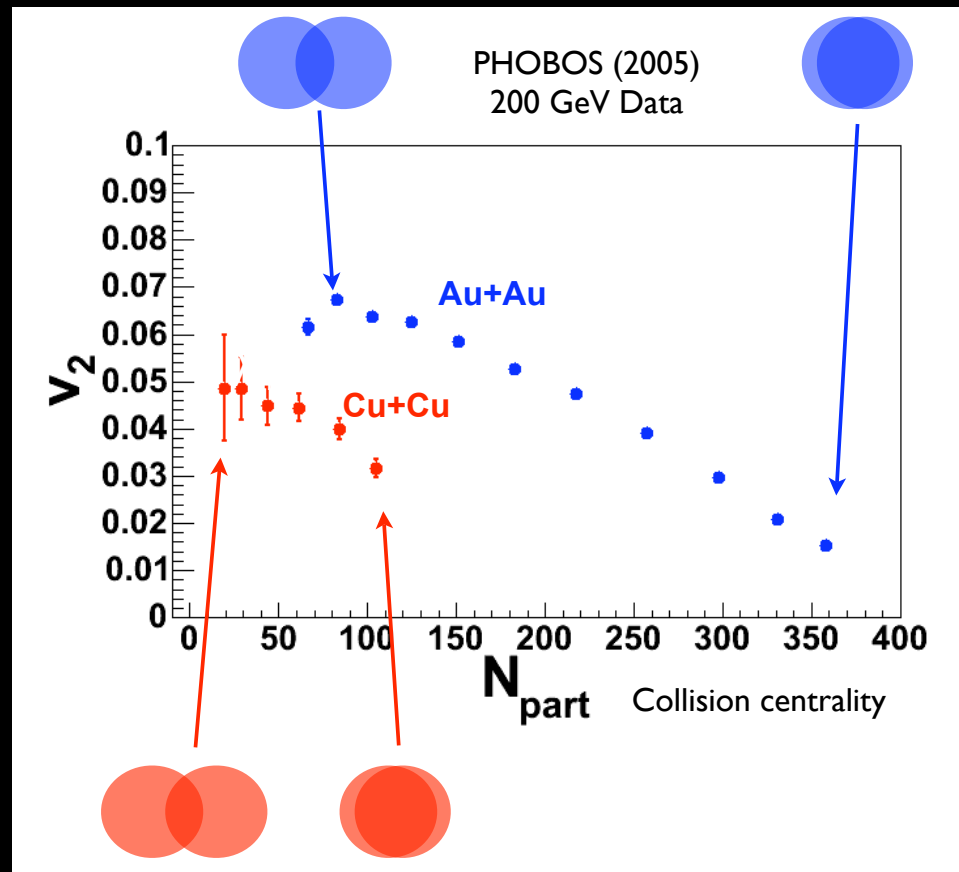
Test connection between geometry and elliptic flow  
by comparing Au+Au to Cu+Cu



Gold  
 $A=197$



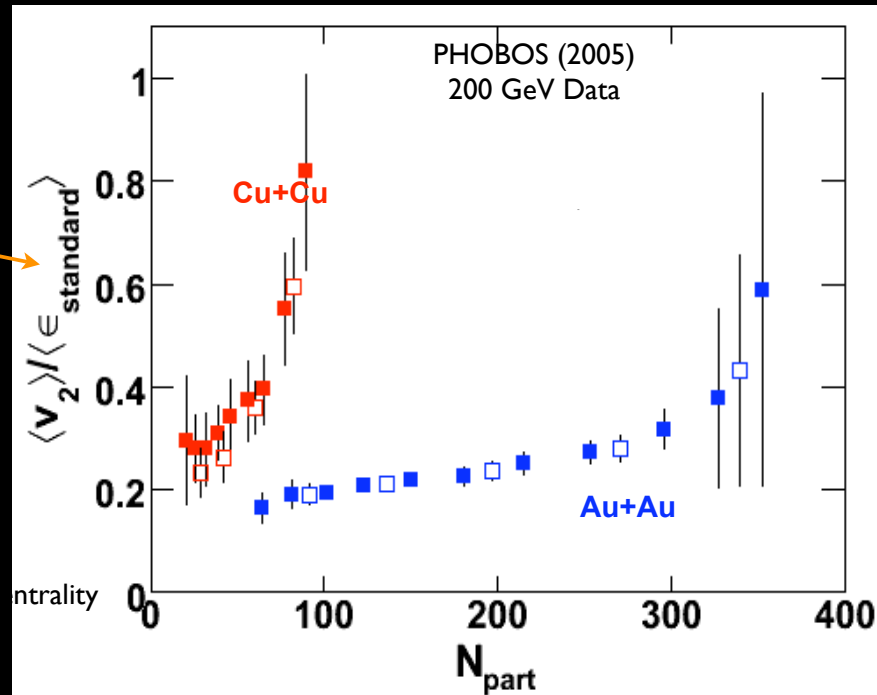
Copper  
 $A=64$



$v_2$  is large even for central Cu+Cu

# Challenge: System Size Scaling

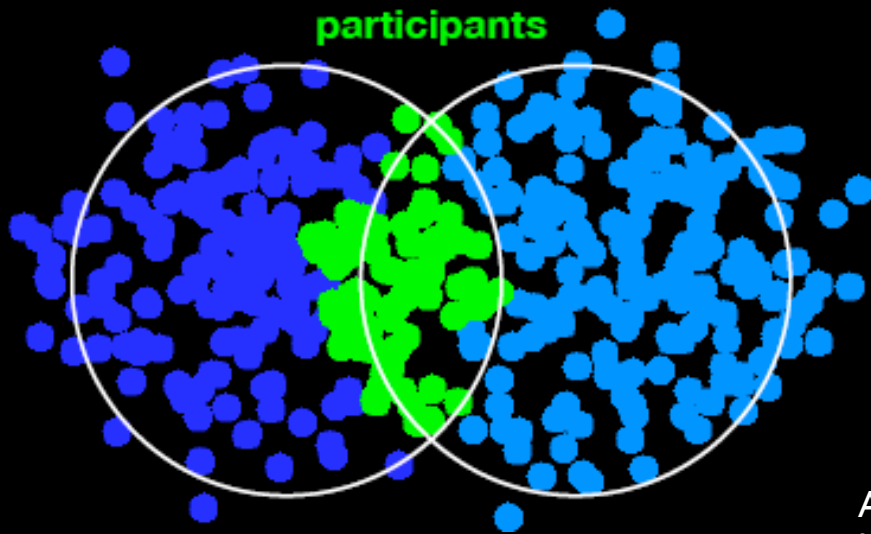
$$\epsilon_{std} = \frac{\sigma_y^2 - \sigma_x^2}{\sigma_y^2 + \sigma_x^2}$$



For same  $N_{part}$  ( $\sim$  same initial density),  $v_2/\epsilon_{std}$  is much larger in Cu+Cu than in Au+Au collisions

# Re-thinking $\epsilon$

At fixed  $b$



In **Glauber MC** model, geometry is sampled by finite number of nucleons



Geometry varies from event-to-event,  
even at fixed  $b$

Aguiar, Hama, Kodama, Osada, hep-ph/0106266 (QM 2001)

Miller, Snellings, nucl-ex/0312008

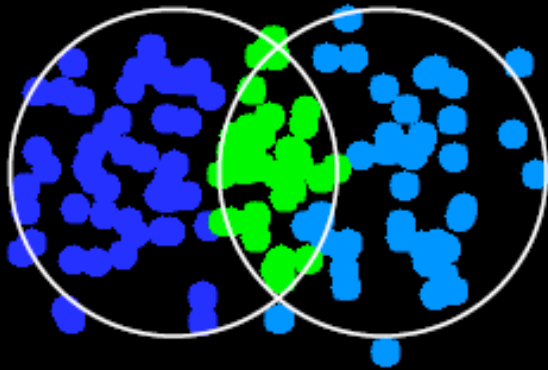
(4 citations until 2005, 28 since then)

Broniowski et al, arXiv:0706.4266

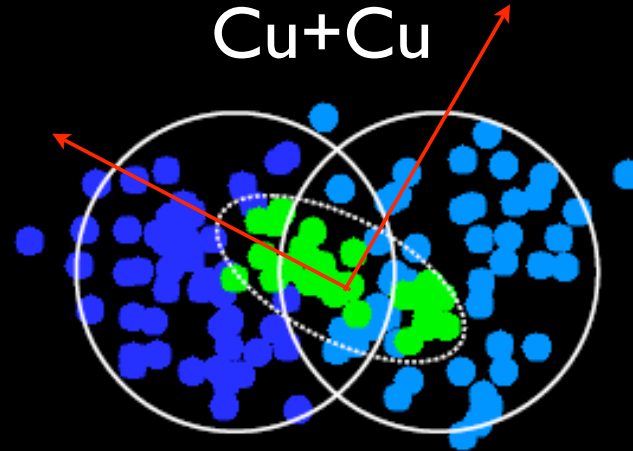


# Collision Geometry Fluctuations

Cu+Cu



Cu+Cu



Plots from Richard Bindel, Maryland,  
using PHOBOS Glauber MC

$$\epsilon_{part} = \frac{\sigma_y'^2 - \sigma_x'^2}{\sigma_y'^2 + \sigma_x'^2} = \frac{\sqrt{(\sigma_y^2 - \sigma_x^2)^2 + 4(\sigma_{xy}^2)^2}}{\sigma_y^2 + \sigma_x^2}$$

“Participant Eccentricity”

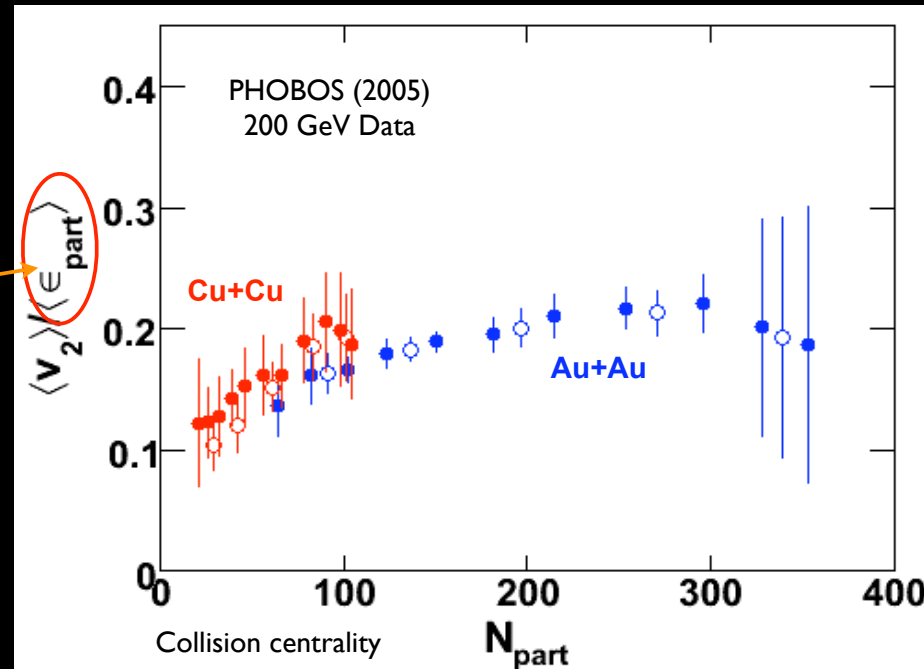
PHOBOS 2005, see also

Broniowski et al, arXiv:0706.4266

If flow is driven by initial matter distribution,  
the orientation (and shape) of that distribution  
should determine direction and magnitude of flow

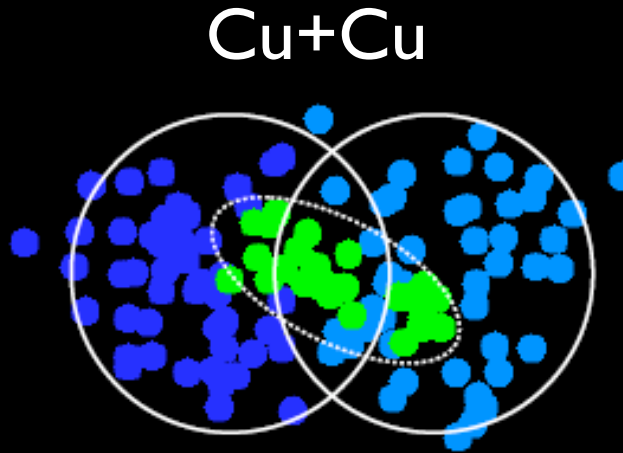
# System Size Scaling

$$\epsilon_{part} = \frac{\sqrt{(\sigma_y^2 - \sigma_x^2)^2 + 4(\sigma_{xy}^2)^2}}{\sigma_y^2 + \sigma_x^2}$$



Re-interpretation of Glauber MC initial states  
yields  $v_2$  scaling between Cu+Cu and Au+Au

# Collision Geometry Fluctuations



How do we know the Glauber shapes and  
shape fluctuations are real?

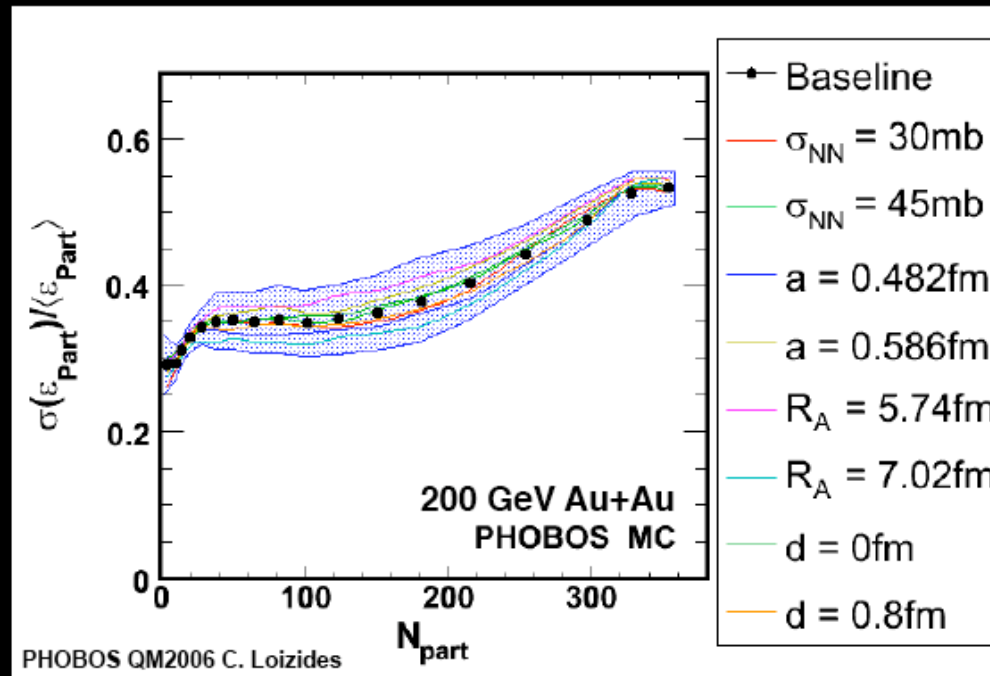
Measure them directly!

If  $v_2 \propto \epsilon$ , then:

$$\frac{\sigma(v_2)}{\langle v_2 \rangle} = \frac{\sigma(\epsilon)}{\langle \epsilon \rangle}$$

i.e. relative fluctuations in  $v_2$  should be  
determined by relative fluctuations in  $\epsilon$

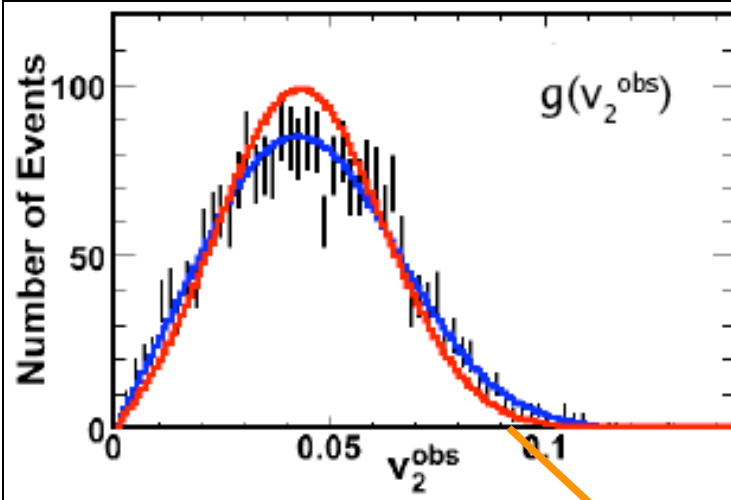
# $\epsilon_{\text{part}}$ Fluctuations in Glauber MC



Large event-by-event variation of  $\epsilon_{\text{part}}$  ( $\sim 40\%$ )

Robust against variation of Glauber MC parameters

# Extracting $v_2$ Fluctuations



Measured

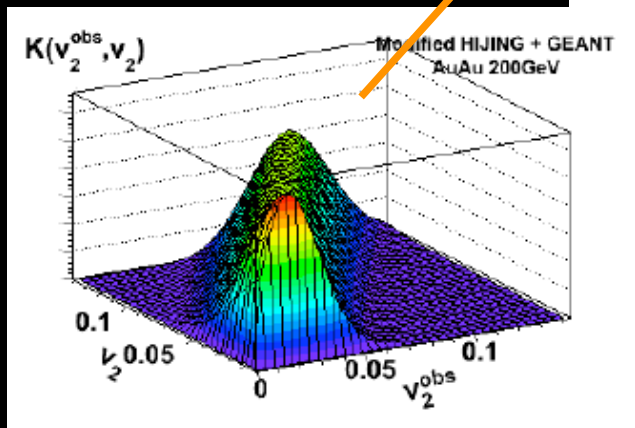
PHOBOS: event-by-event fit of  $(v_2, \phi_0)$  over  $\sim 4\pi$

$$g(v_2^{\text{obs}}) = \int_0^1 K(v_2^{\text{obs}}, v_2) f(v_2) dv_2$$

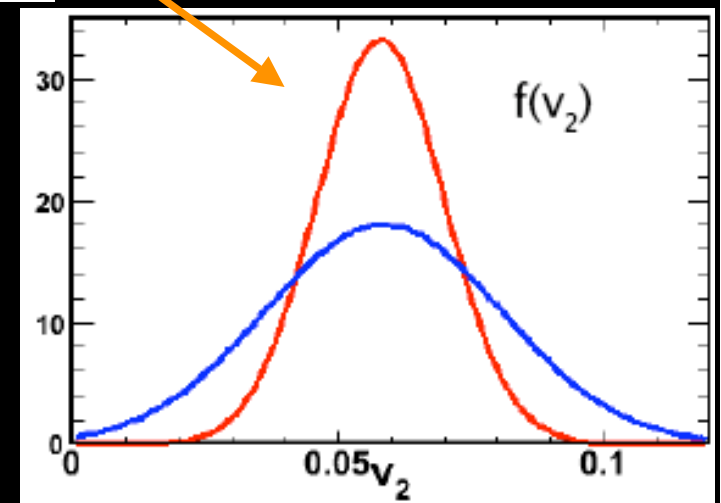
Measured

Constructed from MC

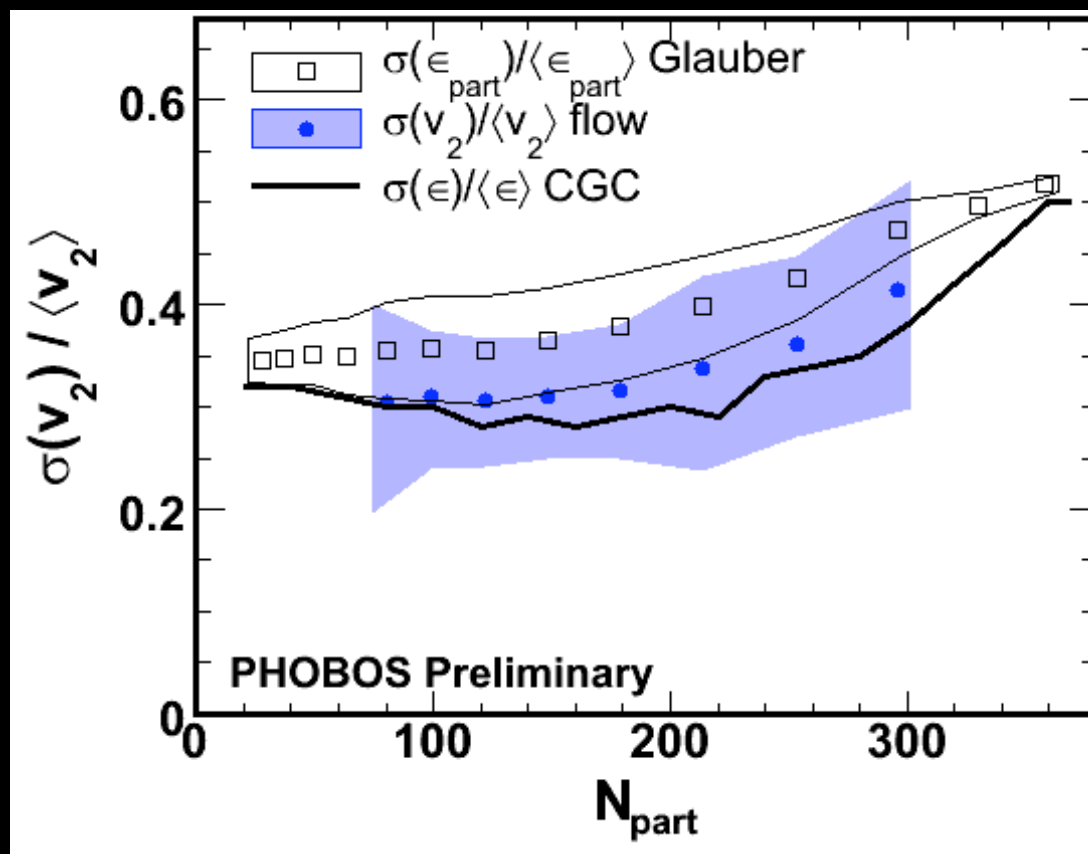
Corrected



Correction function

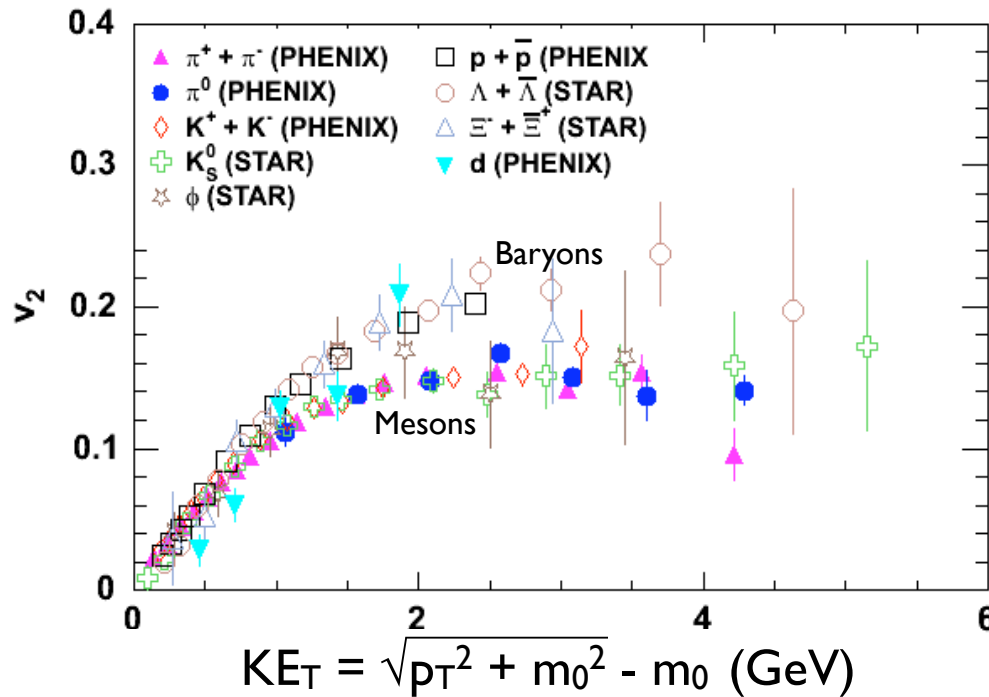


# Observed Elliptic Flow Fluctuations



# What is the nature of this matter?

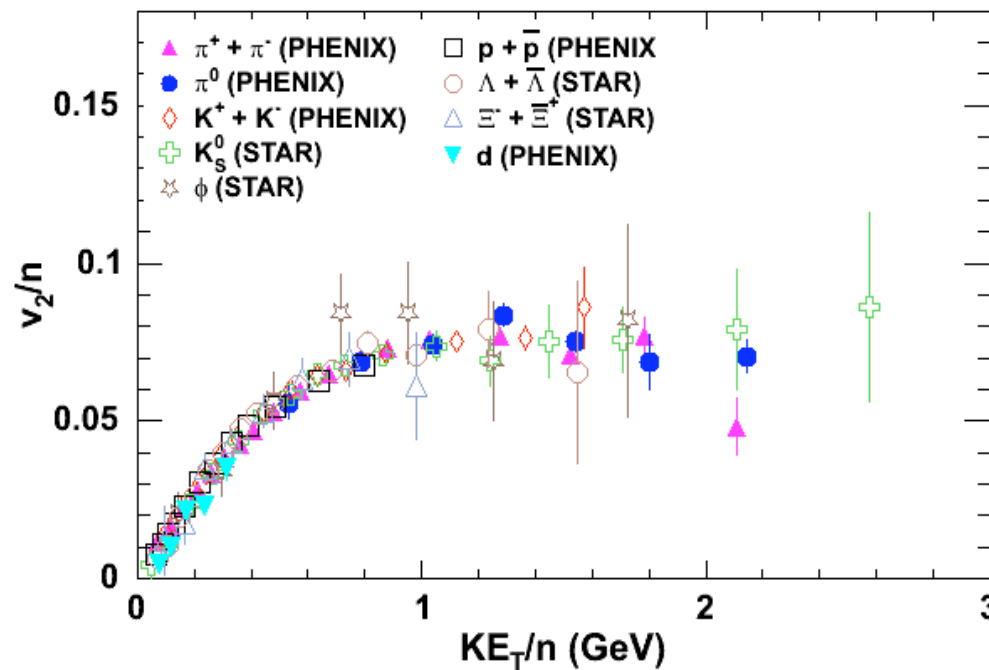
Plot from M. Isaah  
CIPANP '06



Elliptic flow as a function of  
“transverse kinetic energy”

# What is the nature of this matter?

Plot from M. Isaah  
CIPANP '06



Baryons:  $n=3$   
Mesons:  $n=2$

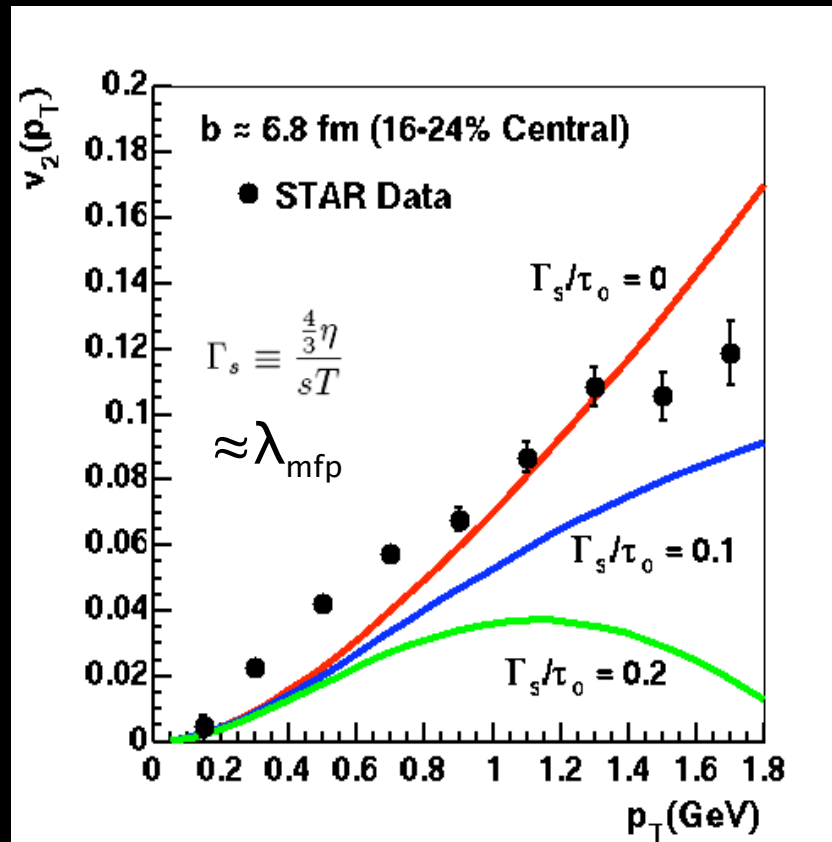
Flow mechanism “knows” about quarks

But: detailed microscopic dynamics that lead to  
“quark-number scaling” are not yet understood



# How well does our fluid flow?

D. Teaney, 2003: Estimated viscous corrections to ideal hydro calculations

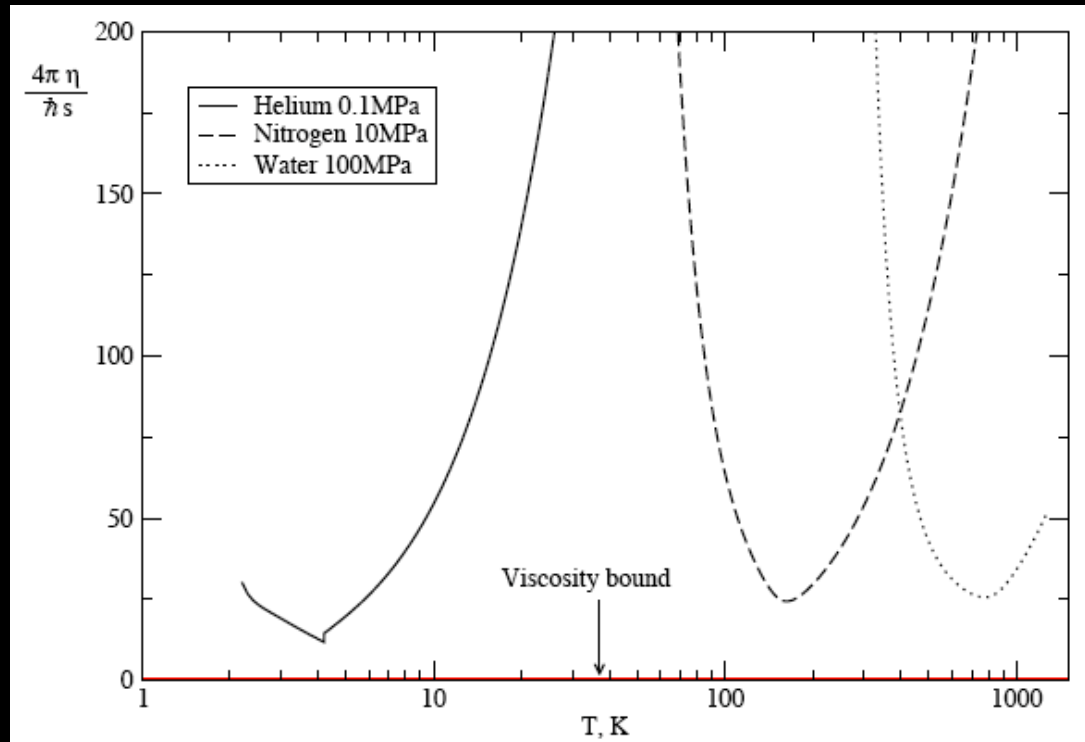


Comparing  
shear viscosity/entropy density,  
RHIC matter is **100× better**  
fluid than water

Large elliptic flow implies  
very small viscosity  $\Leftrightarrow$  small  $\lambda_{\text{mfp}}$   $\Leftrightarrow$  strong coupling

# Perfect liquid at RHIC?

Data from RHIC suggest that  $\eta/s < 0.2$  (possibly  $< 0.1$ )



← RHIC @  $4\pi \eta/s \approx 1$

How can one calculate  $\eta/s$ ?

# How can one calculate $\eta/s$ ?

Perturbative QCD gives  $\eta/s \approx 1$

Liquid  $\rightarrow$  vanishing  $\lambda_{\text{mfp}} \rightarrow$  strong coupling  
pQCD is the right theory, but wrong approximation

Lattice QCD: hard (see later)

String theory:

Shear Viscosity of Strongly Coupled  $\mathcal{N} = 4$  Supersymmetric Yang-Mills Plasma

G. Policastro<sup>1,2</sup>, D.T. Son<sup>3,4</sup>, and A.O. Starinets<sup>1</sup>

<sup>1</sup>*Department of Physics, New York University, New York, New York 10003*

<sup>2</sup>*Scuola Normale Superiore, Piazza dei Cavalieri 7, 56100, Pisa, Italy*

<sup>3</sup>*Physics Department, Columbia University, New York, New York 10027*

<sup>4</sup>*RIKEN-BNL Research Center, Brookhaven National Laboratory, Upton, New York 11973*

(April 2001)

# String Theory to the Rescue?

Using the anti-de Sitter/conformal field theory correspondence, we relate the shear viscosity  $\eta$  of the finite-temperature  $\mathcal{N} = 4$  supersymmetric Yang-Mills theory in the large  $N$ , strong-coupling regime with the absorption cross section of low-energy gravitons by a near-extremal black three-brane. We show that in the limit of zero frequency this cross section coincides with the area of the horizon. From this result we find  $\eta = \frac{\pi}{8} N^2 T^3$ . We conjecture that for finite 't Hooft coupling  $g_{\text{YM}}^2 N$  the shear viscosity is  $\eta = f(g_{\text{YM}}^2 N) N^2 T^3$ , where  $f(x)$  is a monotonic function that decreases from  $\mathcal{O}(x^{-2} \ln^{-1}(1/x))$  at small  $x$  to  $\pi/8$  when  $x \rightarrow \infty$ .

The ratio of shear viscosity to volume density of entropy can be used to characterize how close a given fluid is to being perfect. Using string theory methods, we show that this ratio is equal to a universal value of  $\hbar/4\pi k_B$  for a large class of strongly interacting quantum field theories whose dual description involves black holes in anti-de Sitter space. We provide evidence that this value may serve as a lower bound for a wide class of systems, thus suggesting that black hole horizons are dual to the most ideal fluids.

## AdS/CFT correspondence

Maldacena (1997), Gubser, Klebanov, Polyakov; Witten (1998)

$\mathcal{N} = 4$  Super-Yang-Mills theory in 4d with  $SU(N_C)$   $\longleftrightarrow$  A string theory in 5d AdS

Finite temperature  $\longleftrightarrow$  Black hole in  $AdS_5$

Large  $N_C$  and strong coupling limit  $\longleftrightarrow$  Classical gravity limit

YM observables at infinite  $N_C$  and infinite coupling can be computed using classical gravity

Apply to both dynamical and thermodynamic observables.

# Viscosity Bound

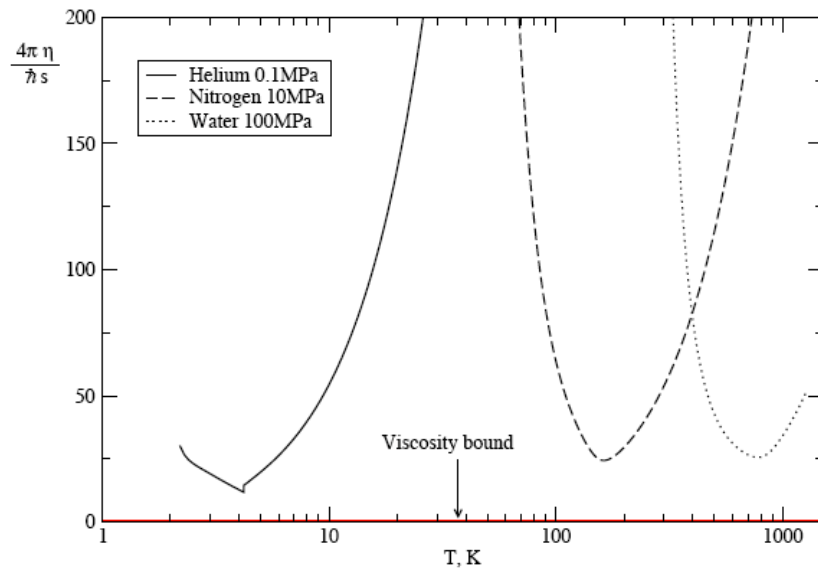


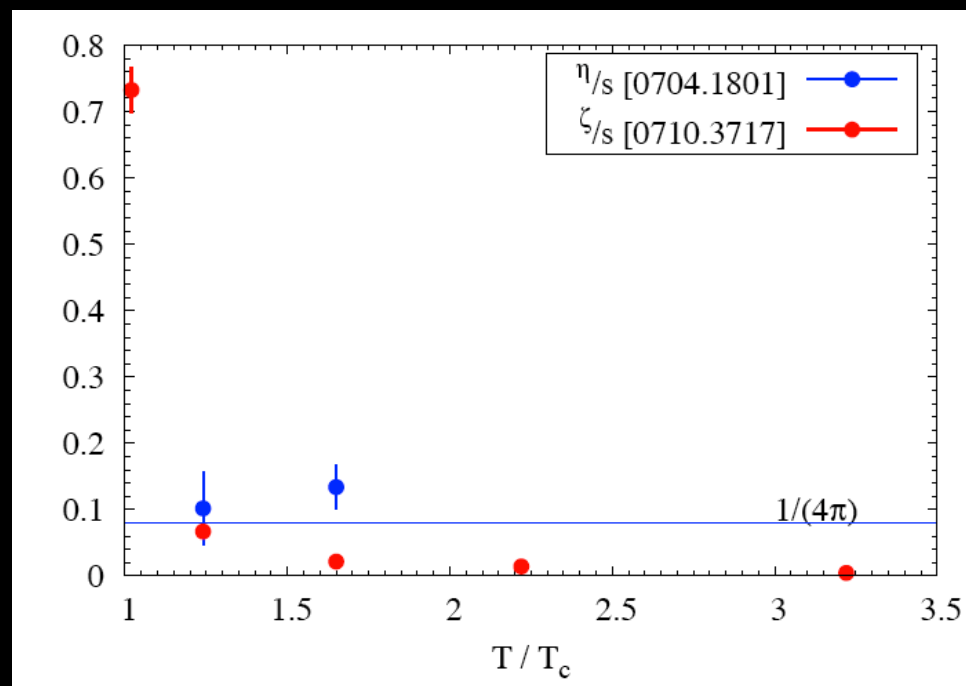
Figure 2: The viscosity-entropy ratio for some common substances: helium, nitrogen and water. The ratio is always substantially larger than its value in theories with gravity duals, represented by the horizontal line marked “viscosity bound.”

All field theories with a gravity dual were found to have  
 $\eta/s > 1/4\pi$

Universal bound?

“Black hole horizons (in 5-d anti-de-Sitter space...) are dual to the most ideal liquids”

# Lattice QCD Calculation of $\eta/s$

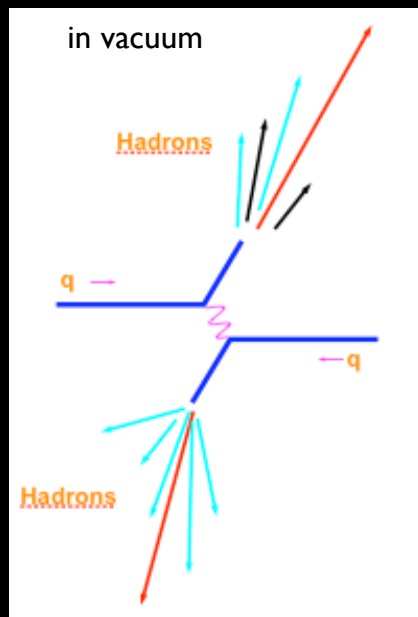


Harvey Meyer (MIT)

Lattice calculations agree with small shear viscosity for QCD plasma

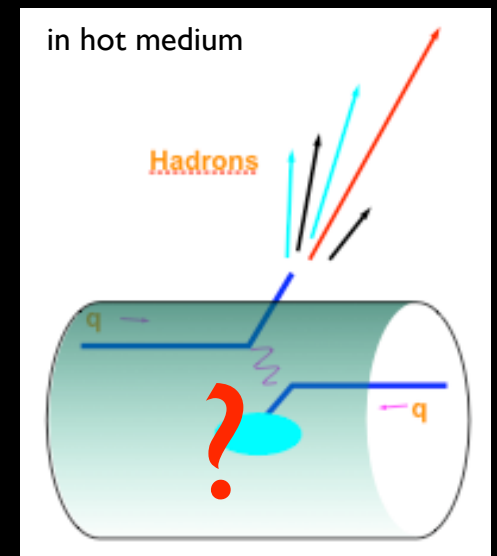
Require bold extrapolations to extract  $\eta$

# What else can one learn about the medium?



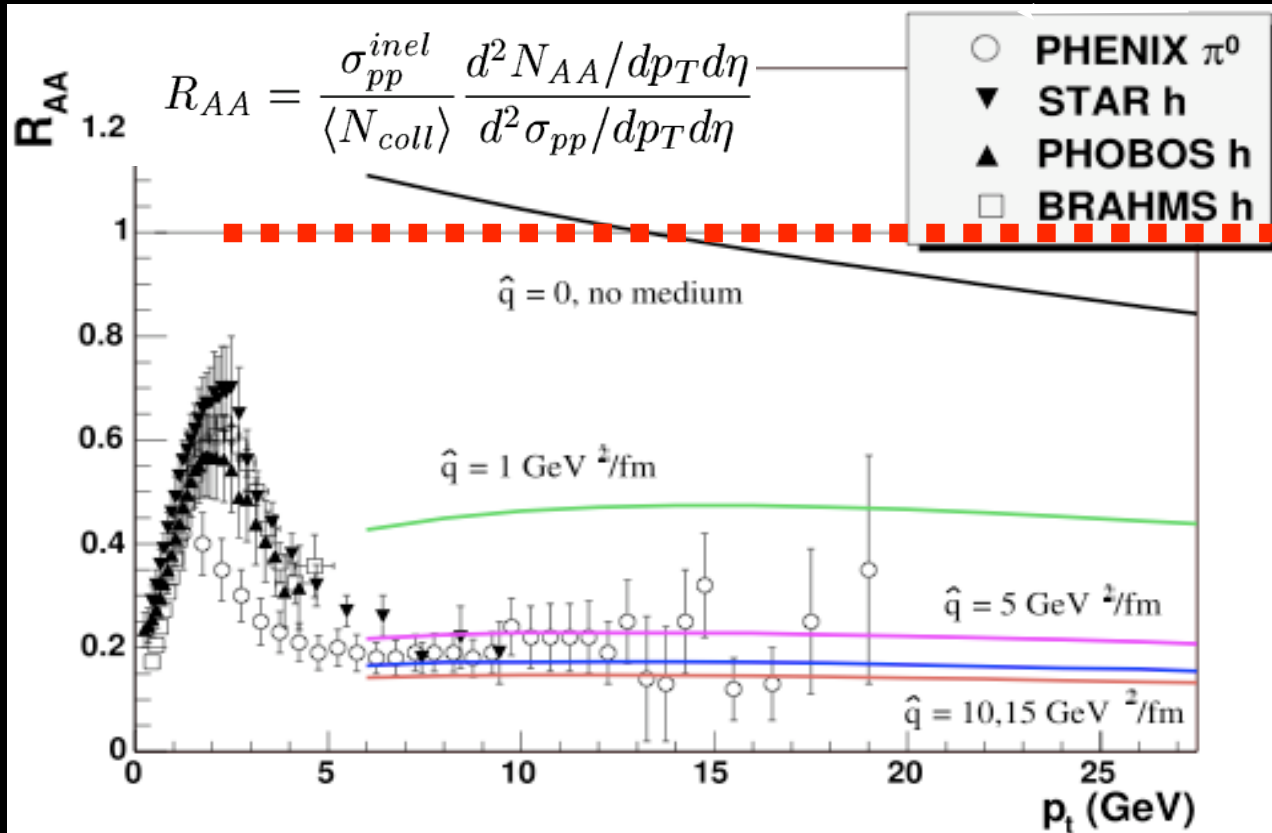
Hadrons at high  $p_T$  ( $\gg T$ )  
originate from  
“fragmentation” of high  $p_T$   
quarks (or gluons)

What happens when high  
 $p_T$  partons traverse the  
medium?





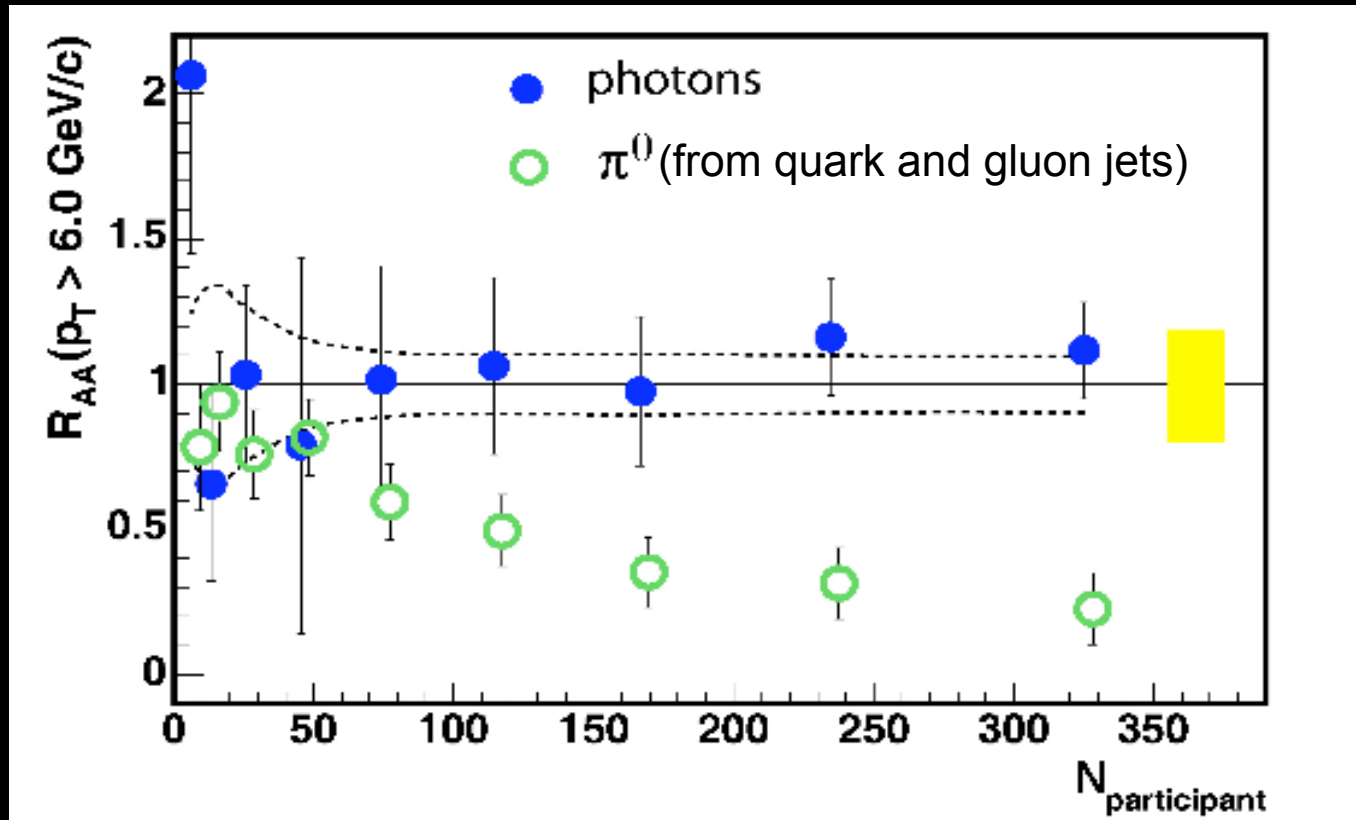
# The Medium is “black”: Jet Quenching



Expected yield in Au+Au,  
relative to p+p

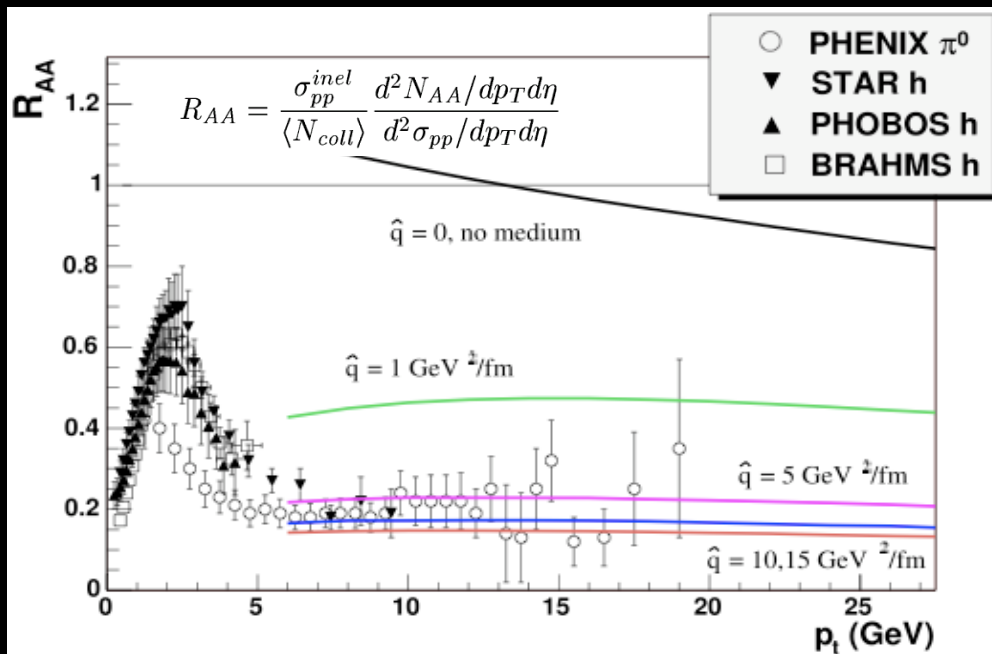
Observe a suppression  
("jet quenching")  
by factor 5-6!

# $N_{\text{coll}}$ -Scaling ?



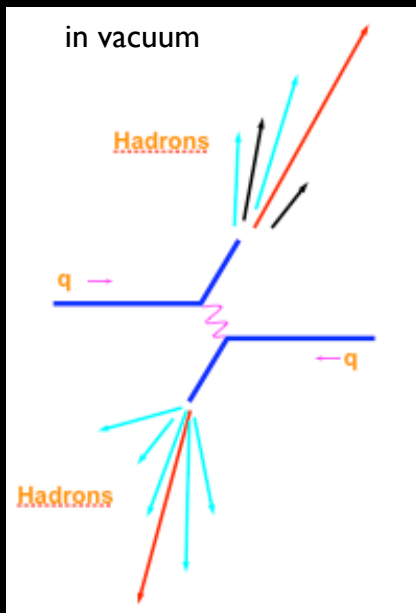
High  $p_T$  photons (which don't suffer energy loss in the medium) are produced with the expected rate relative to p+p

# The Medium is “black”: Jet Quenching



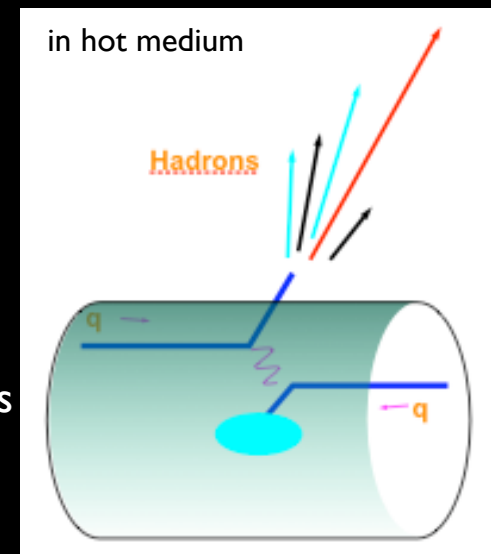
Expected yield in Au+Au,  
relative to p+p

Observe a suppression  
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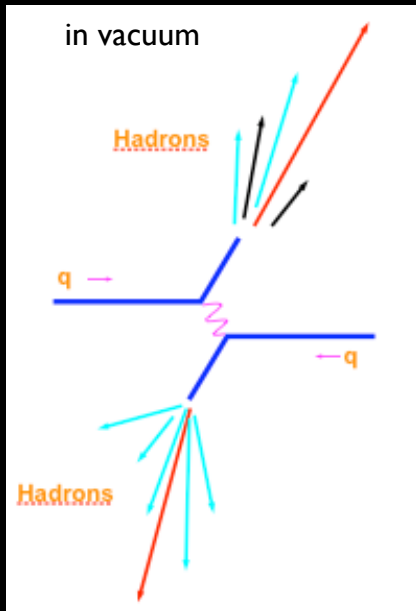
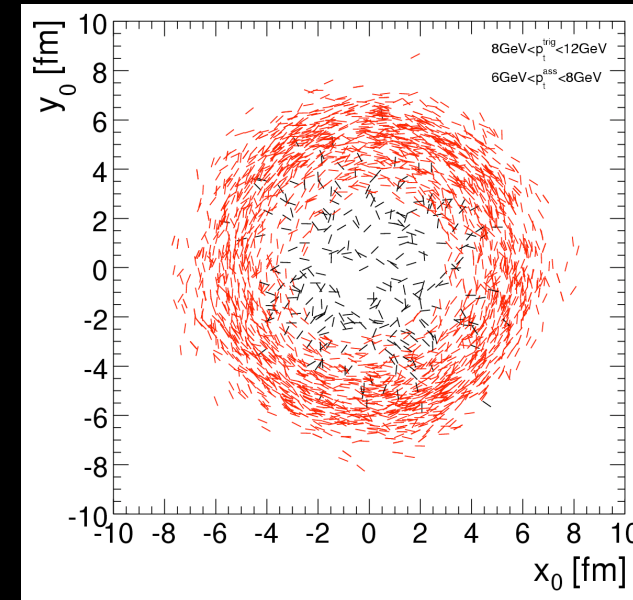
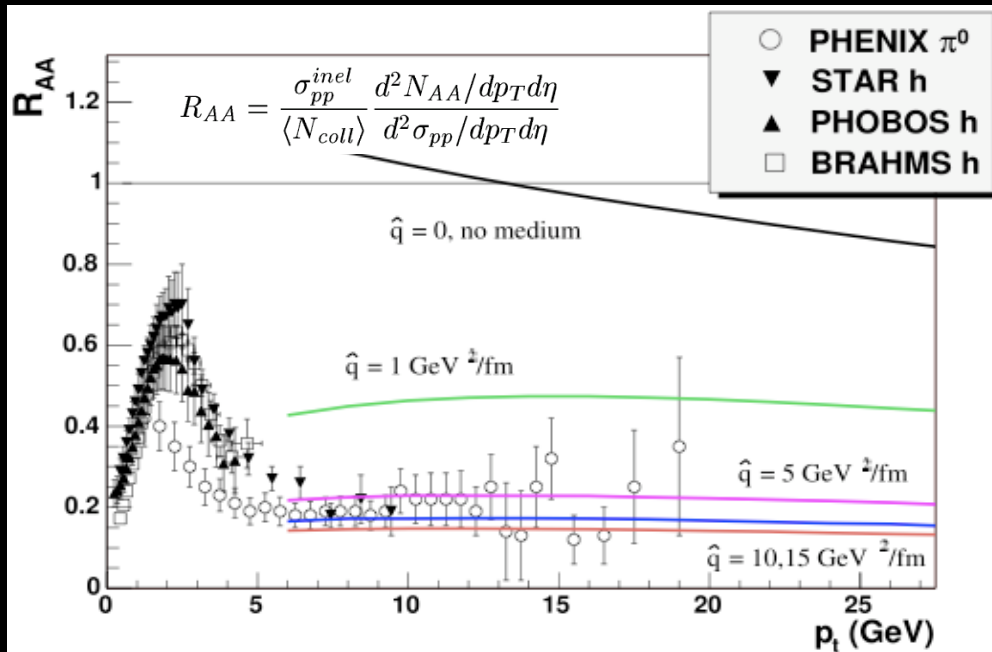


Hadrons at high  $p_T$   
originate from  
“fragmentation” of high  $p_T$   
quarks (or gluons)

In medium, only “surface  
radiation” escapes. Partons  
traversing medium are  
“swallowed” by medium.

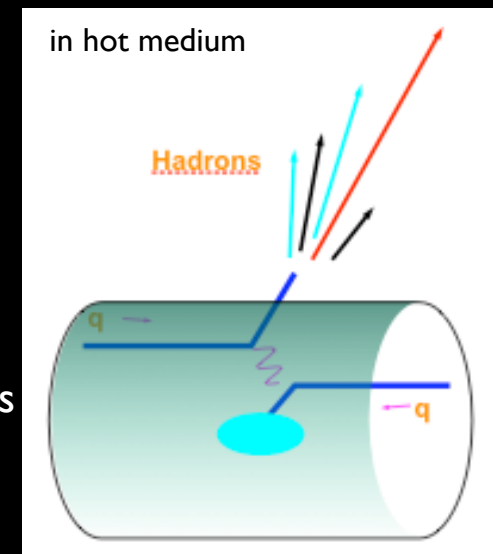


# The Medium is “black”: Jet Quenching

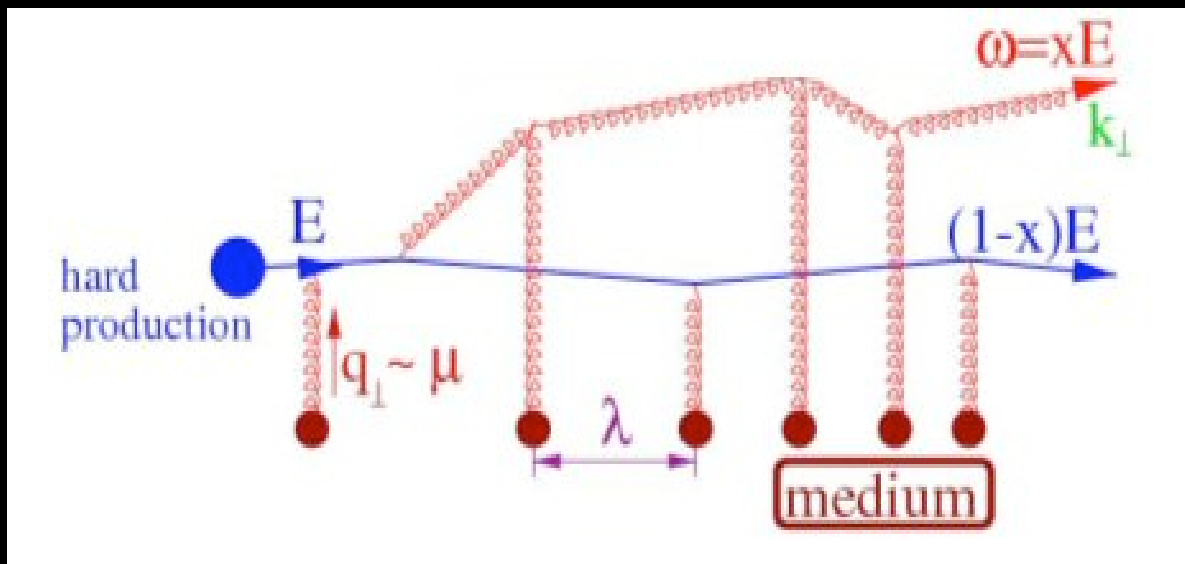


Hadrons at high  $p_T$  originate from “fragmentation” of high  $p_T$  quarks (or gluons)

In medium, only “surface radiation” escapes. Partons traversing medium are “swallowed” by medium.



# Gluon Bremsstrahlung



$$\Delta E \approx -\frac{\alpha_s}{2\pi} N_c \hat{q} L^2$$

$k_T$  kick per unit path length

pQCD calculations give  $\hat{q} \approx 1-3 \text{ GeV}^2/\text{fm}$

# Energy Loss in String Theory

## $\hat{q}$ of $\mathcal{N}=4$ SYM theory

BDMPS transport coefficient reads:  $\lambda = g_{YM}^2 N_c$

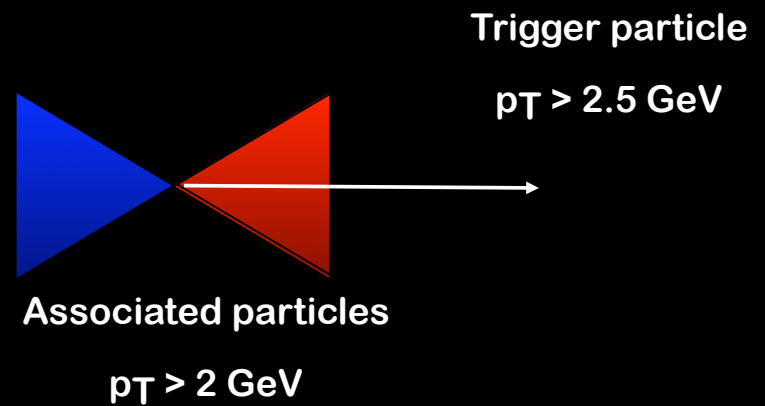
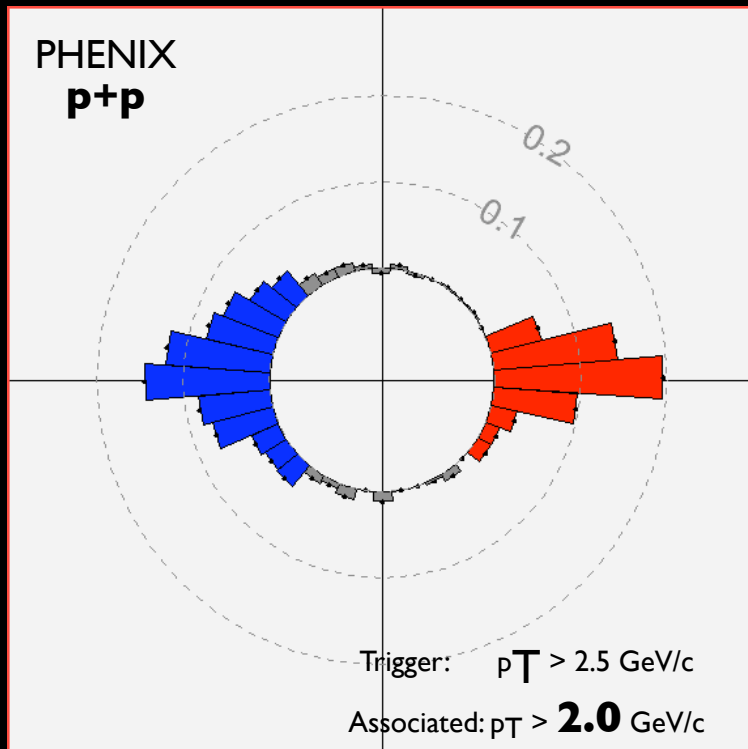
$$\hat{q}_{SYM} = \frac{\pi^{3/2} \Gamma(\frac{3}{4})}{\Gamma(\frac{5}{4})} \sqrt{\lambda} T^3 \approx 26.69 \sqrt{\alpha_{SYM} N_c} T^3$$

- It is **not** proportional to number of scattering centers
- Take:  $N_C = 3, \alpha_s = \frac{1}{2}, T = 300 \text{ MeV}$

$$\hat{q}_{SYM} = 4.5 \text{ GeV}^2/\text{fm}.$$

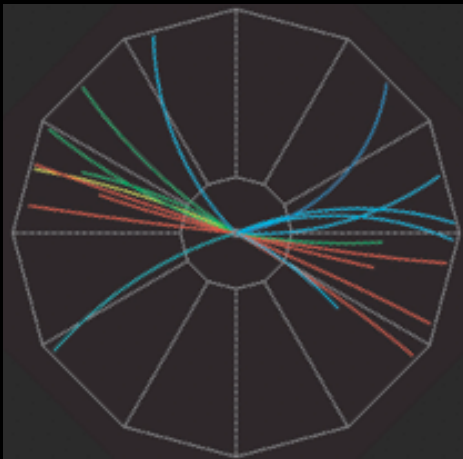
- Experimental estimates: 5-15 GeV<sup>2</sup>/fm

# Jets and Angular Correlations

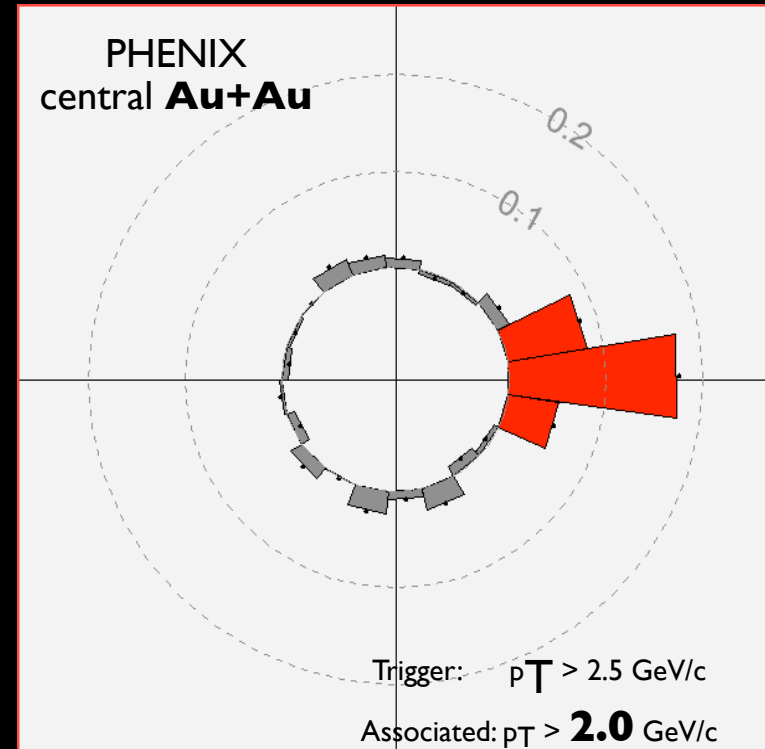
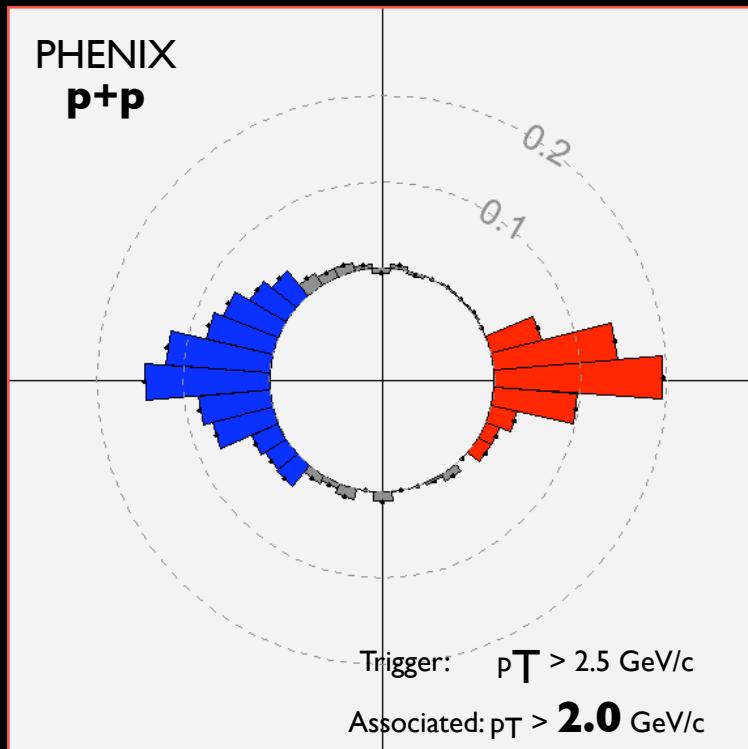


Plot angle of associated particles above  
 $p_T$  threshold relative to trigger

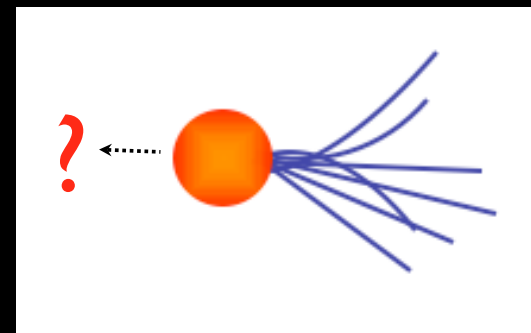
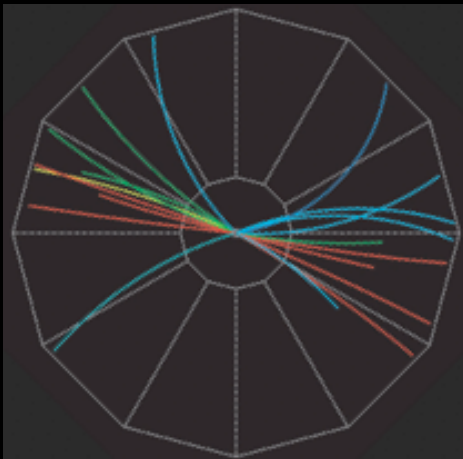
STAR  
200 GeV p+p  
 $p_T > 2.0$  GeV/c



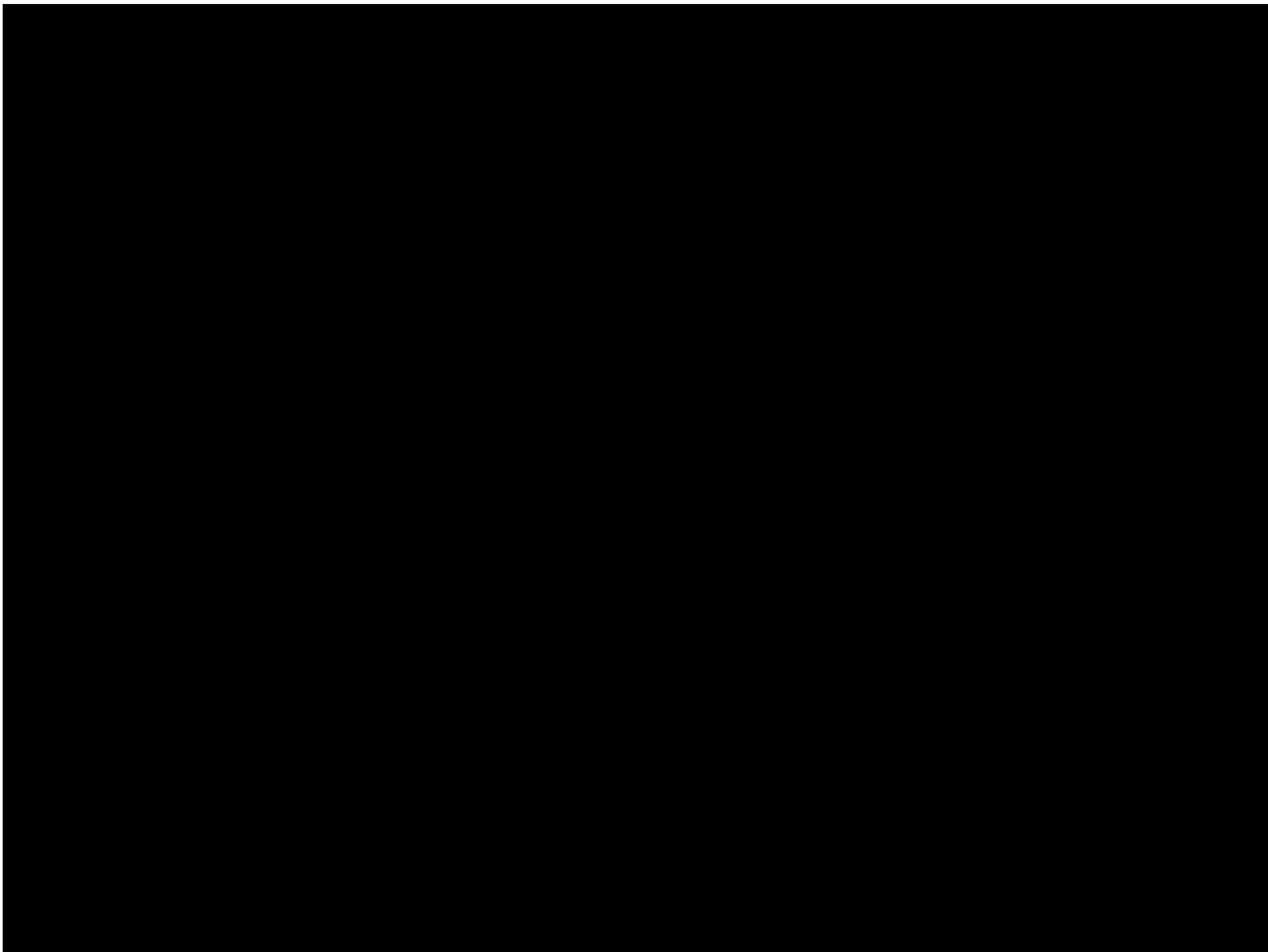
# Jets and Angular Correlations



STAR  
200 GeV p+p  
 $p_T > 2.0$  GeV/c



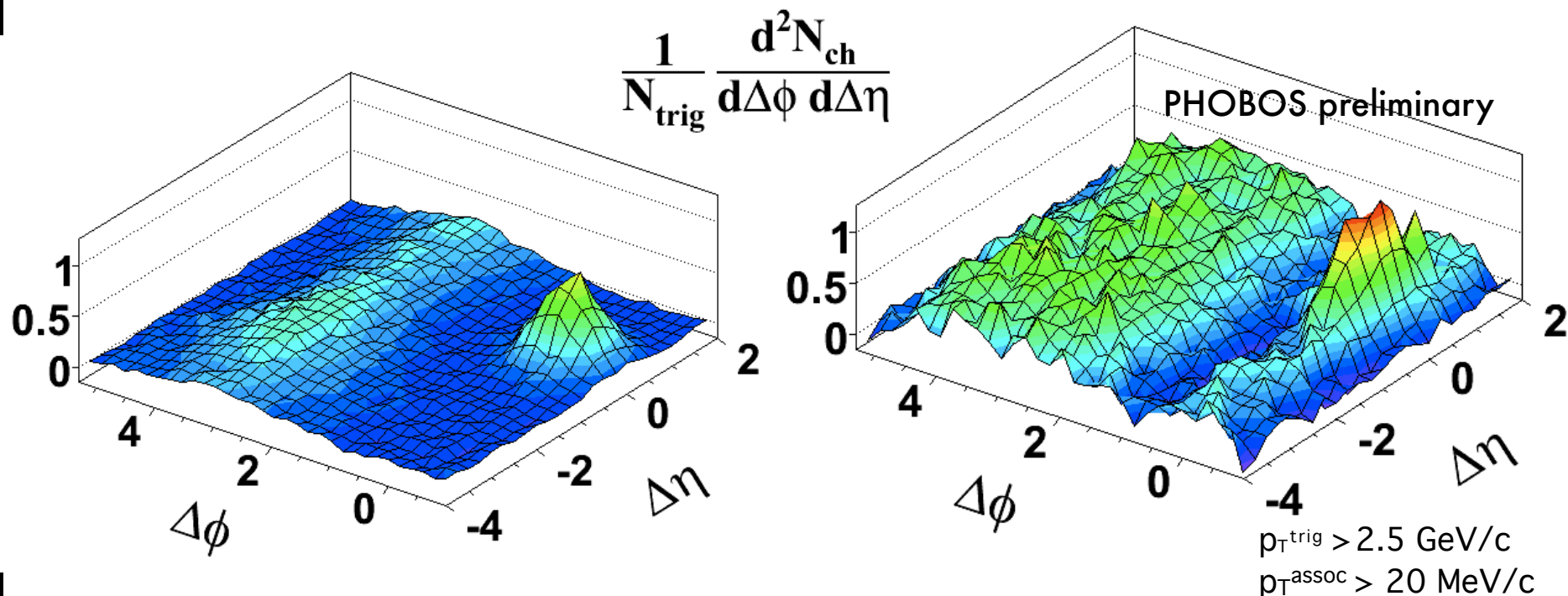




# Where does the energy go?

p+p PYTHIA v6.325

Au+Au 0-30% central



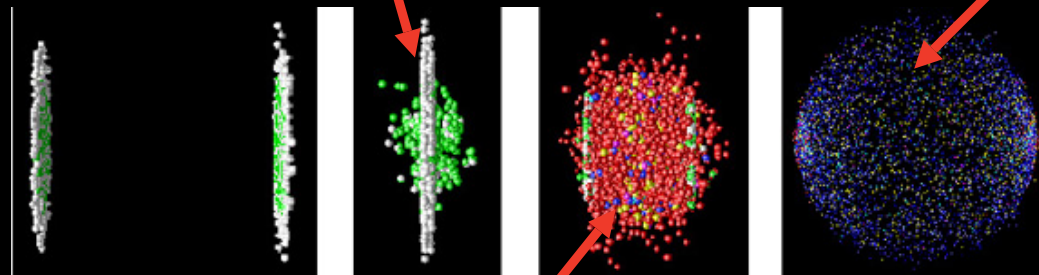
“Lost” energy found in quasi-thermal low  $p_{\text{T}}$  particles,  
even far in rapidity from trigger particle

### Initial Collisions

Hard Scattering takes place [direct  $\gamma$ ]  
High  $p_T$  partons are produced [d+Au]  
Overall Entropy defined [ $dN/d\eta$ ]  
Geometrical asymmetry [Geometry]

### Hadronization

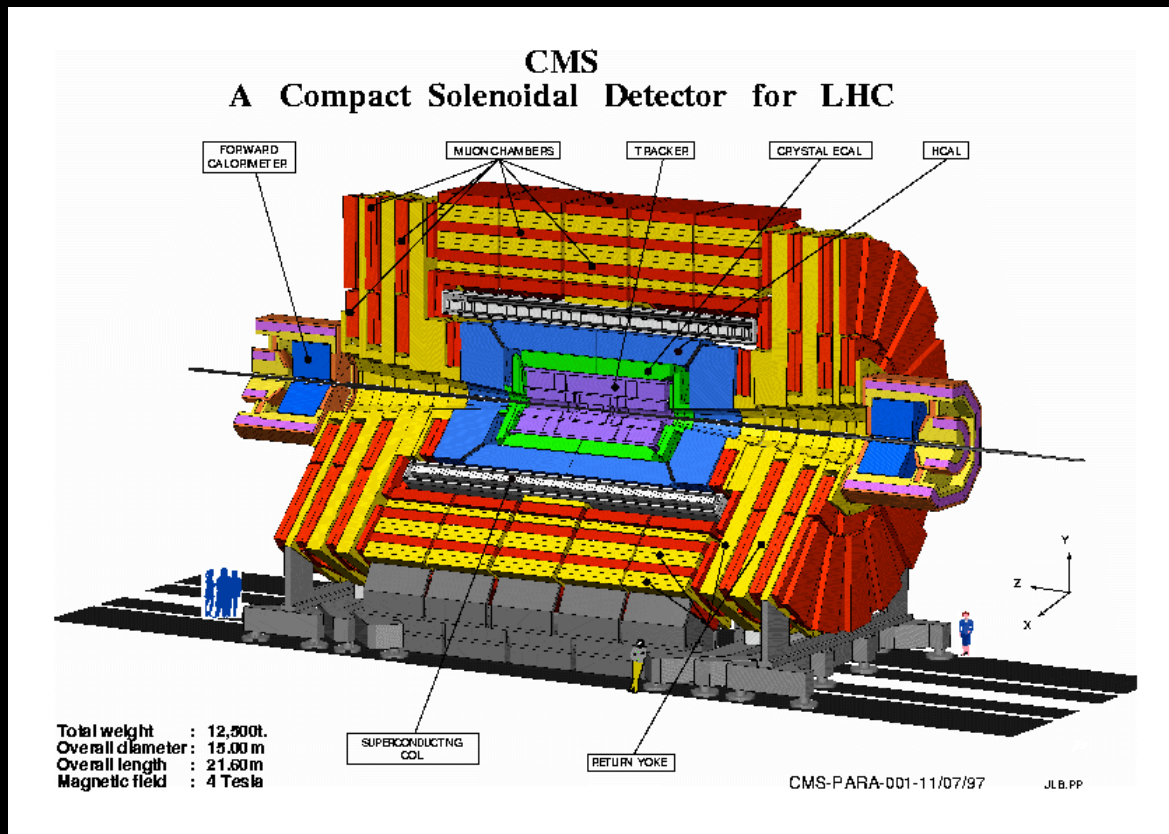
Recombination from quark soup [proton-non suppression, quark-scaling of  $v_2$ ]  
Global statistical hadron formation at  $T_{ch} = 170$  MeV [particle ratios]  
Radial expansion with  $\beta_T \sim 0.6c$  [PID spectra]  
Particle emission after 10fm/c for few fm/c [HBT]



### Early Stage ( $\sim$ few fm/c)

High Density ( $\sim 5$  GeV/fm<sup>3</sup>) [ $dN/d\eta$ , high  $p_T$  suppression]  
Local thermal equilibration [Elliptic Flow  $v_2$ ]  
Pressure driven expansion [Elliptic Flow  $v_2$ , HBT]  
Opaque for fast partons [Back-to-Back jets]

# Heavy Ions at LHC



## CMS: Big experiment

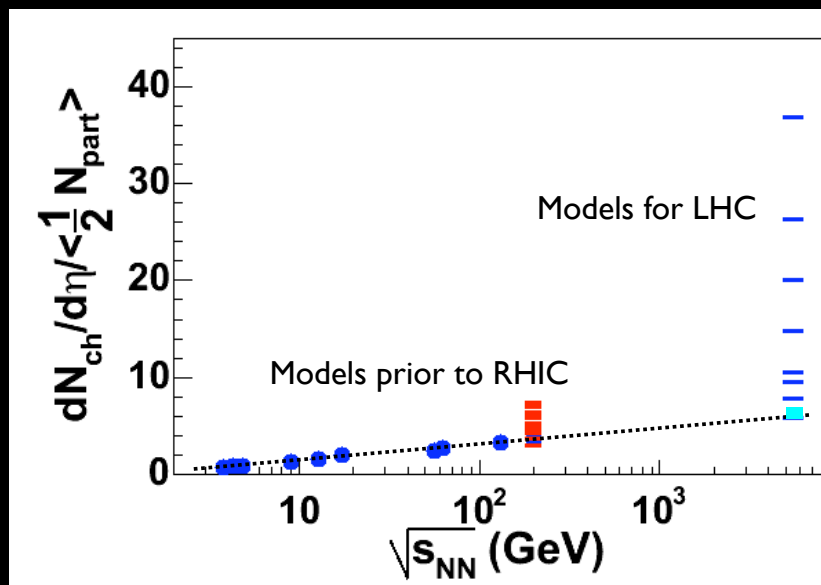
2600 Physicists  
\$500M construction  
Designed for p+p

12500 tons  
1 GHz interaction rate  
1 TByte/sec data flow  
World's largest magnet (2.6 GJ)  
200 m<sup>2</sup> Si Detectors

LHC Tunnel will close Sep 1 2007

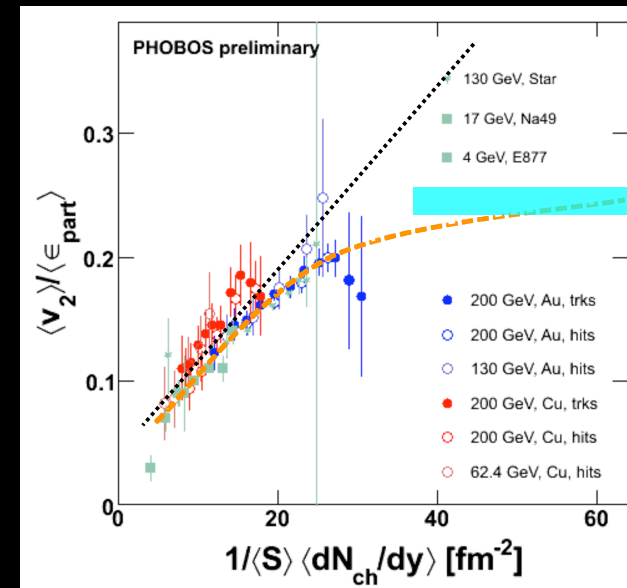
## Heavy ions at LHC

Unprecedented change in initial conditions  
Qualitatively new probes of the medium, e.g. Jets



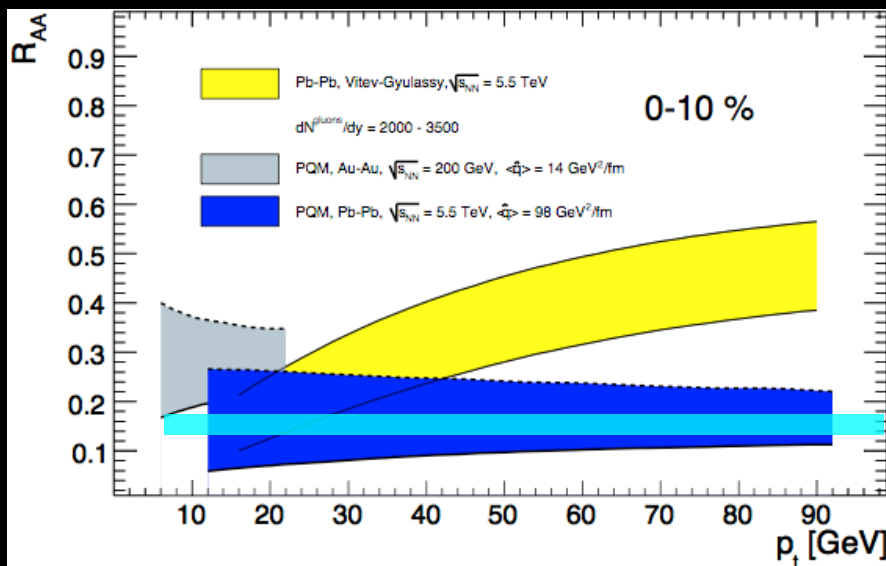
LHC

I Day: Multiplicity  $\Rightarrow$  Initial Density



LHC

I Week: Does elliptic flow saturate, indicating equilibrium?

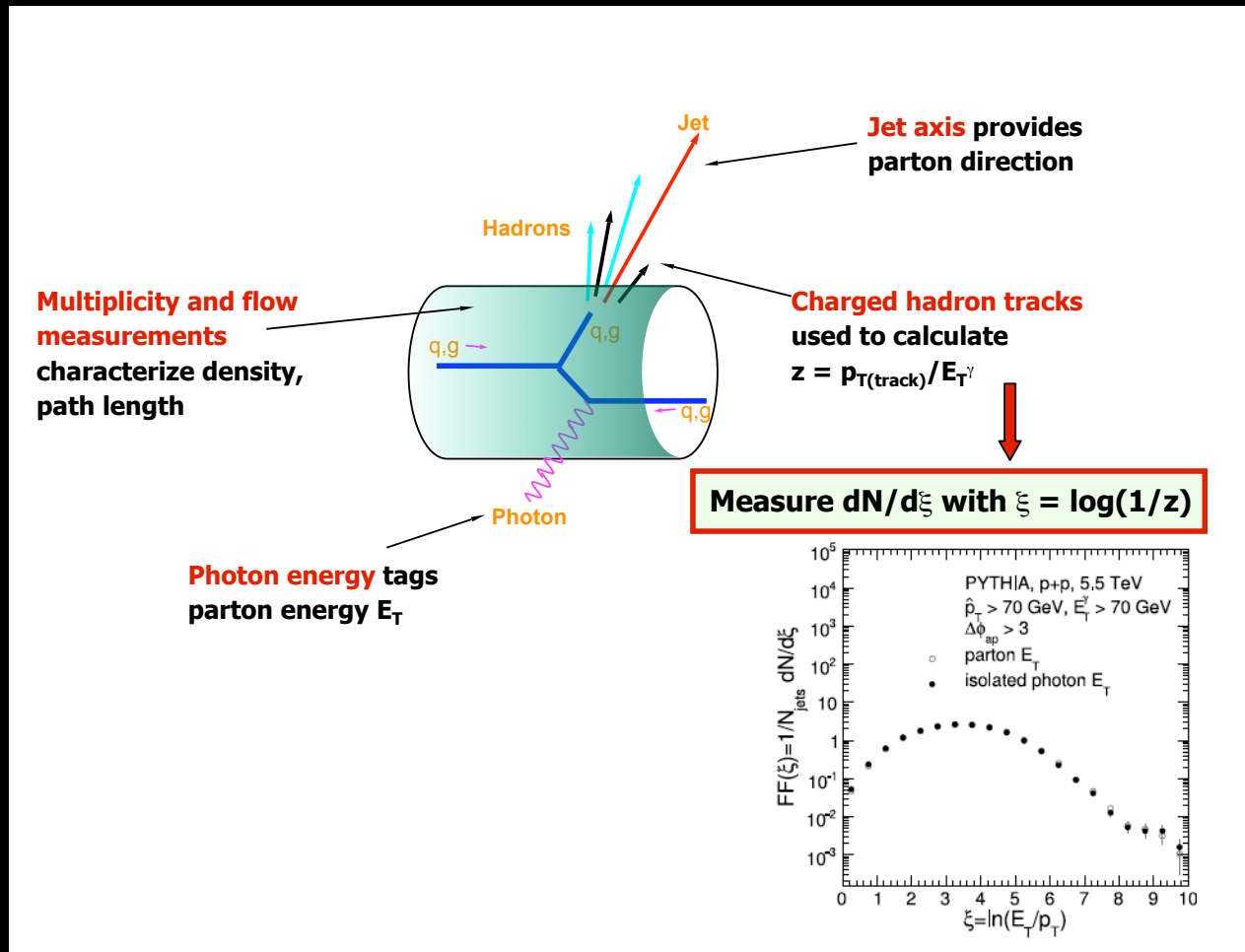


LHC

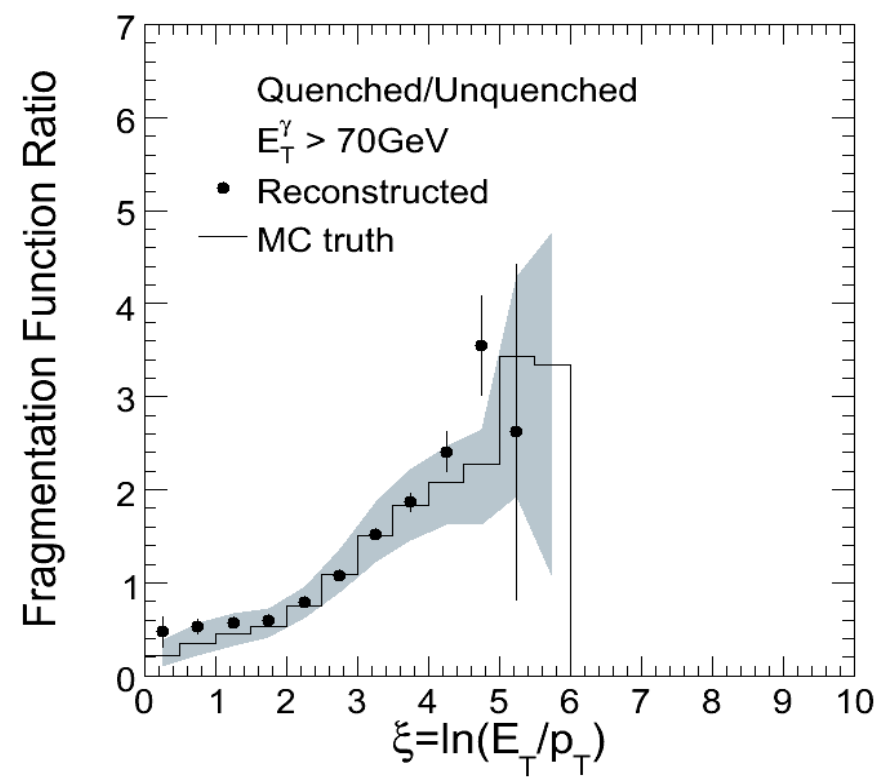
I Month: What is the jet quenching parameter? Is the medium “black”?

Once we have these qualitative answers:  
 $\rightarrow$  program of precision measurements of medium properties

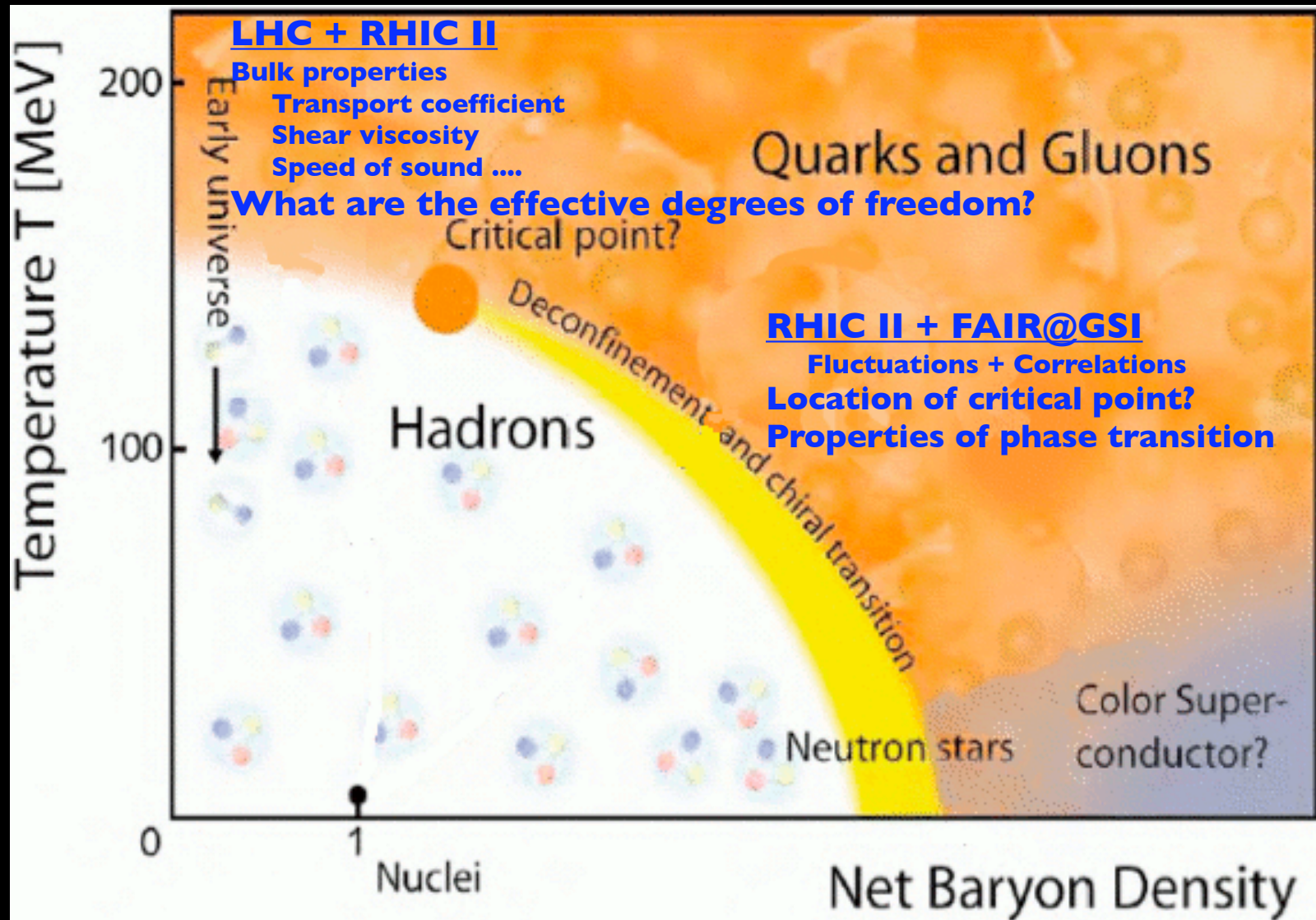
# Measuring energy loss at LHC



Higher collision energy allows to find jets and photons at high  $p_T$  ( $>70 \text{ GeV}/c$ )



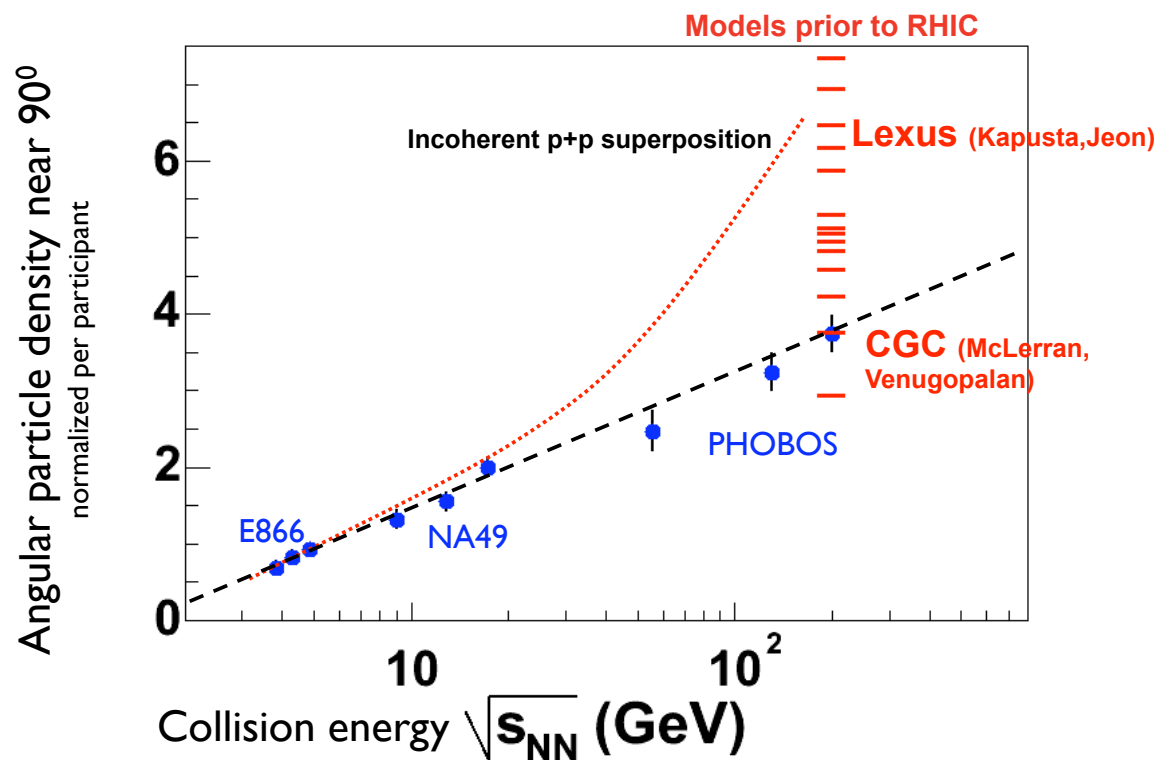




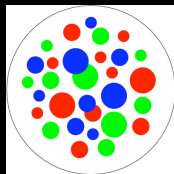




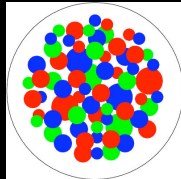
See you  
next week!



*Low Energy*



*High Energy*



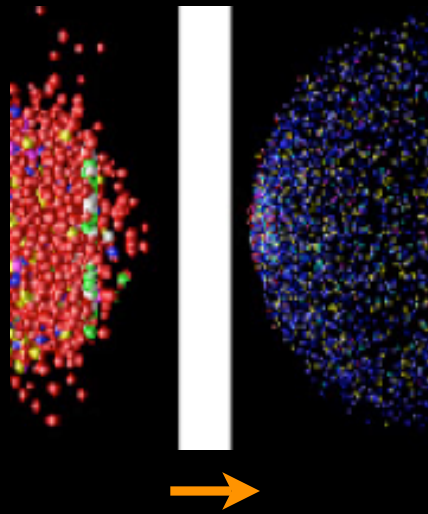
**Parton Saturation**

Particle production by “liberation” of  
gluons already present in incoming nuclei

Effective gluon density increases with energy,  
but saturates when gluons below  $Q_{\text{sat}}$   
overlap in transverse plane

This “Color Glass Condensate” describes  
nuclei at high energies

# Hadronization



# Enhancement of Multi-Strange Baryons

受入  
98-06-144  
高工研圖書室

EUROPEAN ORGANIZATION FOR NUCLEAR RESEARCH

CERN-EP/98-64  
22 April 1998

## Enhancement of central $\Lambda$ , $\Xi$ and $\Omega$ yields in Pb-Pb collisions at 158 A GeV/c

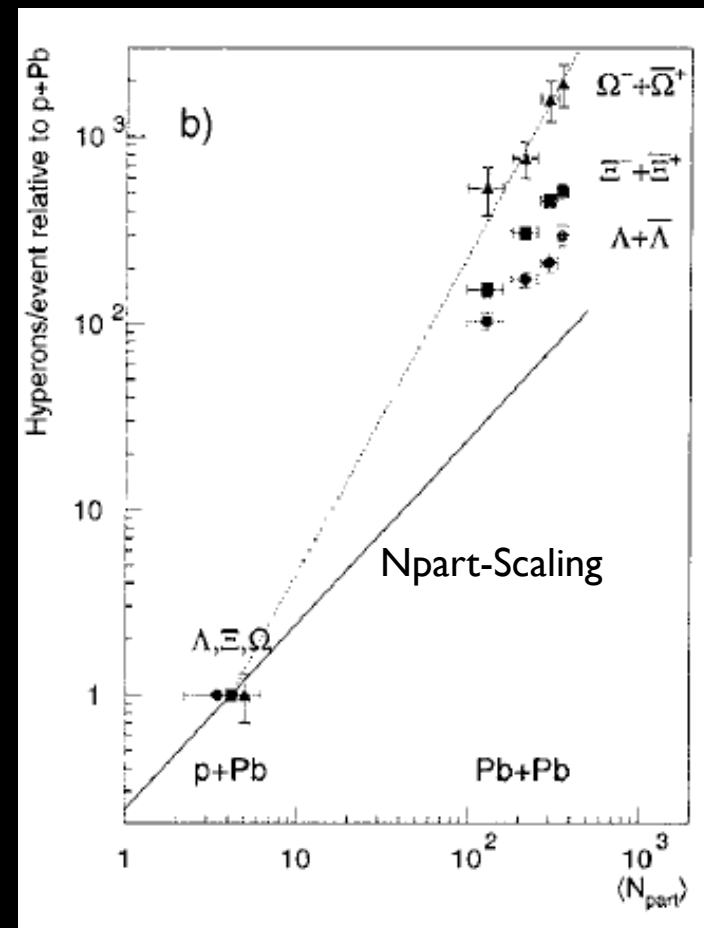
E. Andersen<sup>1)</sup>, F. Antinori<sup>6,12)</sup>, N. Armenise<sup>2)</sup>, H. Bakke<sup>3)</sup>, J. Bán<sup>8)</sup>, D. Barberis<sup>7)</sup>, H. Beker<sup>6)</sup>, W. Beusch<sup>6)</sup>, L.J. Bloodworth<sup>3)</sup>, J. Böhm<sup>14)</sup>, R. Calandro<sup>2)</sup>, M. Campbell<sup>6)</sup>, E. Cantatore<sup>6)</sup>, N. Carrer<sup>12)</sup>, M.G. Catanesi<sup>2)</sup>, E. Ches<sup>6)</sup>, M. Dameri<sup>7)</sup>, G. Darbo<sup>7)</sup>, A. Diczek<sup>13)</sup>, D. Di Bari<sup>2)</sup>, S. Di Liberto<sup>13)</sup>, B.C. Earl<sup>6)</sup>, D. Elia<sup>2)</sup>, D. Evans<sup>4)</sup>, K. Fanebust<sup>3)</sup>, R.A. Fini<sup>2)</sup>, J.C. Fontaine<sup>10)</sup>, J. Ftáčnik<sup>3)</sup>, B. Ghidini<sup>2)</sup>, G. Grella<sup>16)</sup>, M. Guida<sup>16)</sup>, E.H.M. Heijne<sup>6)</sup>, H. Helstrup<sup>4)</sup>, A.K. Holme<sup>6)</sup>, D. Huss<sup>20)</sup>, A. Jacholkowski<sup>2)</sup>, G.T. Jones<sup>5)</sup>, P. Jovanovic<sup>5)</sup>, A. Jusko<sup>8)</sup>, V.A. Kachanov<sup>17)</sup>, T. Kachelhoffer<sup>18)</sup>, J.B. Kinson<sup>5)</sup>, A. Kirk<sup>5)</sup>, W. Klempt<sup>6)</sup>, B.T.H. Knudsen<sup>3)</sup>, K. Knudson<sup>6)</sup>, I. Králik<sup>4)</sup>, J.C. Lassalle<sup>6)</sup>, V. Lenti<sup>2)</sup>, R. Lietava<sup>8)</sup>, R.A. Loconsole<sup>2)</sup>, G. Løvhaugen<sup>6,11)</sup>, M. Lupták<sup>6)</sup>, V. Mack<sup>10)</sup>, V. Manzari<sup>2)</sup>, P. Martinengo<sup>6)</sup>, M.A. Mazzoni<sup>15)</sup>, F. Meddi<sup>13)</sup>, A. Michalon<sup>18)</sup>, M.E. Michalon-Mentzer<sup>18)</sup>, P. Middelkamp<sup>6)</sup>, M. Morando<sup>12)</sup>, M.T. Muciaccia<sup>2)</sup>, E. Nappi<sup>2)</sup>, F. Navach<sup>2)</sup>, P.I. Norman<sup>6)</sup>, B. Osculati<sup>7)</sup>, B. Pastirčák<sup>3)</sup>, F. Pellegrini<sup>12)</sup>, K. Piška<sup>14)</sup>, F. Posa<sup>2)</sup>, E. Quercigh<sup>6)</sup>, R.A. Ricci<sup>7)</sup>, G. Romano<sup>16)</sup>, G. Rosa<sup>16)</sup>, L. Rossi<sup>7)</sup>, H. Rotscheidt<sup>6)</sup>, K. Šafařík<sup>6)</sup>, S. Saladino<sup>2)</sup>, C. Salvo<sup>7)</sup>, L. Šándor<sup>4,8)</sup>, T. Scognetti<sup>2)</sup>, G. Segato<sup>12)</sup>, M. Sené<sup>13)</sup>, R. Sené<sup>13)</sup>, S. Simone<sup>2)</sup>, A. Singovski<sup>17)</sup>, W. Snoeys<sup>3)</sup>, P. Staroba<sup>14)</sup>, S. Szafran<sup>13)</sup>, M. Thompson<sup>5)</sup>, T.F. Thorsteinsen<sup>3)</sup>, G. Tomasichio<sup>2)</sup>, G.D. Torrieri<sup>6)</sup>, T.S. Tveter<sup>11)</sup>, J. Urbán<sup>8)</sup>, G. Vassiliadis<sup>11)</sup>, M. Venables<sup>5)</sup>, O. Villalobos Baillie<sup>6)</sup>, T. Virgili<sup>16)</sup>, A. Volte<sup>13)</sup>, M.F. Votruba<sup>6)</sup> and P. Závada<sup>14)</sup>

### Abstract

$\Lambda$ ,  $\Xi$  and  $\Omega$  yields and transverse mass spectra have been measured at central rapidity in Pb-Pb and p-Pb collisions at 158 A GeV/c. The yields in Pb-Pb interactions are presented as a function of the collision centrality and compared with those obtained from p-Pb collisions. Strangeness enhancement is observed which increases with centrality and with the strangeness content of the hyperon.

To be submitted to *Physics Letters B*

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- <sup>18)</sup> Centre de Recherches Nucléaires, Strasbourg, France



Large Enhancement for e.g.  
 $\Omega$ /participant, relative to p  
+Pb

# Spectrum of Produced Hadrons

$$\langle n_j \rangle = \frac{(2J_j + 1)V}{(2\pi)^3} \int d^3p \left[ e^{\sqrt{p^2 + m_j^2}/T + \mu \cdot \mathbf{q}_j/T} \pm 1 \right]^{-1}$$

Diagram illustrating the variables in the equation for the spectrum of produced hadrons:

- Yield** (indicated by a red arrow pointing to  $\langle n_j \rangle$ )
- Temperature** (indicated by a cyan arrow pointing to  $T$ )
- Mass** (indicated by a red arrow pointing to  $m_j$ )
- Chemical Potential** (indicated by a cyan arrow pointing to  $\mu$ )
- Quantum Numbers** (indicated by a red arrow pointing to  $J_j$  and  $\mathbf{q}_j$ )

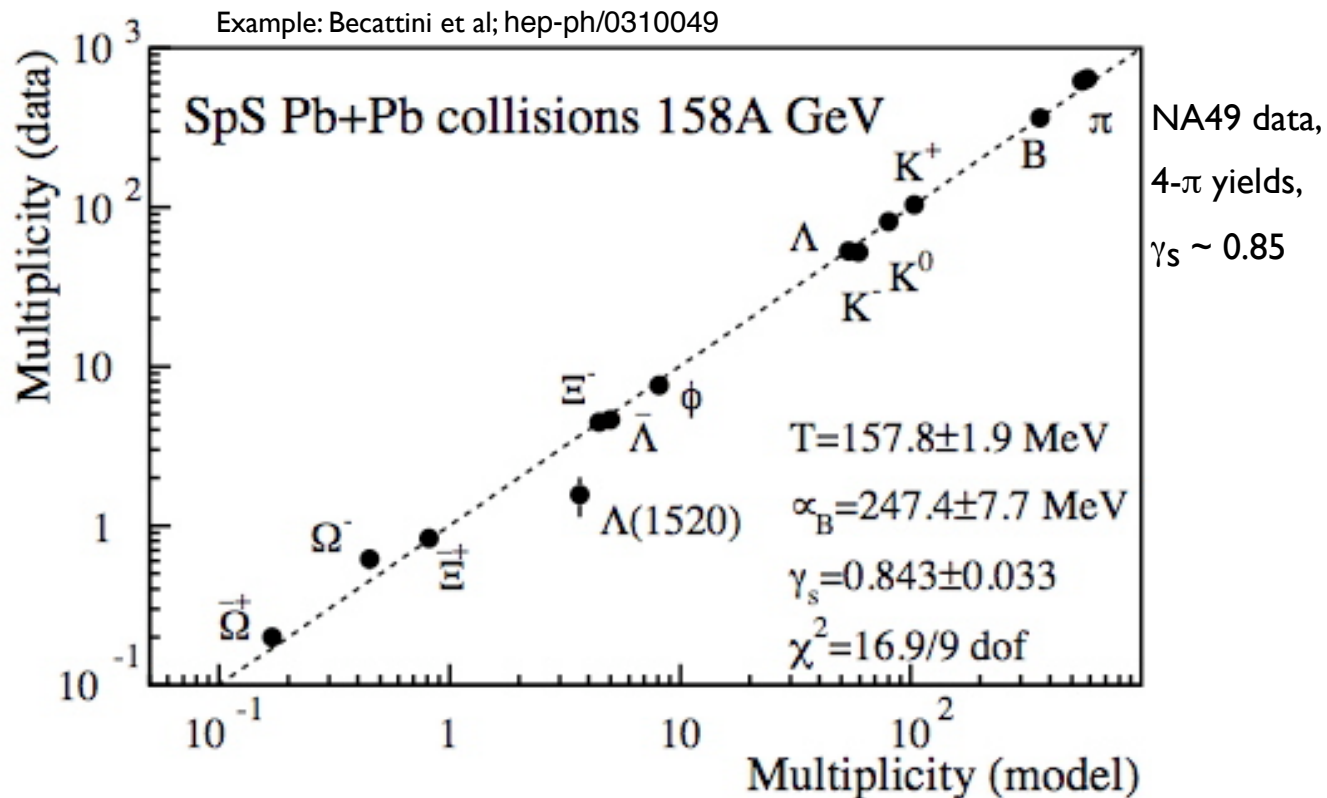
- Statistical *Description* of Observed Yields in Gibbs Grand-Canonical Ensemble

- Many Different Implementations

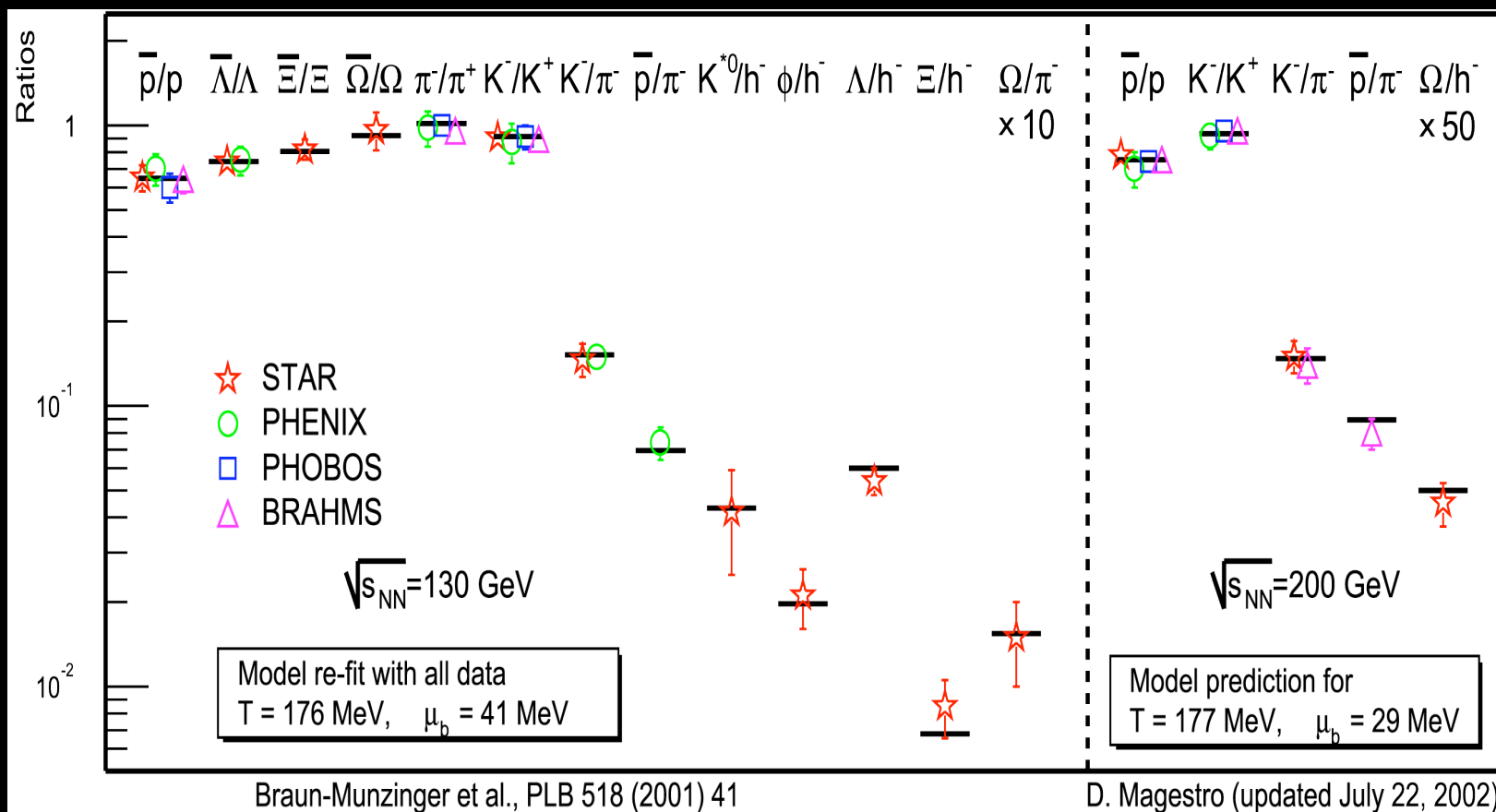
c.f. Hagedorn, Becattini, Braun-Munzinger, Cleymans, Heinz, Letessier, Mekijan, Rafelski, Redlich, Satz, Sollfrank, Stachel, Tounsi + many others

- Mid-Rapidity vs  $4-\pi$  yields
- Non-Equilibrium ( $\gamma_s, \gamma_q$ )
- Numerical Implementation
- Here: Common Features of Different Approaches

# Spectrum of Produced Hadrons

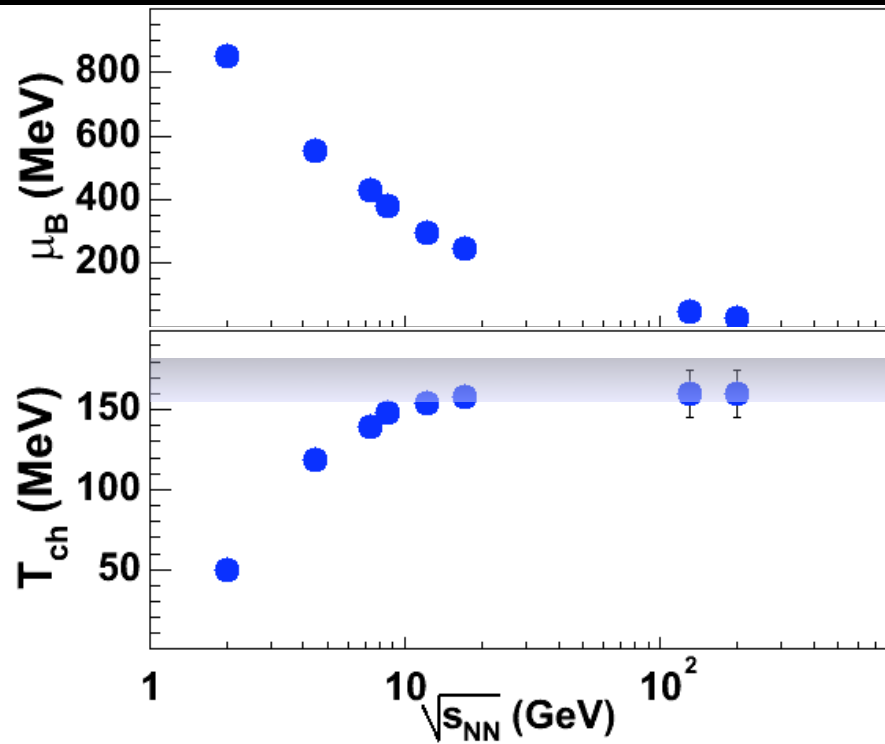


# Statistical Model Fit



**Relative Abundance: Two Parameters !**

# “Thermal Fit” Parameters vs sqrt(s)



$\mu_B$  drops with collision energy

$T_{ch}$  approaches limiting value



Size (Mass, Volume)

Microcanonical ensemble. All conservation laws including energy-momentum (angular momentum, parity), charges enforced.

$V > 20 \text{ fm}^3, M > 10 \text{ GeV}$  (F. Liu et al., Phys. Rev. C 68 (2003) 024905)  
F. B., L. Ferroni, talk in ISMD 2003)

Canonical ensemble. Energy and momentum conserved on average, charges exactly. Temperature is introduced

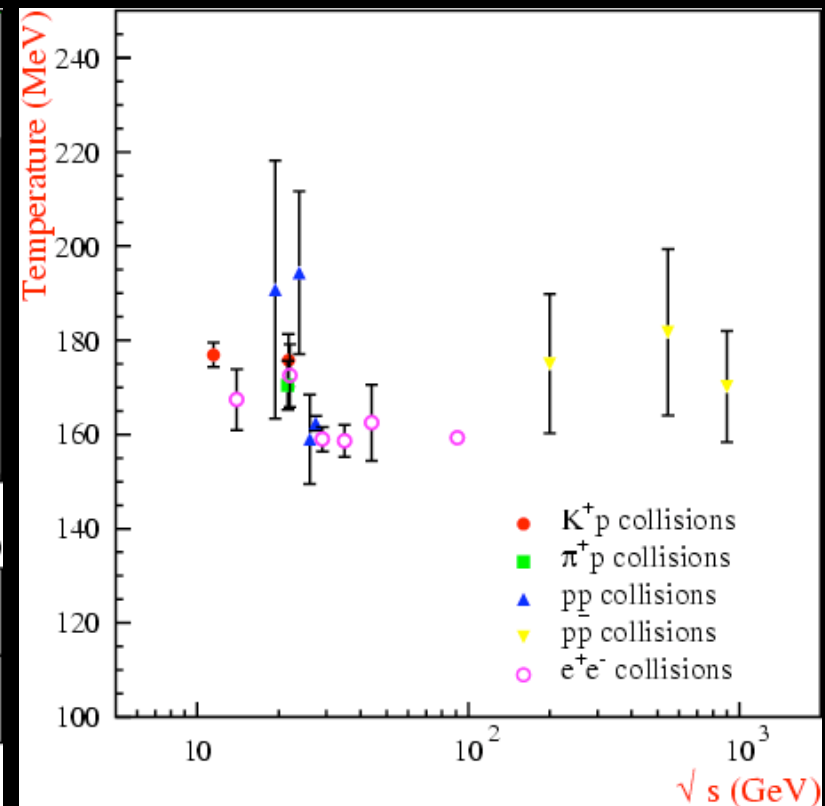
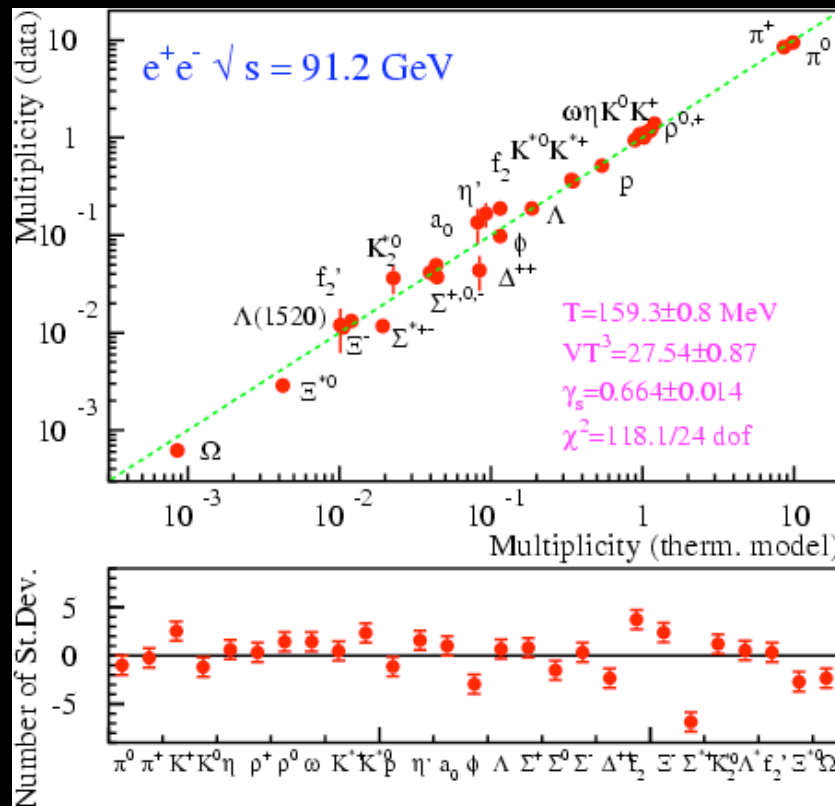
$V > 100 \text{ fm}^3, M > 50 \text{ GeV}$  (A. Keranen, F.B., Phys. Rev. C 65 (2002) 044901)

Grand-canonical ensemble. Also charges are conserved on average. Chemical potentials are introduced

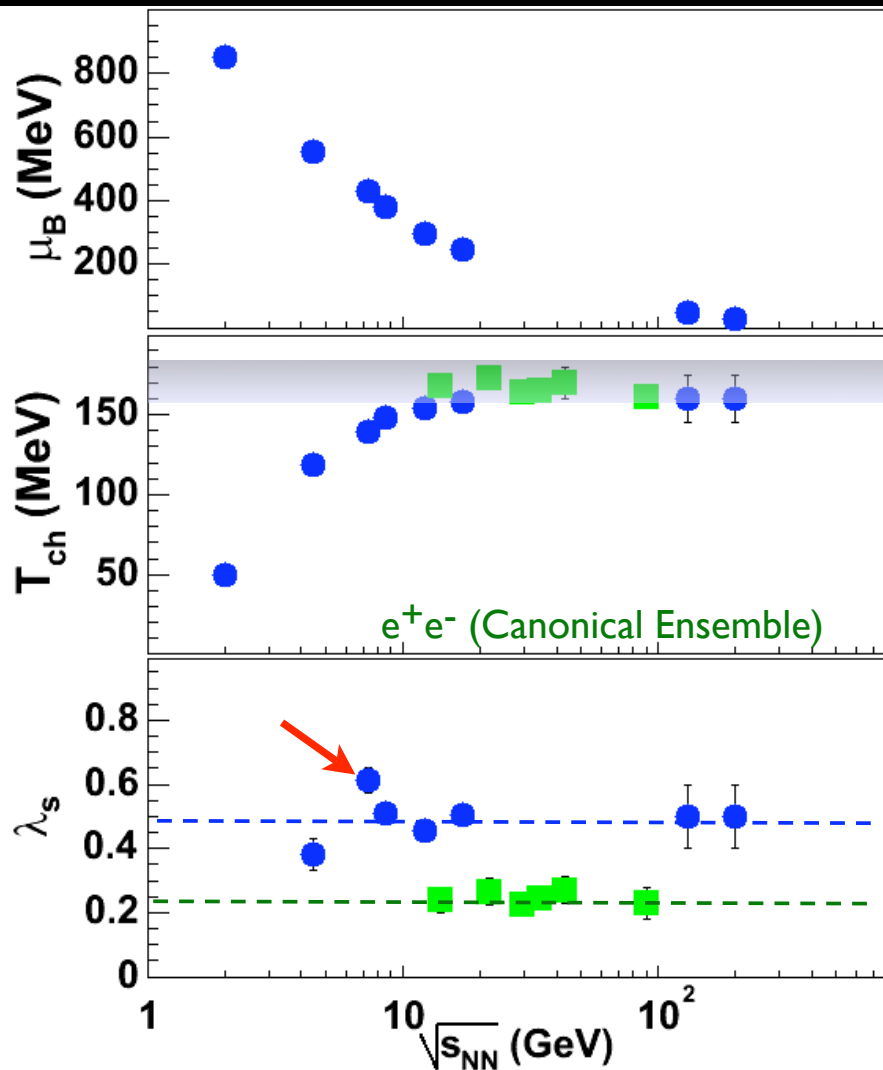
Difficulty of computing

from Francesco Becattini

# Statistical Model for Elementary Collisions



# “Thermal Fit” Parameters vs sqrt(s)



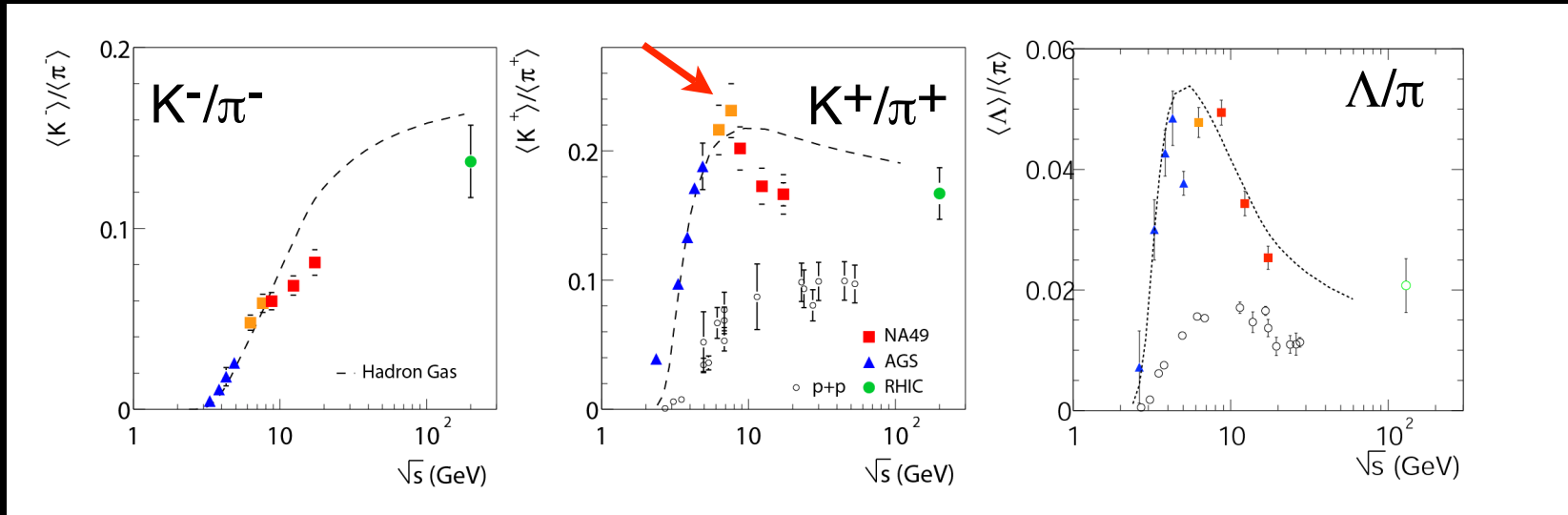
e<sup>+</sup>e<sup>-</sup> hadronizes at same  $T_{ch}$

Are we looking at a *local*  
or *global* property?

Strangeness enhancement unique to AA

Global (or at least large)  
correlation volume

# The “Horn”



NA49: Sharp maximum in  $K/\pi$  ratio at low SPS energy

Broad maximum from  $\mu_B$  and p+p  $\sqrt{s}$  dependence,  
but no sharp structure in models