

Leading Order SCET Feynman Rules

1 Leading Order Lagrangian

The leading order Lagrangian, with a single set of quark and gluon collinear modes in the n direction and quark and gluon ultrasoft modes, in general covariant gauge, is,

$$\mathcal{L}^{(0)} = \mathcal{L}_{us}^{(0)} + \sum_n \mathcal{L}_{n\xi}^{(0)} + \mathcal{L}_{ng}^{(0)},$$

$$\mathcal{L}_{n\xi}^{(0)} = e^{-ix \cdot \mathcal{P}} \bar{\xi}_n \left(in \cdot D + i \not{D}_{n\perp} \frac{1}{i\bar{n} \cdot D_n} i \not{D}_{n\perp} \right) \frac{\not{n}}{2} \xi_n,$$

$$\mathcal{L}_{ng}^{(0)} = \frac{1}{2g^2} \text{Tr} \{ [i\mathcal{D}^\mu, i\mathcal{D}^\nu] [i\mathcal{D}_\mu, i\mathcal{D}_\nu] \} + \tau \text{Tr} \{ ([i\mathcal{D}_{us}^\mu, A_{n\mu}]^2) \} + 2\text{Tr} \{ \bar{c}_n [i\mathcal{D}_\mu^{us}, [i\mathcal{D}^\mu, c_n]] \},$$

$$\mathcal{L}_{us}^{(0)} = \bar{\psi}_{us} i \not{D}_{us} \psi_{us} - \frac{1}{2} \text{Tr} \{ G_{us}^{\mu\nu} G_{\mu\nu}^{us} \} + \tau_{us} \text{Tr} \{ (i\partial_\mu A_{us}^\mu)^2 \} + 2\text{Tr} \{ \bar{c}_{us} i\partial_\mu iD_{us}^\mu c_{us} \},$$

$$iD_{n\perp}^\mu = \mathcal{P}_\perp^\mu + gA_{n\perp}^\mu, \quad i\bar{n} \cdot D_n = \bar{\mathcal{P}} + g\bar{n} \cdot A_n,$$

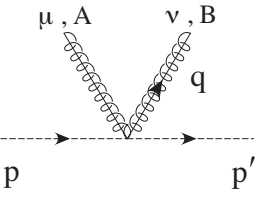
$$i\mathcal{D}^\mu = \frac{n^\mu}{2} (\bar{\mathcal{P}} + g\bar{n} \cdot A_n) + (\mathcal{P}_\perp^\mu + gA_{\perp,n}^\mu) + \frac{\bar{n}}{2} (in \cdot \partial + gn \cdot A_n + gn \cdot A_{us}),$$

$$i\mathcal{D}_{us}^\mu \equiv \frac{n^\mu}{2} \bar{\mathcal{P}} + \mathcal{P}_\perp^\mu + \frac{\bar{n}^\mu}{2} in \cdot \partial + \frac{\bar{n}}{2} gn \cdot A_{us}, \quad iD_{us}^\mu = i\partial_\mu + A_{us}^\mu$$

2 Leading Order Feynman Rules

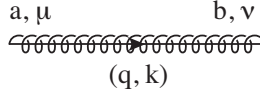
The above Lagrangian leads to the Feynman rules

$$\begin{array}{l}
 \begin{array}{c} (p, p_r) \\ \text{-----} \blacktriangleright \text{-----} \end{array} \\
 \begin{array}{c} \mu, A \\ \text{-----} \blacktriangleright \text{-----} \\ \text{~~~~~} \updownarrow \text{~~~~~} \end{array} \\
 \begin{array}{c} \mu, A \\ \text{-----} \blacktriangleright \text{-----} \blacktriangleright \text{-----} \\ \text{~~~~~} \updownarrow \text{~~~~~} \\ p \qquad \qquad p' \end{array}
 \end{array}
 = \begin{array}{l}
 i \frac{\not{n}}{2} \frac{\bar{n} \cdot p}{n \cdot p_r \bar{n} \cdot p + p_\perp^2 + i0} \\
 ig T^A n_\mu \frac{\not{n}}{2} \\
 ig T^A \left[n_\mu + \frac{\gamma_\mu^\perp \not{p}_\perp}{\bar{n} \cdot p} + \frac{\not{p}'_\perp \gamma_\mu^\perp}{\bar{n} \cdot p'} - \frac{\not{p}'_\perp \not{p}_\perp}{\bar{n} \cdot p \bar{n} \cdot p'} \bar{n}_\mu \right] \frac{\not{n}}{2}
 \end{array}$$

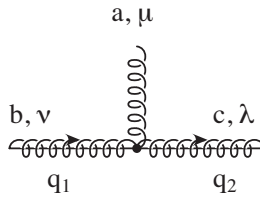


$$= \frac{ig^2 T^A T^B}{\bar{n} \cdot (p-q)} \left[\gamma_\mu^\perp \gamma_\nu^\perp - \frac{\gamma_\mu^\perp \not{p}_\perp}{\bar{n} \cdot p} \bar{n}_\nu - \frac{\not{p}'_\perp \gamma_\nu^\perp}{\bar{n} \cdot p'} \bar{n}_\mu + \frac{\not{p}'_\perp \not{p}_\perp}{\bar{n} \cdot p \bar{n} \cdot p'} \bar{n}_\mu \bar{n}_\nu \right] \frac{\bar{\gamma}}{2}$$

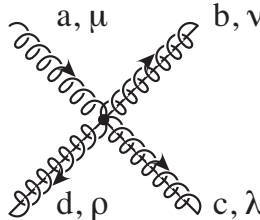
$$+ \frac{ig^2 T^B T^A}{\bar{n} \cdot (q+p')} \left[\gamma_\nu^\perp \gamma_\mu^\perp - \frac{\gamma_\nu^\perp \not{p}_\perp}{\bar{n} \cdot p} \bar{n}_\mu - \frac{\not{p}'_\perp \gamma_\mu^\perp}{\bar{n} \cdot p'} \bar{n}_\nu + \frac{\not{p}'_\perp \not{p}_\perp}{\bar{n} \cdot p \bar{n} \cdot p'} \bar{n}_\mu \bar{n}_\nu \right] \frac{\bar{\gamma}}{2}$$



$$= \frac{-i}{\bar{n} \cdot q \bar{n} \cdot k + q_\perp^2 + i0} \left(g_{\mu\nu} - (1-\tau) \frac{q_\mu q_\nu}{\bar{n} \cdot q \bar{n} \cdot k + q_\perp^2} \right) \delta_{a,b}$$

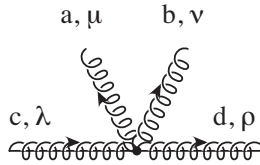


$$= g f^{abc} n_\mu \left\{ \bar{n} \cdot q_1 g_{\nu\lambda} - \frac{1}{2} \left(1 - \frac{1}{\tau_{us}} \right) [\bar{n}_\lambda q_{1\nu} + \bar{n}_\nu q_{2\lambda}] \right\}$$



$$= -\frac{1}{2} ig^2 n_\mu \left\{ f^{abe} f^{cde} (\bar{n}_\lambda g_{\nu\rho} - \bar{n}_\rho g_{\nu\lambda}) \right.$$

$$\left. + f^{ade} f^{bce} (\bar{n}_\nu g_{\lambda\rho} - \bar{n}_\lambda g_{\nu\rho}) + f^{ace} f^{bde} (\bar{n}_\nu g_{\lambda\rho} - \bar{n}_\rho g_{\nu\lambda}) \right\}$$



$$= \frac{1}{4} ig^2 n_\mu n_\nu \bar{n}_\rho \bar{n}_\lambda \left(1 - \frac{1}{\tau_{us}} \right) \left\{ f^{ace} f^{bde} + f^{ade} f^{bce} \right\}$$

Where collinear quarks are denoted by dashed lines, ultrasoft gluons are springs, collinear gluons are springs with a line and τ is the covariant collinear gauge fixing parameter and τ_{us} is the similar ultrasoft gauge fixing parameter.

The Feynman rules with just ultrasoft fields are the same as those in QCD, as are the Feynman rules for the interactions between n -collinear gluons. There are no Lagrangian interactions at this order between collinear fields in different sectors.