SU(3) gauge theory with 12 flavours in a twisted box

C.-J. David Lin

National Chiao-Tung University, Hsinchu, Taiwan

MIT

14/02/2014

Collaborators

- Kenji Ogawa (NCTU, Taiwan)
- Hiroshi Ohki (KMI, Nagoya U., Japan)
- Alberto Ramos (DESY Zeuthen, Germany)
- Eigo Shintani (U. of Mainz, Germany)

Outline

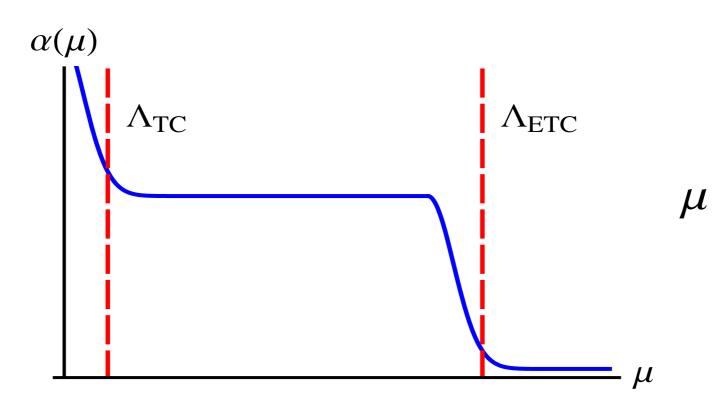
- Motivation.
- Step scaling.
- Two schemes on the twisted box:
 - ★ Twisted Polyakov Loop (TPL) scheme.
 - ★ Wilson flow (WF) scheme.
- Numerical (preliminary) results.
- Outlook.

Why SU(3) with many flavours

- Infrared conformality as an interesting fieldtheoretic problem.
- Walking technicolour model building.
- Understanding of the relation/distinction between confinement and chiral symmetry-breaking scales in QCD.

 $\alpha(\mu)$

Walking technicolour

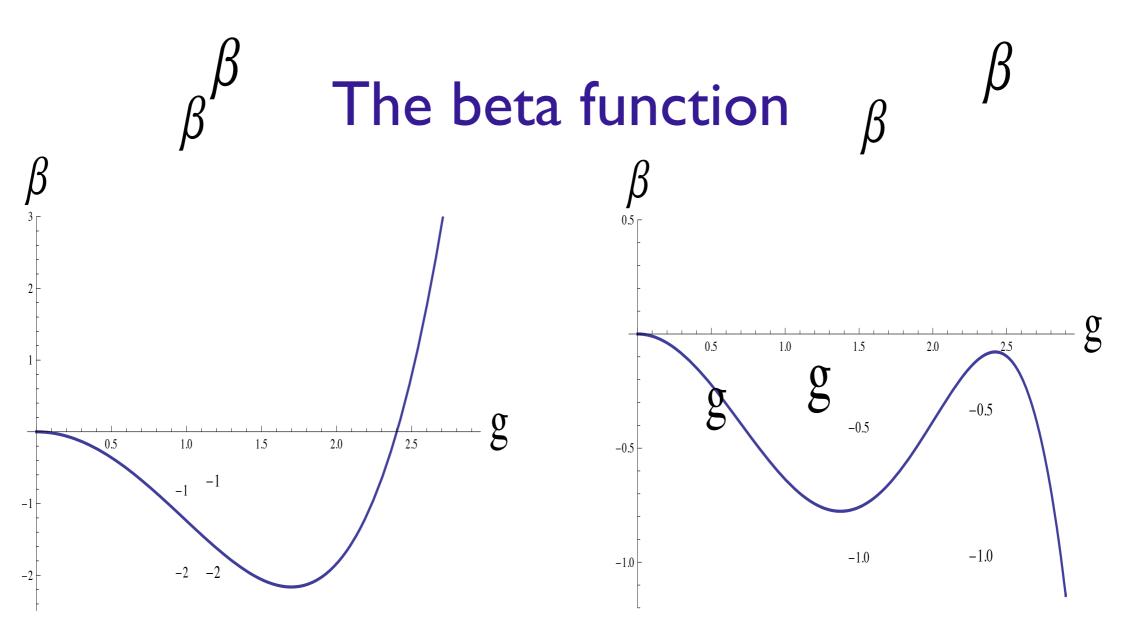


 μ

With large mass anomalous dimension
 Solve the FCNC and S-param problems.

•
$$\Lambda_{\rm ETC}/\Lambda_{\rm TC} \sim 10^2 \sim 10^3$$

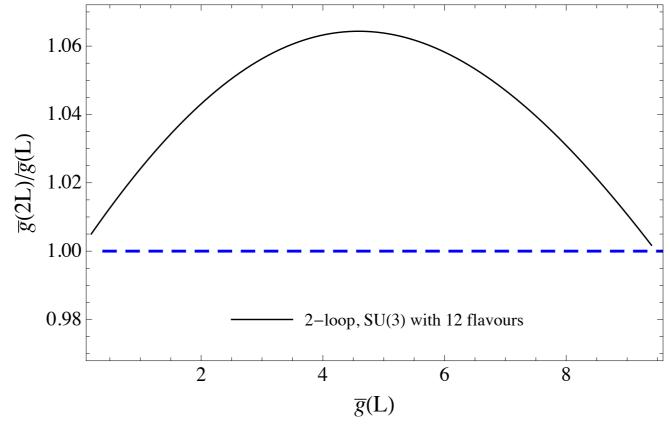
— Compare to typical $L/a \sim 30$.



- Large- N_f gauge theories with asymptotic freedom.
- Need a scale to generate a gap Just below the conformal window $N_f^* < N_f < N_f^{af}$

Lattice strategy for the search of IRFP

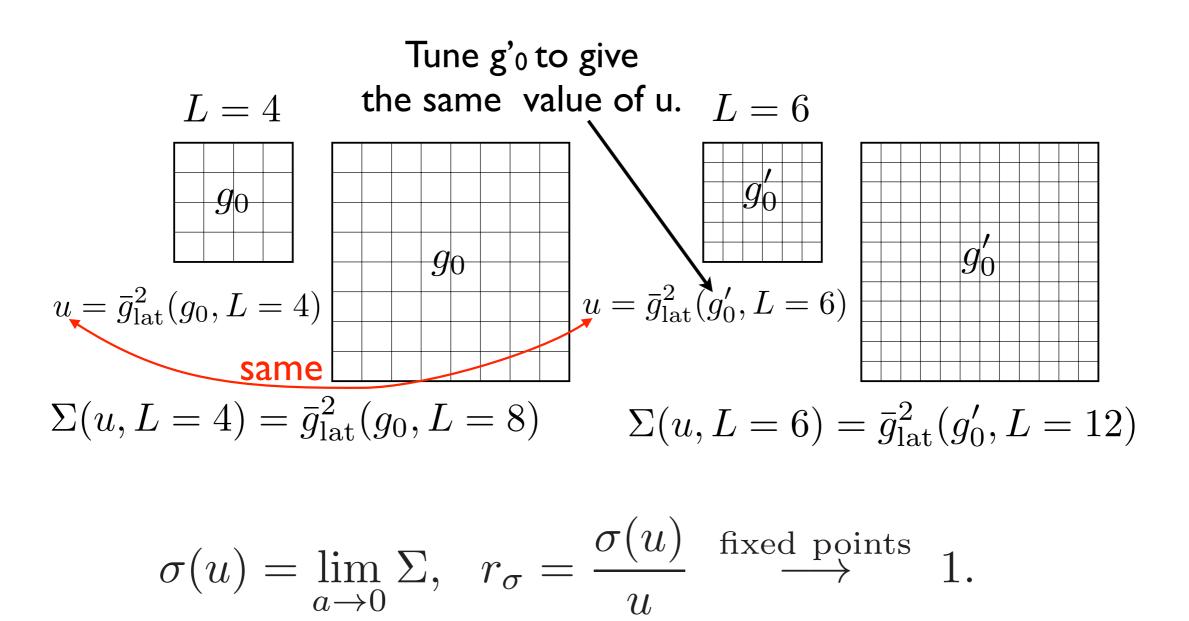
- Spectrum: Large finite-volume effects?
- Finite-size scaling *a'la* M. Fisher : universal curves?
- Running coupling: (slow) running within error?



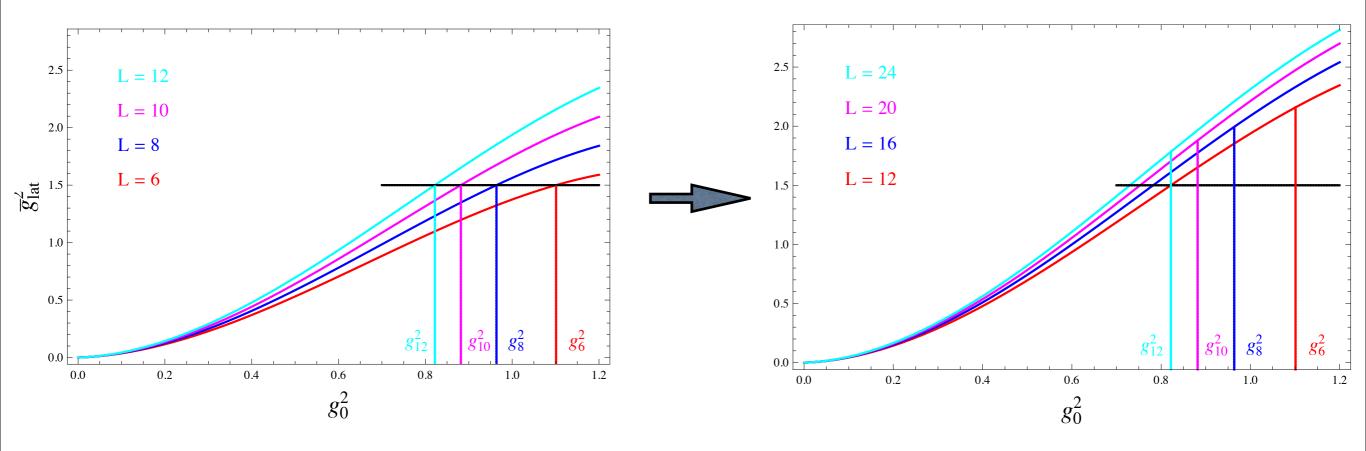
Need high-precision calculations.

The step-scaling method The idea

M. Luscher, P. Weisz, U. Wolff, 1991.



The step-scaling method The practice



Massless unimproved staggered fermions with Wilson's plaquette gauge action.

• Compute \bar{g}_{lat}^2 at many g_0^2 for each volume, and then interpolate volume by volume.

• Very challenging to pin down percentage-level effects in $r_{\sigma} = \frac{\sigma(u)}{\omega}$.

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Bare-coupling interpolation

• Impose the non-decreasing constraint,

$$u_{\text{latt}} = f(u_0) = \int du_0 \left(\sum_{m=0}^{N_{\text{deg}}} c_m u_0^m \right)^2 = \sum_{n=0}^{N_h} h_n u_{0, \mathbf{j}}^n \ u_0 \equiv \frac{1}{\beta} = \frac{g_0^2}{6}$$

in order to avoid the Runge phenomenon.

• Impose the perturbation-theory constraint,

$$h_0 = 0, \ h_1 = 6 \ (\text{then } c_0 = \sqrt{6}).$$

Continuum extrapolation

- Using various polynomials in $\left(\frac{a}{L}\right)^2$.
- Central issue in controlling the systematic error.
- Can we go IR enough before hitting any bulk phase transition?

Twisted box removing the zero modes

• Gauge field:

G.'t Hooft, 1979

 $U_{\mu}(x+\hat{\nu}L) = \Omega_{\nu}U_{\mu}(x)\Omega_{\nu}^{\dagger}, \ \nu = 1, 2,$

where the twist matrices Ω_{ν} satisfy

 $\Omega_1 \Omega_2 = e^{2i\pi/3} \Omega_2 \Omega_1, \ \Omega_\mu \Omega_\mu^\dagger = 1, \ (\Omega_\mu)^3 = 1, \ \mathsf{Tr}(\Omega_\mu) = 0.$

- Fermion: If $\psi(x + \hat{\nu}L) = \Omega_{\nu}\psi(x)$ $\Rightarrow \psi(x + \hat{\nu}L + \hat{\rho}L) = \Omega_{\rho}\Omega_{\nu}\psi(x) \neq \Omega_{\nu}\Omega_{\rho}\psi(x)$
- The fermion "smell" dof: $N_s = N_c$ G. Parisi, 1983 $\psi^a_{\alpha}(x + \hat{\nu}L) = e^{i\pi/3}\Omega^{ab}_{\nu}\psi^b_{\beta}(x)(\Omega_{\nu})^{\dagger}_{\beta\alpha}.$

TPL scheme

- Polyakov loops in the twisted directions:
 P₁(y, z, t) = Tr(Π_jU₁(j, y, z, t)Ω₁e^{2iyπ/3L})
 with gauge and translation invariance.
- The renormalised coupling constant: $g_{\mathsf{TP}}^{2}(L) = \frac{1}{k} \frac{\langle \sum y, zP_{1}(y, z, L/2)P_{1}^{*}(0, 0, 0) \rangle}{\langle \sum x, yP_{3}(x, y, L/2)P_{3}^{*}(0, 0, 0) \rangle},$ where $k = \frac{1}{2k} \sum_{n=1}^{\infty} \frac{(-1)^{n}}{2k} \sim 0.031847$

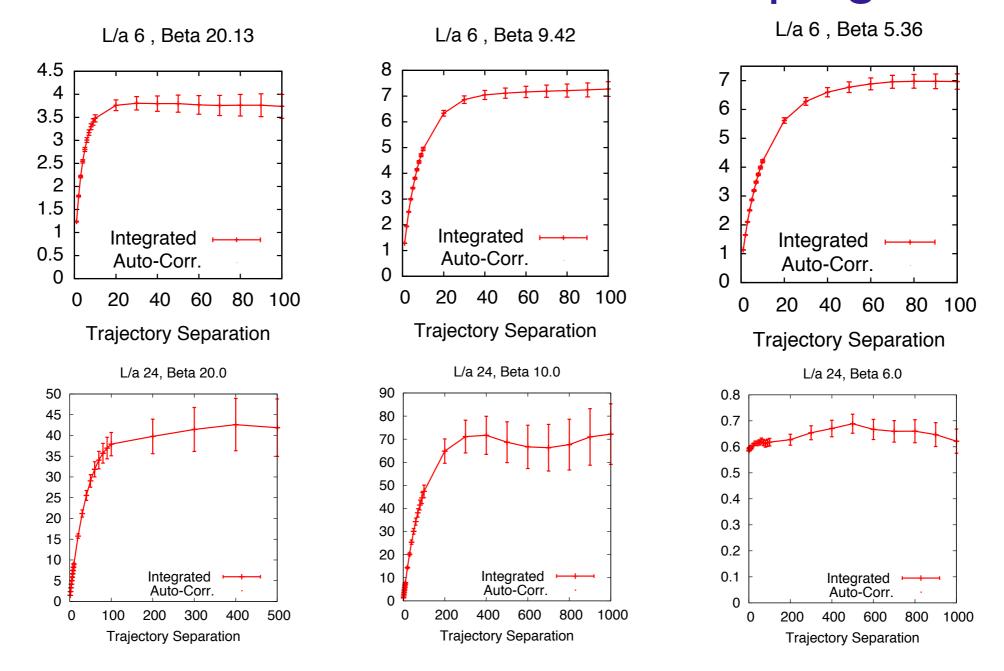
where
$$k = \frac{1}{24\pi^2} \sum_{n=-\infty}^{\infty} \frac{(-1)}{n^2 + (1/3)^2} \sim 0.0318$$

• Special feature:

At
$$L \to \infty$$
, $g_{\text{TP}}^2 \to \frac{1}{k} \sim 32$ if there is no IRFP.

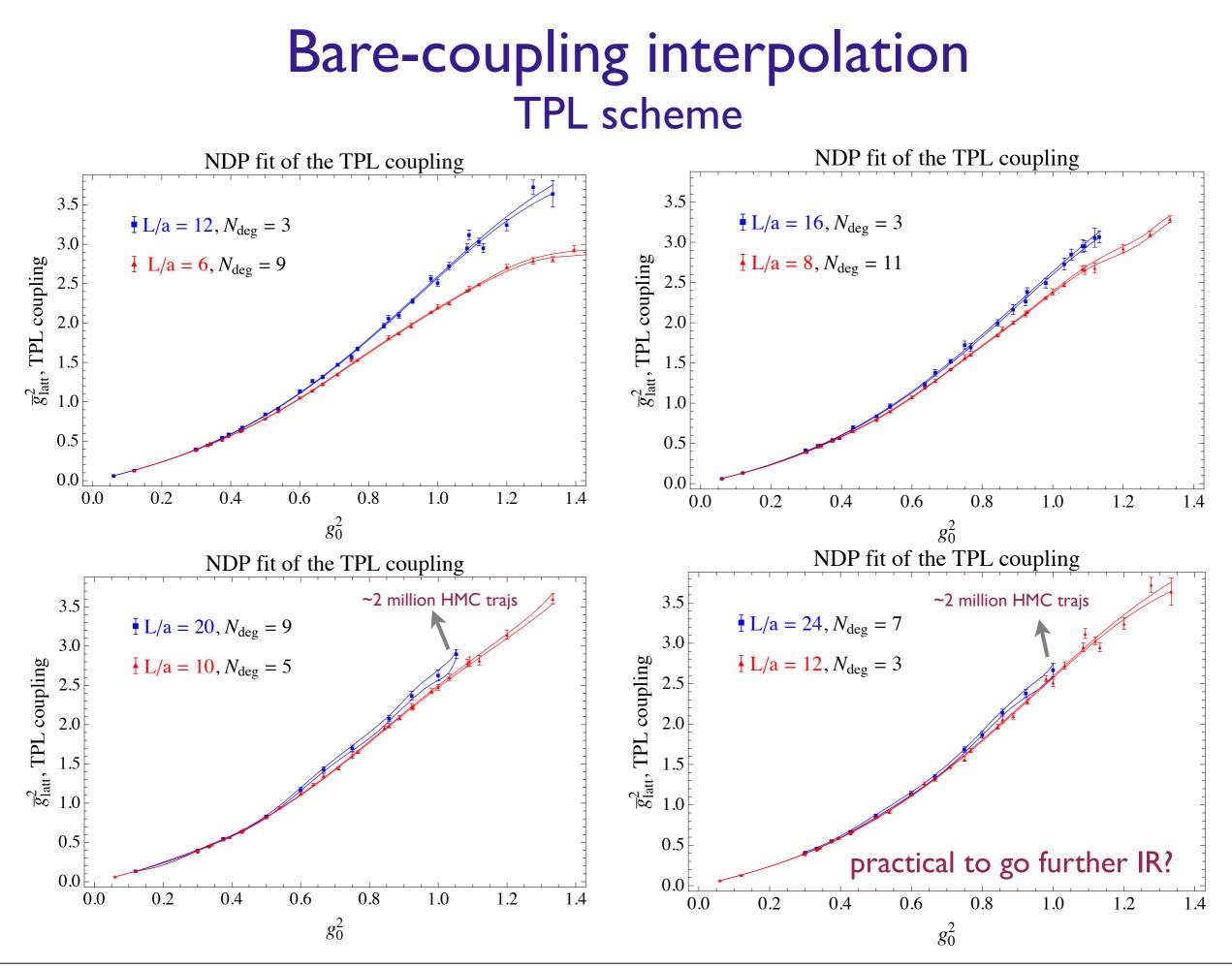
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Challenge in using the TPL scheme Autocorrelation of the coupling

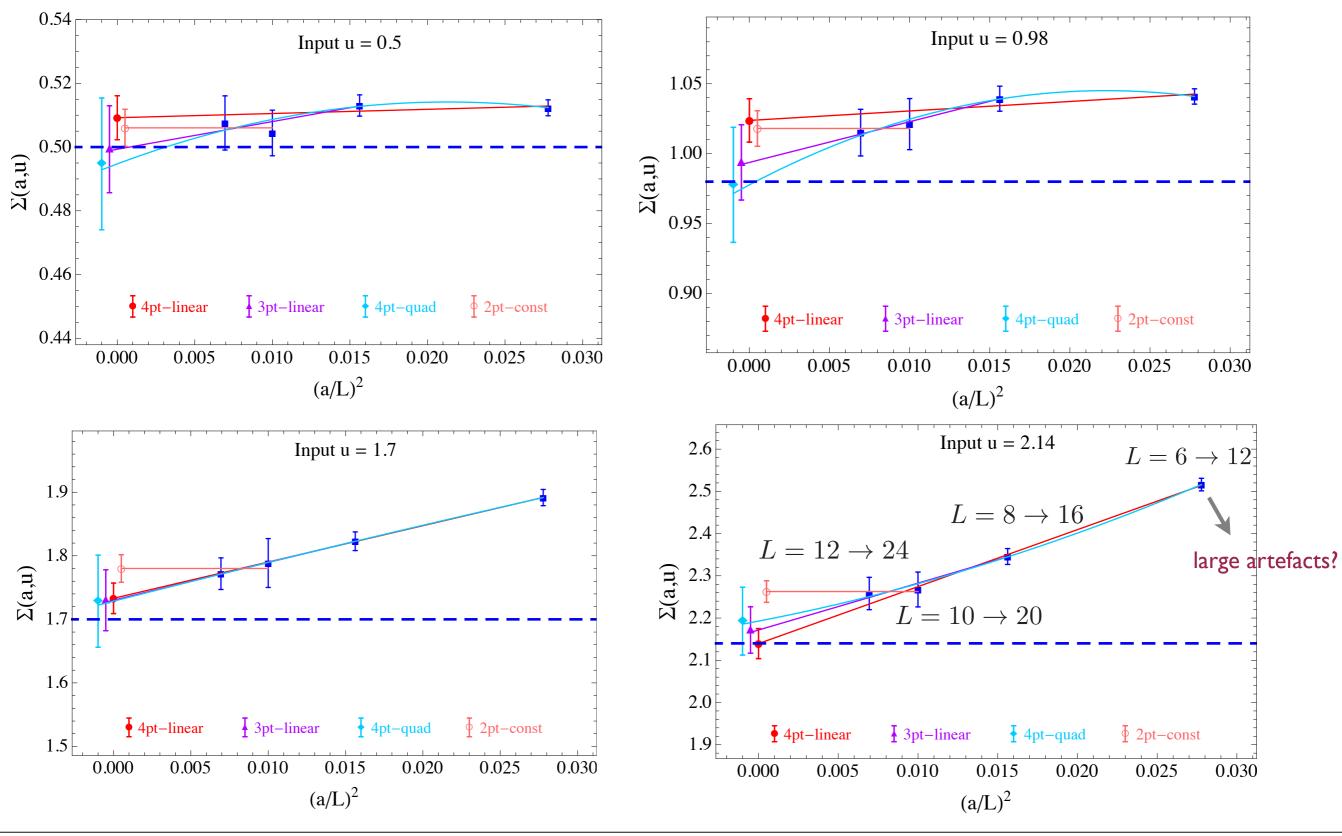


• Autocorrelation time grows with physical volume.

Very challenging to have good statistics for large volumes at low beta.

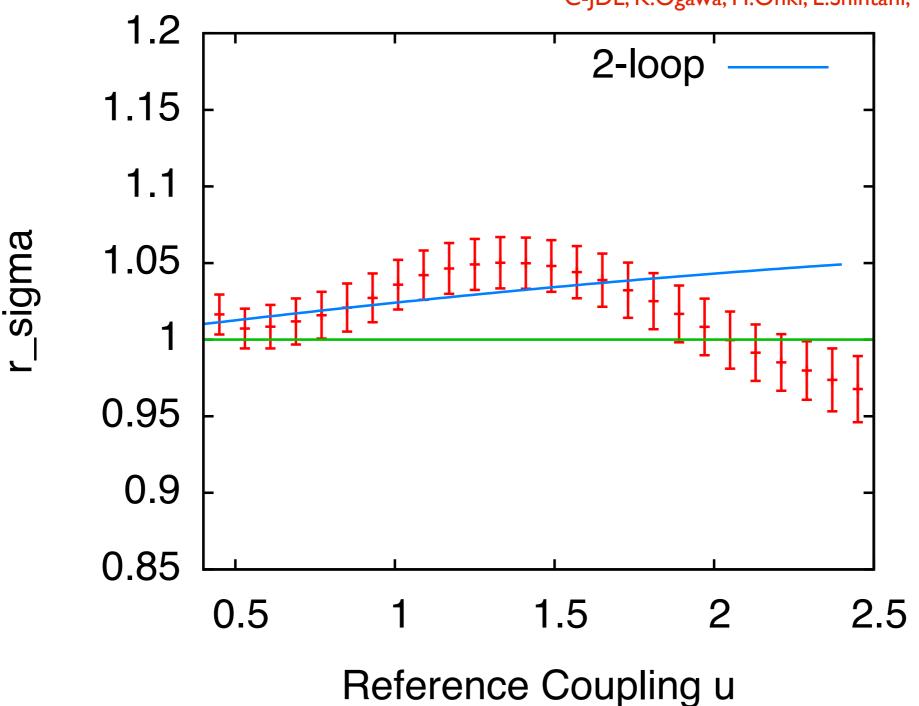


Continuum extrapolation TPL scheme



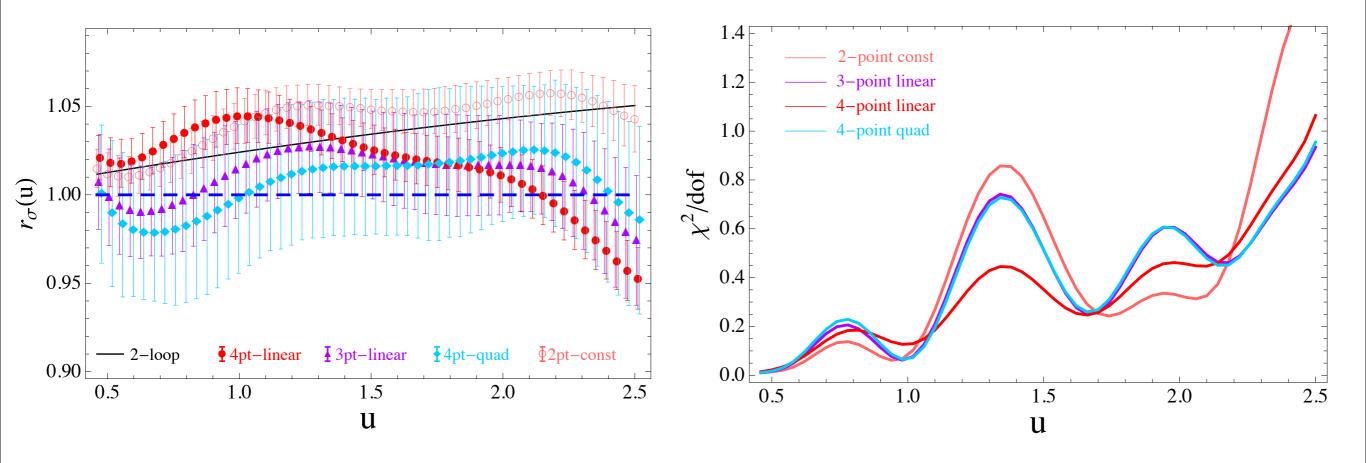
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Result without the L/a=24 lattices TPL scheme



C-JDL, K.Ogawa, H.Ohki, E.Shintani, JHEP 1208 (2012) 096

Result with the L/a=24 lattices TPL scheme



Systematic error was severely underestimated without the L/a=24 data.

The Wilson flow

• Diffusion of the gauge fields:

 $\dot{V}_t(x,\mu) = -g_0^2 \left\{ \partial_{x,\mu} S_w(V_t) \right\} V_t(x,\mu), \ V_t(x,\mu) \big|_{t=0} = U(x,\mu).$

- The radius of diffusion is $\sqrt{8t}$.
- Local operators are also diffused.

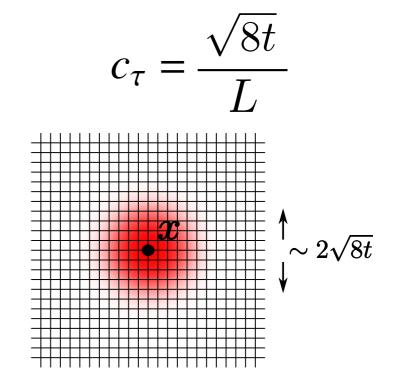


Figure taken from M.Luscher, Lattice 2013

 \mathcal{U}_{\cap}

 $\mathcal{U} \cap$

The Wilson flow scheme

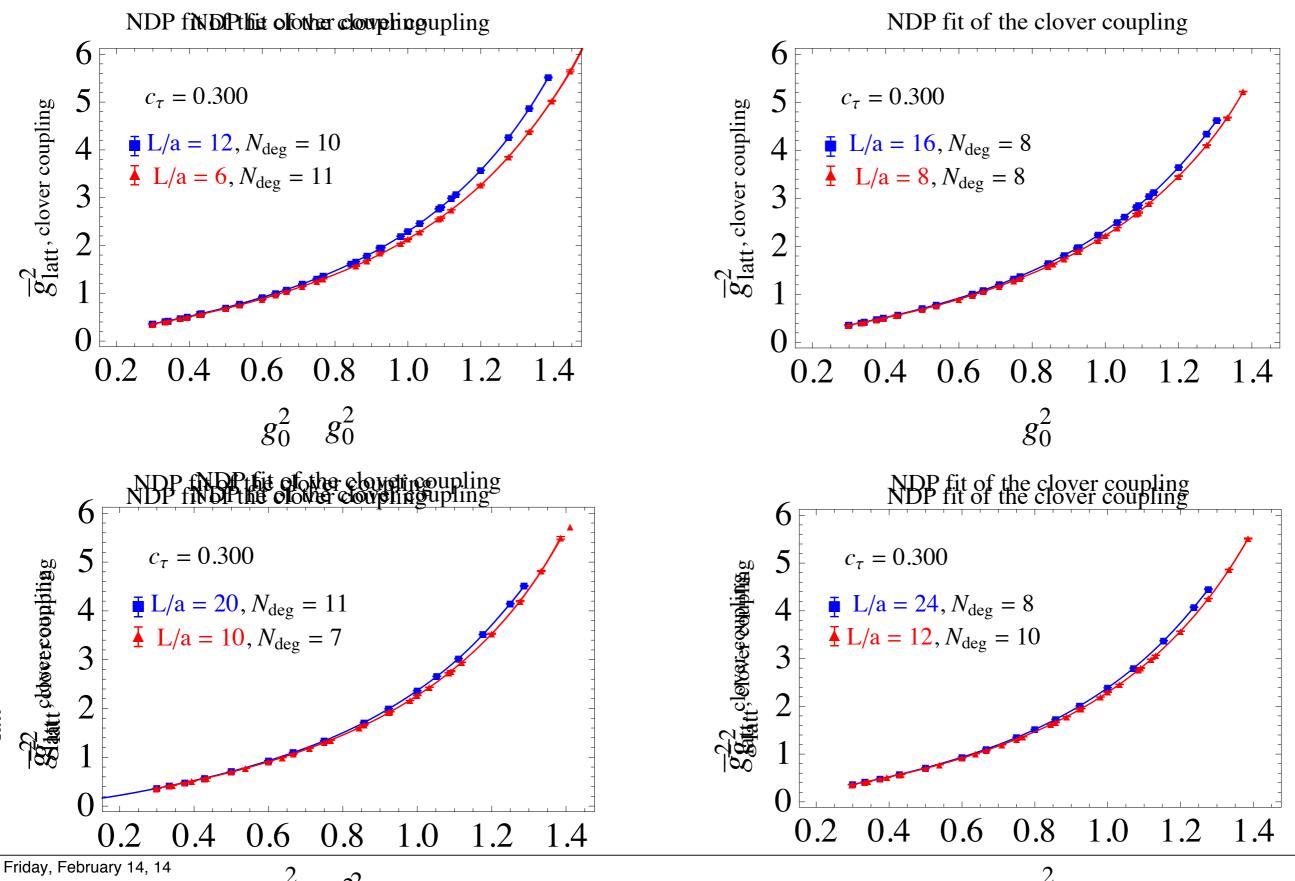
- The quantity, $\langle E(t) \rangle = \frac{1}{4} \langle G_{\mu\nu}(t) G_{\mu\nu}(t) \rangle$, is finite when expressed in terms of renormalised coupling at positive flow time.
- In a colour-twisted box, can define,

$$\overline{g}_{\rm GF}^2(L) = \mathcal{N}^{-1} t^2 \langle E(t) \rangle = \overline{g}_{\rm MS}^2 + \mathcal{O}(\overline{g}_{\rm MS}^4),$$

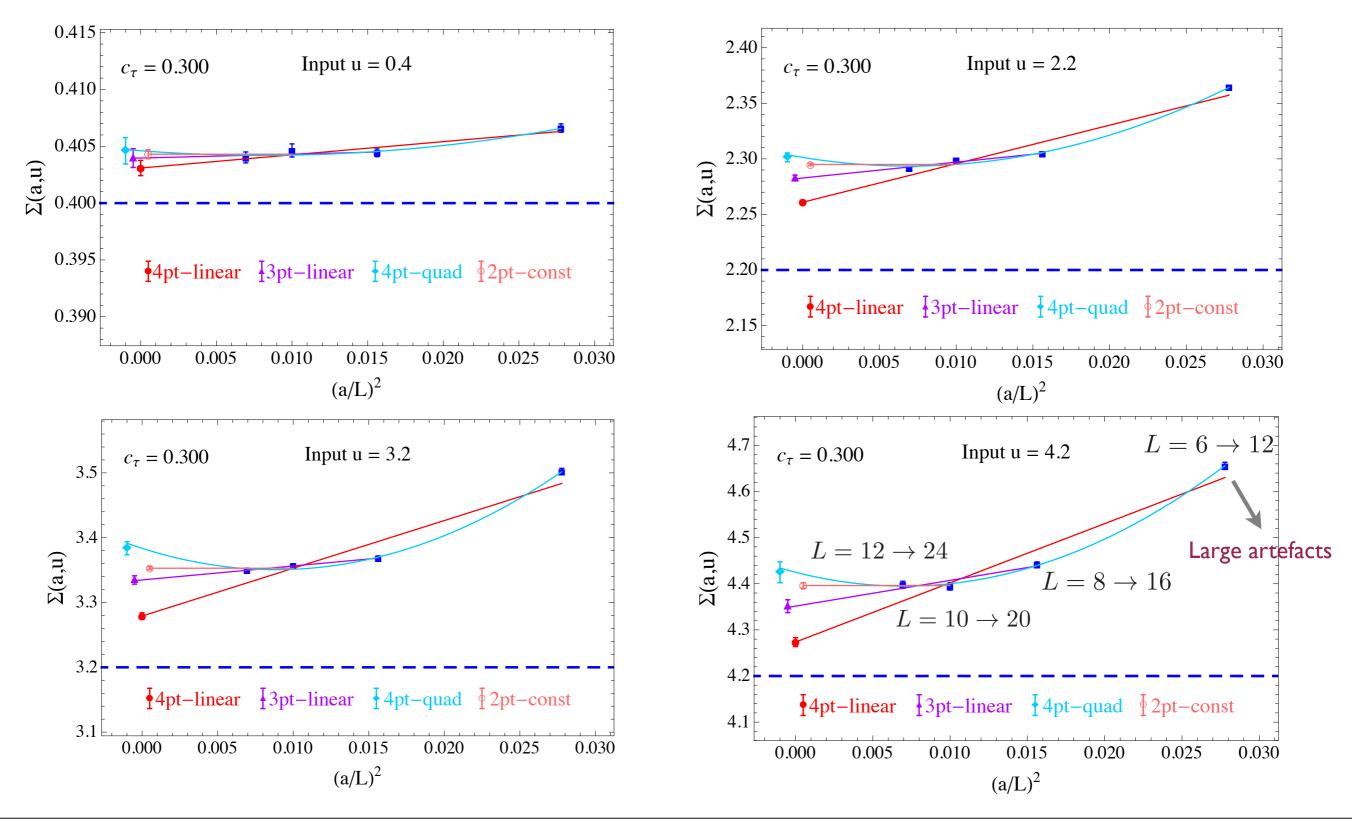
where \mathcal{N} can be computed in perturbation theory.

- Use the clover operator, $\downarrow \downarrow \downarrow \downarrow \downarrow$, to extract $\langle E(t) \rangle$.
- Autocorrelation time ~25 HMC trajectories for all simulations.

Bare-coupling interpolation Wilson flow scheme

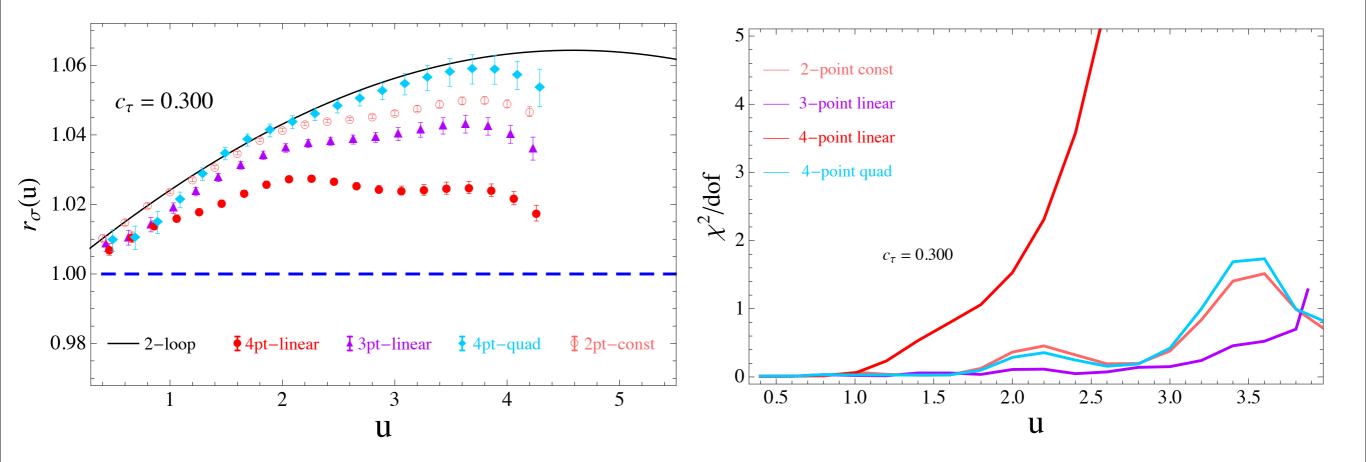


Continuum extrapolation Wilson flow scheme

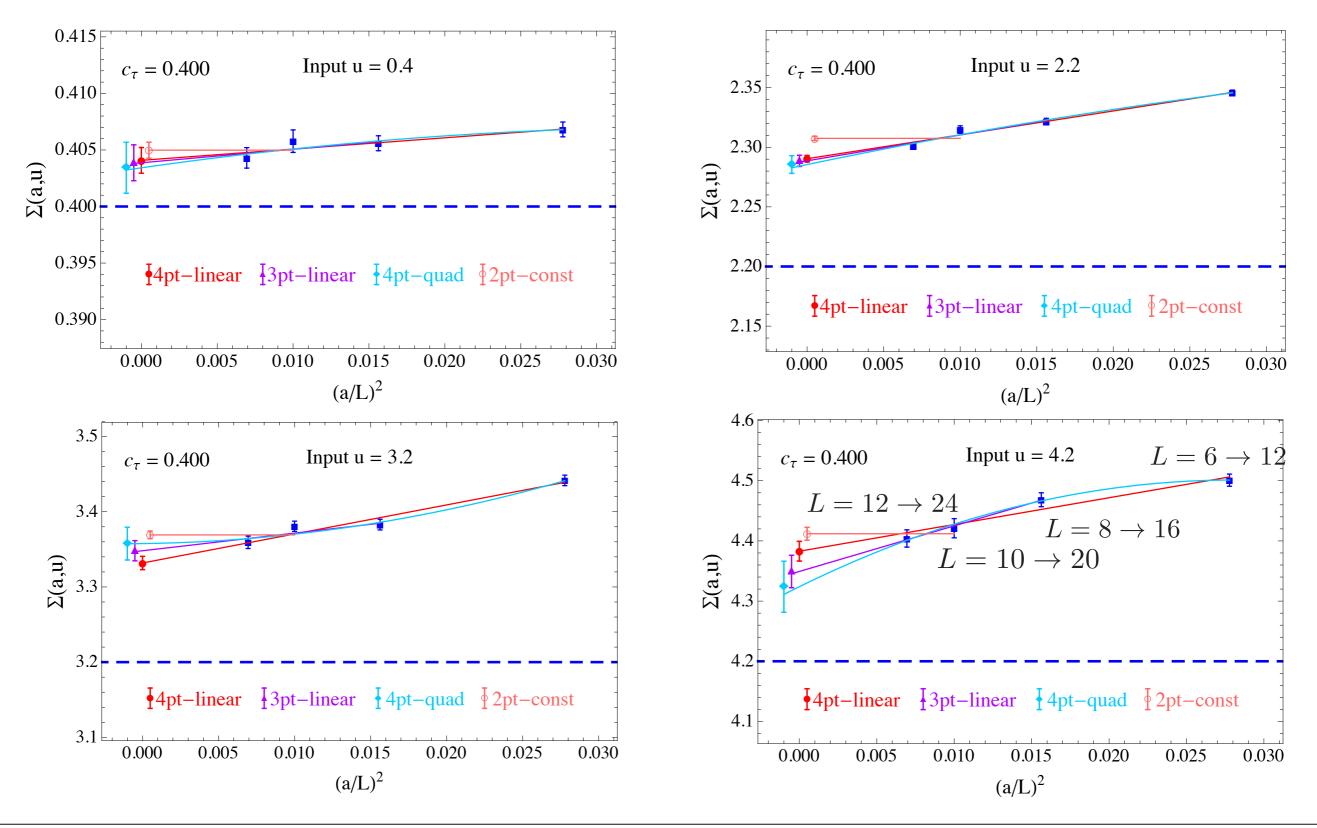


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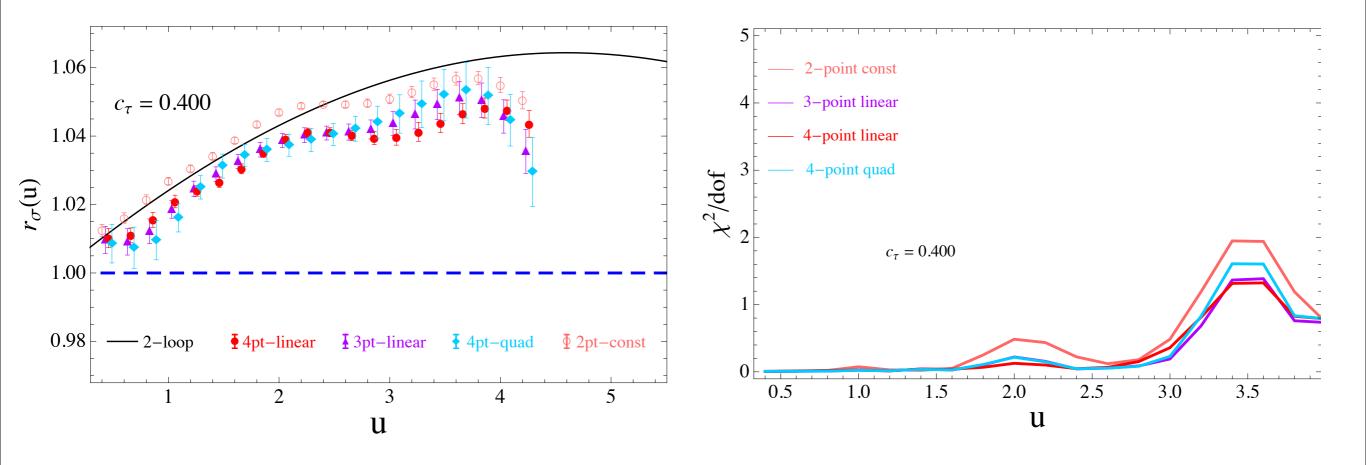
Preliminary result Wilson flow scheme



Continuum extrapolation Wilson flow scheme



Preliminary result Wilson flow scheme

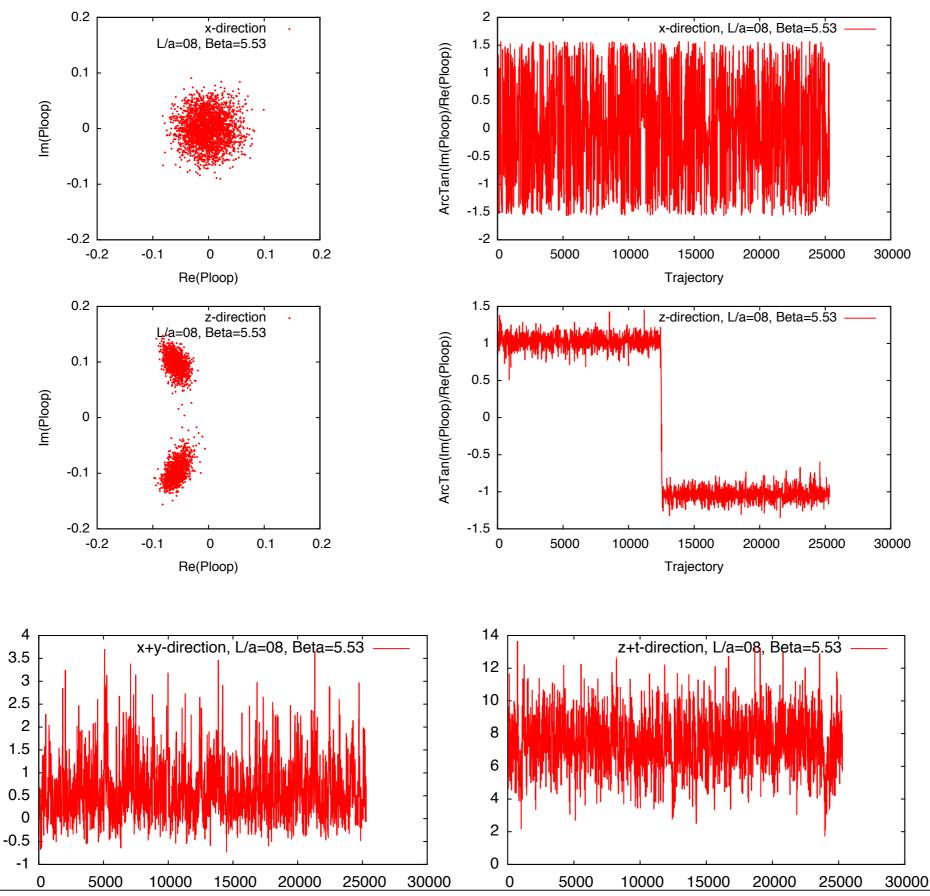


Remarks and outlook

- Calculation in the TPL scheme shows no definite conclusion for IR comformality in SU(3) gauge theory with 12 flavours hitherto.
- It is very challenging to use the TPL scheme to study the evolution of the coupling in the IR.
- On the other hand, the Wilson flow scheme offers a very nice/promising tool.
- We are currently generating data to go further IR

Backup slides

Vacuum structure



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